



Tailoring the Band Gap in Codoped GaN Nanosheet From First Principle Calculations

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To develop an efficient photocatalyst with intense visible light absorption and high charge mobility is important but still remains a problem. In this work, we have explored the electronic properties of C-monodoped, C-Ge, and C-Sn codoped GaN nanosheets by the hybrid density functional theory in order to find the excellent photocatalytic materials. Results indicate the monodoping of C introduces unoccupied impurity states inside the band gap that serve on recombination centers. Thus, the C monodoping is not suitable to ameliorate visible light absorption. Moreover, the C-Ge and C-Sn codoping not only successfully reduce the band gap of nanosheet GaN but also avoid the unoccupied impurity states. The charge-compensated C-Ge and C-Sn codoped GaN nanosheets are energetically favorable for hydrogen evolution but not insufficient to produce oxygen, indicating that they could serve as Z-scheme photocatalysts. In particular, the minimum defect formation energy of C-Ge is negative and lowest. The C-Ge codoped GaN is one of the most prospective candidates for the decomposition of hydrogen from water.

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1. INTRODUCTION

Gallium nitride (GaN) has attracted considerable interest in recent years for its excellent photoelectric properties. With a large band gap (3.04 eV) (Bastos et al., 2018), GaN is widely used as short-wavelength light-emitting diodes (Ha et al., 2007) and room-temperature laser diodes (Feng et al., 2017). The conduction band edge of bulk GaN is at 0.5 eV higher than the redox potential of H⁺/H₂, which implies that the GaN can decompose hydrogen from water (Maeda et al., 2005; Kida et al., 2006). Its wide band gap, however, makes GaN photochemically active only under UV light irradiation. In order to extend the visible light response of GaN, some efforts have been made based on band gap engineering strategies. It was found that Cr-O codoping significantly narrows the band gap of GaN, but this leads to an intermediate band entering the band gap (Pan et al., 2010), which increases the probability of carrier capture. Similar results were obtained for V-O codoped GaN (Meng et al., 2012).

According to the recent theoretical calculations (Freeman et al., 2006), the monolayer GaN has a honeycomb structure similar to that of graphene (Castro Neto et al., 2009), and boron nitride nanosheets (Golberg et al., 2010). These two-dimensional (2D) atomic-layer nanomaterials are interesting topics for their picturesque properties, which would have advantages over bulk structures in many cases (Onen et al., 2016). Particularly, the 2D GaN has a large specific surface area and shorter carrier migration distance etc., and it can remain in high-temperature

1

environments (Sun et al., 2017). The unique properties of nanosheets GaN motivate our research into whether a tunable band gap could be obtained via chemical modification. In this work, the C-monodoping and C-Ge and C-Sn codoping were used to modulate the band structure of GaN nanosheet. We focused our studies on the structure, electronic properties, and minimum defect formation energies of the doped GaN nanosheets.

2. COMPUTATIONAL METHOD

Our first-principles calculations were executed by using density functional theory (DFT) within generalized gradient approximation (GGA) (White and Bird, 1994) parameterized by Perdew, Burke, and Ernzerhof (PBE) (Ernzerhof and Scuseria, 1999), as realized in VASP code (Kresse and Furthmüller, 1996; Kresse and Joubert, 1999). Projector augmented wave (PAW) (Blöchl, 1994) potentials were employed to treat interactions between ions and valance electrons. The electron configuration of N (2s²p³), Ga (3d¹⁰4s²4p¹), C (2s²2p²), Ge (3d¹⁰4s²4p²), and Sn (5s²5p²) were chosen to construct the potential function, respectively. All the calculations were carried out with a planwave cutoff energy of 550 eV. We adopted vacuum space of 20 Å along the z-direction in order to wipe out the interaction between the periodic images. All the atomic structure energy convergence tolerance for self-consistent iteration was set to 1.0×10^{-5} eV, and the convergence criteria force was set to 0.01 $eVÅ^{-1}$. The Brillouin zone integrations were approximated with a 9 \times 9 \times 1 k-point (Monkhorst and Pack, 1976) for the GaN nanosheet structure relaxation and electronic properties calculations. The GGA-PBE algorithm will underrate the band gap value of the semiconductor. The accurate electronic structure was calculated by the Heyd-Scuseria-Ernzerhof (HSE06) (Heyd et al., 2003; Heyd and Scuseria, 2004) hybrid functional method. The HSE06 expression for the exchange-correlation energy was written a:

