

**Charge-Shift Bonding in Xenon Hydrides: An NBO/NRT
Investigation on $\text{HXeY} \cdots \text{HX}$ ($\text{Y} = \text{Cl}, \text{Br}, \text{I}$; $\text{X} = \text{OH}, \text{Cl}, \text{Br},$
 $\text{I}, \text{CCH}, \text{CN}$) via H-Xe Blue-Shift Phenomena**

Supporting Information

**Guiqiu Zhang,*¹ Yue Su,¹ Xiaoran Zou,¹ Lei Fu,¹ Junjie Song,¹ Dezhan Chen¹
and Chuanzhi Sun*¹**

Table S1. Calculated bond lengths $R_{\text{H-Xe}}$ and $R_{\text{Xe-Y}}$ (in Å) and monomer-to-complex frequency blue shifts (in cm^{-1}) of H-Xe stretching mode for hydrogen-bonded complexes $\text{HXeY} \cdots \text{HX}$ ($\text{Y} = \text{Cl, Br, I}$; $\text{X} = \text{OH, Cl, Br, I, CN, CCH}$) as well as the angle (in $^\circ$) between HXeY and HX in Structure A at the MP2/def2-TZVPPD level of theory, with available experimental blue shifts. The data in parentheses were calculated at the CCSD(T) level of theory.

Monomers/Complexes	$R_{\text{H-Xe}}$	$R_{\text{Xe-Y}}$	Angle($^\circ$)	H-Xe blue shifts	H-Xe blue shifts (exp)
HXeCl	1.666	2.616			
$\text{HXeCl} \cdots \text{H}_2\text{O}$	1.646	2.671	75	101 (118) ^c	82 ^c
$\text{HXeCl} \cdots \text{HCl}$	1.645	2.672	82	103	116 ^d
$\text{HXeCl} \cdots \text{HBr}$	1.644	2.674	82	102	
$\text{HXeCl} \cdots \text{HI}$	1.646	2.670	82	114	
$\text{HXeCl} \cdots \text{HCN}$	1.648	2.661	92	88	
$\text{HXeCl} \cdots \text{HCCH}$	1.655	2.642	79	51	
HXeBr	1.679	2.774			
$\text{HXeBr} \cdots \text{H}_2\text{O}$	1.656	2.825	71	116 (148) ^c	101 ^c
$\text{HXeBr} \cdots \text{HCl}$	1.657	2.822	77	110	122 ^d
$\text{HXeBr} \cdots \text{HBr}$	1.656	2.823	77	110	149 ^d
$\text{HXeBr} \cdots \text{HI}$	1.658	2.819	78	90	
$\text{HXeBr} \cdots \text{HCN}$	1.658	2.817	88	106	
$\text{HXeBr} \cdots \text{HCCH}$	1.667	2.798	75	58	
HXeI	1.708	2.976			
$\text{HXeI} \cdots \text{H}_2\text{O}$	1.679	3.022	67	139 (188) ^c	138 ^c
$\text{HXeI} \cdots \text{HCl}$	1.682	3.017	72	122 (167) ^b	94, 111, 155 ^b
$\text{HXeI} \cdots \text{HBr}$	1.681	3.017	72	122 (149) ^a	110, 157 ^a
$\text{HXeI} \cdots \text{HI}$	1.683	3.012	72	107 (120) ^a	75, 96 ^a
$\text{HXeI} \cdots \text{HCN}$	1.680	3.015	83	132	
$\text{HXeI} \cdots \text{HCCH}$	1.693	2.996	70	67 (104) ^b	49, 55 ^b

^a From ref. Tsuge et al., 2013.

^b From ref. Zhu et al., 2015.

^c From ref. Tsuge et al., 2014.

^d From ref. Lignell et al., 2008.

Table S2. Second-order perturbation energies($E^{(2)}$, kcal mol⁻¹) due to donor-acceptor interactions in HXeY at the MP2/def2-TZVPPD level of theory.

Molecules	$n_Y \rightarrow \sigma^*_{H-Xe}$	$n_H \rightarrow \sigma^*_{Xe-Y}$	$n_{Xe} \rightarrow \hat{\sigma}^*_{H-Y}$
HXeCl	64.57	3957.00	1363.55
HXeBr	65.75	5187.90	1102.19
HXeI	70.75	5125.55	845.16

Explanation of the Procedure Employed to Calculate the BO of the Xe-H Bond

1. Conceptual Model in NBO/NRT

In the NBO/NRT framework, the fundamental starting point for a rational electronic theory of bonding is the Lewis-structure representation of the shared and unshared electrons in each atomic valence configuration, as formulated by Lewis. Based on the Lewis-structure model, perturbation theory is used to calculate systematically the corrections that bring the Lewis-structure model into an improved Lewis-structure model. In NBO/NRT language, it is the natural Lewis-structure model.

Subsequently, bonding analysis can be dissected into localized and delocalized contributions.

Take HXeY as an illustrative example. HXeY could be described as the hybrid of three resonance structures: $\text{H}-\text{Xe}^+ \text{Y}^-$, $\text{H}^- \text{Xe}^+-\text{Y}$, and H^+Y^- . When we focus on the H-Xe bond, the structure I provides a localized contribution due to electron-sharing bonding. In structure II, it is a delocalized contribution arising from a donor-acceptor interaction. In chemical language, it is dative bonding. Structure III is a long-bonding structure. There is no contribution to the H-Xe bond strength, because there are neither localized contributions nor delocalized contributions to the H-Xe bonding.

2. Calculate the Weighting of Resonance Structure

A quantitative resonance theory can help us to find the weighting of resonance structure α , according to

$$D(\text{true}) = \sum_{\alpha} \omega_{\alpha} D\alpha^{(L)}$$

where $D(\text{true})$ is the true density matrix of the system of interest. $D\alpha^{(L)}$ is corresponding to the density matrix of the resonance structure α . Note that $\sum_{\alpha} \omega_{\alpha} = 1$.

3. Calculate Each Property of the System

In the NRT of framework, each property $\langle P \rangle_{\text{true}}$ of the true delocalized system can be represented in resonance-averaged form

$$\langle P \rangle_{\text{true}} = \sum_{\alpha} \omega_{\alpha} \langle P \rangle_{\alpha}$$

where $\langle P \rangle_{\alpha}$ is the value of the property for natural resonance structure α .

Still take HXeY as an example. The H-Xe bond strength $D(\text{H-Xe})$ can be written as $D(\text{H-Xe}) = \omega_{\text{I}} D_{\text{I}} + \omega_{\text{II}} D_{\text{II}} + \omega_{\text{III}} D_{\text{III}}$.

From the conceptual analysis of the natural Lewis-structure, we know $D_{\text{III}} = 0$. Thus, $D(\text{H-Xe}) = \omega_{\text{I}} D_{\text{I}} + \omega_{\text{II}} D_{\text{II}}$

4. Calculate Bond Order of H-Xe

When using the bond order to reflect the H-Xe bond strength, D_{I} and D_{II} can be written as $D_{\text{I}} = k_{\text{I}} \omega_{\text{I}} b_{\text{I}}$, $D_{\text{II}} = k_{\text{II}} \omega_{\text{II}} b_{\text{II}}$, where b_{I} represents the normal covalent H-Xe bond order in resonance structure I, b_{II} is the H-Xe dative bond order in resonance structure II.

If $k_{\text{I}} = k_{\text{II}} = k$, and $b_{\text{I}} = b_{\text{II}} = 1$, we obtain $b(\text{H-Xe}) = \omega_{\text{I}} + \omega_{\text{II}}$ for our studied HXeY.

It is worthwhile noting that the present NBO/NRT methods could not provide the ratio of $k_{\text{I}} / k_{\text{II}}$, and that the adequacy in calculating the BO of the H-Xe bond has been tested through indirect comparisons with experimental results.