$$E_{XC}^{HSE} = \chi E_X^{SR}(\mu) + (1 - \chi) E_X^{PBE,SR}(\mu) + E_X^{PBE,LR}(\mu) + E_C^{PBE}$$
(1)

The exchange-correlation energy includes the short- and longranged parts. Here, the SR is the short-range parts, and LR stands for the long-ranged. χ represents the mixing coefficient and is set to 0.25. μ indicates the screening parameter and is set to 0.2. In addition, the total density of states (TDOS) and the projected density of states (PDOS) are calculated by the tetrahedral method with Blöchl correction (Blöchl et al., 1994).

3. RESULTS AND DISCUSSION

3.1. Pure GaN Nanosheet

The GaN nanosheet is the (001) facet of the optimized bulk, which has a honeycomb graphite structure. This structure is shown in **Figures 1A,B**. The Ga-N bond length is 1.849 Å and is smaller than in the bulk GaN case (1.966 Å) (Sahin et al., 2009). Since the sp² bond is stronger than the sp³ bond, the bond length in the GaN nanosheet becomes shorter. The **Figure 1C** displays the band structure of the GaN nanosheet with the band gap of 3.41 eV using the HSE06 functional. The VBM is located

at the k point while the CBM is at the Γ point, which means that the GaN nanosheet is the indirect band gap semiconductor. According to the Figure 1D, the GaN nanosheet has nonmagnetic semiconducting character, and the VBM and CBM are predominantly derived from N 2p, Ga 4p, and Ga 3d orbits. The hybridization between the N 2p, Ga 4p, and Ga 3d orbits are located at the bottom of the conduction band. The bond length d_{Ga-N} , lattice constant *a*, and band gaps calculated by the PBE and HSE06 functionals are given in Table 1. Our theoretical data have exhibited good agreement with previous studies (Maruska and Tietjen, 1969; Armstrong et al., 2005; Zevan et al., 2008; Sahin et al., 2009; Xia et al., 2013; Bastos et al., 2018), which indicates that the calculation methods in this study are reasonable. It can be found that the PBE algorithm will underestimate the band gap of GaN nanosheet. Therefore, the HSE06 algorithm is adopted to study the electronic structures of the GGA-PBE optimized phases.

3.2. Monodoping in the GaN Nanosheet

In this section, we studied the effect of C doping on the electronic properties of the nanosheet GaN. For $GaN_{1-x}C_x$, the nanosheets (with x = 0.125) were modeled by replacing aneutral N atom (A site in **Figure 1A**) with a C atom in the $2 \times 2 \times 1$ supercell of the GaN nanosheet. Because the charge number of the doping system is less than that of the pure system, we considered the spin-polarization effect in the calculation.

For the C modoped GaN nanosheet, the bond distance between the C atom and the nearest-neighboring Ga atom is 1.911 Å, which is close to the bond distance (1.849 Å) of Ga-N in nanosheet GaN. The minor difference may be explained by the ionic radius difference between C and N atom. The DOS and band structure of GaN nanosheets with the C doping presented in Figures 2A,B. For both C monodoped systems, the VBM is mainly contributed by mixing of the N 2p, C 2p, Ga 4p, and Ga 3d states, whereas the CBM is primarily contributed by the mixing of the between Ga 4p, N 2p, Ga 3d, and C 2p states. C dopant introduces the unoccupied impurity states on the spin down direction. The impurity state formed by a strong hybrid of C 2p and Ga 4p is very close to the conduction band of GaN. The generation of the impurity state may be due to the fact that the C and N are elements of the same cycle, and the porbital energy of C is thus higher than that of N. Electrons in the valence band can easily migrate to the conduction band through these impurity states. Although the band gaps for C doping GaN nanosheet decrease to 1.84 eV, the monodoped system has partially unoccupied impurity states at top of Fermi level. These impurity states can facilitate electron-hole recombination and have destructive influence on the photocatalytic activity. It indicates that C monodoped GaN nanosheet is not suitable for a visible light photocatalyst.

3.3. Codoping in the GaN Nanosheet

To overcome the unoccupied impurity states, the anion-cation passivated codoping was utilized to maintain the charge balance (Gai et al., 2009). In this section, we have discussed the influence of codoping C-Ge and C-Sn on the electronic properties of the supercell $2 \times 2 \times 1$ GaN nanosheet. The C atom substitutes





TABLE 1 | The obtained lattice constant a, bond length d_{Ga-N} , and band gap E_g with the PBE and HSE06 functional.

Structure		a (Å)	d _{Ga-N}	HSE06- E_g (eV)	PBE- <i>E</i> _g (eV)
Nanosheet	This work	3.202	1.849	3.41	2.44
	Other work	3.216 (Xia et al., 2013)	1.851 (Xia et al., 2013)	3.44 (Armstrong et al., 2005)	2.27 (Sahin et al., 2009)

the N atom (A site in **Figure 1A**), and the Ge (or Sn) atom substitutes the Ga atom (B site in **Figure 1A**), respectively. For C-Ge and C-Sn codoping, the N-Ge bond length is 1.829 Å, and the N-Sn bond distance is 2.00 Å, decreasing 0.02 Å, and increasing 0.151Å, respectively, with comparison with the N-Ga bond length (1.849 Å). The C-Ge bond length is 1.810 Å and the C-Sn bong length is 1.986 Å. The DOS and band structure

for the C-Ge are plotted in **Figure 3**. According to the DOS, we found that the VBM is dominated by C 2p, Ga 4p, N 2p, and Ga 3d states, whereas the CBM is dominated by N 2s, Ge 4s, Ga 4s, and N 2p states. Since the band structures display that the VBM and CBM are located at the Γ point, the C-Ge system is the direct band gap semiconductor. The calculated band gap of C-Ge codoped GaN nanosheet is 1.83 eV. The





DOS of C-Sn doping is plotted in **Figure 4B**. It indicates that the VBM is dominated by C 2p, Ga 4p, N 2p, and Ga 3d, whereas the CBM is dominated by Sn 5s, N 2p, Ga 4p, and Sn 4d. These configurations have strong hybridization at the top of the valence band and at the bottom the conduction band. According to **Figure 4A**, the calculated band gap of C-Sn codoped GaN nanosheet is 1.85 eV. The interesting phenomenon is that the impurity states are completely occupied for codoped GaN nanosheets.

As discussed above, those results indicate that the nanosheet GaN changes from an indirect band gap to a direct band gap by C-Ge and C-Sn codoping. The nanosheet GaN band gap decreases from 3.41 to 1.83 eV in the C-Ge codoping case and to 1.85 eV in the C-Sn codoping one. Smaller band gap values and a direct band gap are beneficial for excited state electronics, as they allow for greater ease with which to jump to the conduction band. At the same time, the smaller band gap value is favorable for moving the absorption spectrum to the visible light region. Therefore, the C-Ge and C-Sn codoped nanosheets GaN are likely to remarkable candidates for visible light photocatalyst.

3.4. Defect Formation Energy

We have calculated the defect formation energy (E_f) for evaluating the stability of doping system in each case. The defect formation energy is defined as (Van de Walle and Neugebauer, 2004; Castleton et al., 2006):

$$E_f = E(doped) - E(GaN) - \sum_i n_i \mu_i$$
(2)

Here, E(doped) and E(GaN) indicate the total energies of doped and pure nanosheet GaN, respectively. n_i represents the number of the atoms that has been added to $(n_i>0)$ or removed from $(n_i<0)$ the pure system for the construction doped nanosheet GaN, while μ_i is the chemical potential of the element *i*, and its value μ_i rests with the synthetic environment

$$\mu_{Ga} + \mu_N = \mu_{GaN(nanosheet)} \tag{3}$$

Here, $\mu_{GaN(nanosheet)}$ represents the chemical potential, which is equal to the energy of the unit cell. As the chemical potential of



Ga and N cannot exceed that of the Ga bulk $\mu_{Ga(bulk)}$ and N gas $\mu_{N(gas)}$, the following rules must be obeyed

$$\mu_{Ga}^{min} \le \mu_{Ga} \le \mu_{Ga(bulk)} \tag{4}$$

$$\mu_N^{\min} \le \mu_N \le \mu_{N(gas)} \tag{5}$$

To form the nanosheet GaN naturally, the minima of μ_{Ga} and μ_N will satisfy

$$\mu_{Ga}^{min} = E(Ga_n N_n) - E(Ga_{n-1} N_n) \tag{6}$$

$$\mu_N^{min} = E(Ga_n N_n) - E(Ga_n N_{n-1}) \tag{7}$$

Here, $E(Ga_nN_n)$ is the total energy of perfect system with n primitive cells, $E(Ga_{n-1}N_n)$ the total energy with Ga defect, and $E(Ga_nN_{n-1})$ the total energy with N defect. We set the $\Delta\mu_{Ga}=\mu_{Ga}-\mu_{Ga(bulk)}$ and $\Delta\mu_N=\mu_N-\mu_{N(gas)}$. In this article, only the minimum defect formation energies of doped GaN nanosheets have been estimated; e.g., the defect formation energy take the minimum value when μ_{Ga} equals $\mu_{Ga(bulk)}$. The minimum formation energies for doping GaN nanosheets are shown in **Figure 5**, and the results are outlined in **Table 2**.

From the above discussion, we know that the minimum defect formation energy for the C monodoping is positive, which means that the C monodoping is more difficult to form because the electronegativity of C is less than that of N. The C-Ge and C-Sn codoping minimum defect formation energies are significantly lower than C monodoping. Importantly, the minimum defect formation energy of the C-Ge codoped is negative and lowest, which is greatly reduced by the opposite charge states of C-Ge atoms.

3.5. Band Edge Alignment

It is well-known that an excellent photocatalyst should not only have an appropriate band gap but also suitable band edge positions, i.e., The water reduction level and oxidation level must

TABLE 2 | The calculated band gaps (E_{gap}), minimum defect formation energies (E_{f}), and chemical potentials of N and Ga.

Structure	E _{gap} (eV)	E _f (eV)	${\bigtriangleup \mu_{Ga}}$ (eV)	${\bigtriangleup \mu_{N}}$ (eV)	
C(x = 0.125)	1.84	1.39	-0.36	-0.36	
C-Sn	1.85	0.55	-0.36	-0.36	
C-Ge	1.83	-1.71	-0.36	-0.36	



lie between the CBM and VBM in order to thermodynamically allow for both the hydrogen evolution reaction and oxygen evolution. On the basis of the band edge positions, some photocatalysts are suitable for only hydrogen evolution, and some photocatalysts are good for only oxygen evolution, which is called the Z-scheme photocatalysis system. Here, we have considered the situation of the standard water reduction level (-4.44 eV) and oxidation level (-5.67 eV) with respect to the vacuum level (Chakrapani et al., 2007). A large vacuum space has been used to calculate the absolute position of the band edge for quantitatively comparing with the water redox potentials. As the calculated vacuum level is a relative value, we have shifted the energy value of the CBM (VBM) for each system by subtracting this value to obtain the absolute band edge position.

The calculated positions of the CBM and VBM of the pure and doped GaN nanosheets with respect to the standard hydrogen electrode potentials are indicated in **Figure 6**. Compared to the band edge potentials of pure GaN, both CB and VB edges of codoped systems are increased. The calculated CBM and VBM of C-Ge (C-Sn) codoped nanosheet are located at -3.50 eV (-3.55 eV) and -5.33 eV (-5.40 eV), respectively. The CBM of C-Ge (C-Sn) is 0.90 eV (0.85 eV) above the water reduction level, and the VBM of C-Ge (C-Sn) is 0.34 eV (0.27 eV) above the water oxidation level. Thus, the C-Sn and C-Ge codoped GaN nanosheets are energetically favorable for hydrogen evolution but not insufficient to produce oxygen, indicating that they could serve as Z-scheme photocatalysts.

3.6. Optical Properties

Optical properties play a critical role in evaluating the performance of photocatalysts in the visible light region. Thus,



FIGURE 7 | The optical absorption spectra of pure and codoped GaN nanosheets.



we have investigated the optical absorption spectra of the pure and codoped GaN nanosheets. The absorption coefficient $\alpha(\omega)$ satisfies the following relation (Onen et al., 2016)

$$\alpha(\omega) = \sqrt{2}\omega \sqrt{\sqrt{\varepsilon_1^2(\omega) + \varepsilon_2^2(\omega)} - \varepsilon_1(\omega)}$$
(8)

where $\varepsilon_1(\omega)$ is the real part that is obtained using the Kramers-Kronig transformation, and $\varepsilon_2(\omega)$ is the imaginary part, which has been calculated by summing over a large amount of empty states. **Figure 7** shows the obtained absorption coefficient of pure and codoped GaN nanosheets. It can be seen that the C-Sn and C-Ge codoped can harvest a longer wavelength of visible light spectrum compared to the pure GaN nanosheet for visible light photocatalysis, which can be attributed to the reasonable band gap reduction.

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3.7. Dynamical Stability

Phonon calculations provide a criterion for the stability of one structure and indicate, through soft modes, structural instability (Ahmad Khandy et al., 2019; Xuan et al., 2019; Ahmad Khandy and Chai, 2020). We therefore calculated phonon dispersion curves for the codoped GaN nanosheet. The C-Ge codoping system was considered, as it has a negative formation energy and can be stably formed. **Figure 8** shows the phonon dispersion curves (PDCs) of C-Ge codoping system obtained using GGA scheme, which contains 24 phonon branches composed of Ga_3GeN_3C cell, including three acoustic modes and 21 optical modes. It is obvious from the figure that there is no virtual frequency in this doping system. It can therefore be concluded that the C-Ge codoped GaN system has dynamic stability.

4. CONCLUSIONS

The electronic properties of C monodoped, C-Ge, and C-Sn codoped GaN nanosheets have been investigated at the HSE06 level of theory. For C monodoping, the effective band gap values of the $GaN_{1-x}C_x(x = 0.125)$ nanosheet can be decreased. However, the C monodoping will introduce the unoccupied impurity states, which leads to the decrease of carrier concentration. The band gap values for the C-Ge and C-Sn codoping GaN nanosheets are 1.83 and 1.85 eV, respectively. Narrow band gap semiconductors can absorb and utilize visible light, and semiconductors can thus be used in solar energy conversion and photocatalysis. The C-Ge and C-Sn codoping GaN nanosheets are direct band gap semiconductors, which make them more advantageous over other materials. Anionic (C) and cationic (Ge or Sn) compensated cooping, greatly reducing the minimum formation energy of doping. However, the minimum defect formation energy of the C-Ge codoping GaN nanosheet is the smallest. The C-Ge and C-Sn codoped systems may be a potential candidate for photocatalytic water splitting to generate hydrogen because of their appropriate band gap and band edge positions. More importantly, the result of the optical absorption spectral analysis shows that the C-Ge and C-Sn codoped GaN nanosheets absorb a longer wavelength of the visible light spectrum as compared to the pure GaN nanosheet. Those results imply that the C-Ge codoping GaN nanosheet could be promising candidate as a visible light photocatalyst for water splitting.

DATA AVAILABILITY STATEMENT

All datasets generated for this study are included in the article/supplementary material.

AUTHOR CONTRIBUTIONS

CK: acquisition, analysis, and interpretation of data for the work. CT: drafting the work or revising it critically for important intellectual content. All authors contributed to the article and approved the submitted version.

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Conflict of Interest: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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