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Valter Moretti, Università degli Studi di Trento, Italy We describe the procedure of zeta regularization of an infinite product, in particular following the technique we introduced in previous works (13).

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The problem of regularizing infinite products is a long standing one in mathematical and theoretical physics, as well as in mathematics, of course. Several either formal or rigorous techniques have been introduced. One that obtained an undiscussed favor is the zeta regularization. The idea is simple to write down in a formal way (so as it as been used for a long times by physicists), and has got a rigorous mathematical formulation. This produced a huge number of interesting (and uninteresting) applications, with many useful (and non-useful) results. In this notes, we briefly recall the main features of the zeta regularization procedure, in particular following the approach introduced in our works.

Given a sequence  $S = \{a_n\}_{n=1}^{\infty}$  of non-vanishing complex numbers, arranged by increasing module, the formal product

$$\prod_{n=1}^{\infty} a_n,$$

appears in the series of formal equalities

$$\log \prod_{n=1}^{\infty} a_n = \sum_{n=1}^{\infty} \log a_n = -\left. \frac{d}{ds} \sum_{n=1}^{\infty} a_n^{-s} \right|_{s=0}$$

This suggests to call zeta regularized product of the  $a_n$ , the number

$$\mathrm{e}^{-\frac{\mathrm{d}}{\mathrm{d}s}\sum_{n=1}^{\infty}a_n^{-s}\Big|_{s=0}}.$$

In particular, if *S* is either the sequence of the eigenvalues of a matrix *A*, or the discrete spectrum of a differential operator *A*, we call *zeta regularized determinant* of *A* the number  $det_{\xi}A$  defined by

$$\log \det_{\zeta} A = - \left. \frac{\mathrm{d}}{\mathrm{d}s} \sum_{n=1}^{\infty} a_n^{-s} \right|_{s=0}.$$

This formal approach is based on the possibility of definining the function of the (complex) variable *s*:

$$\zeta(s,S) = \sum_{n=1}^{\infty} a_n^{-s},$$

called the *zeta function* associated to the sequence *S*, and on the analytic properties of it. In particular,  $\zeta(s, S)$  should be defined at s = 0 with its first derivative. As a matter of fact, it is not strictly necessary that  $\zeta(s, S)$  is defined at s = 0: it is sufficient that there exists a Lauren–Mac Laurin series for it.

Zeta functions have been deeply investigated in mathematics, mainly in analysis and number theory, and this study provided several different types of sequences for which the zeta regularization process works, and therefore possible applications to physics. In general, the series definition of the zeta function is analytic in a region of the complex plane that does not contain the origin, so analytic continuation is necessary. The problem is thus the one of finding out a suitable representation of the zeta function, that permits to describe its analytic continuation. The natural example is the Riemann zeta function

$$\zeta_{\mathrm{R}}(s) = \sum_{n=1}^{\infty} n^{-s},$$

where the sequence is  $S = \{n\}_{n=1}^{\infty}$ . As well known, we have the Hermite formula

$$\zeta_{\rm R}(s) = \frac{1}{2} + \frac{1}{s-1} + 2 \int_0^\infty \frac{(1+t^2)^{-\frac{s}{2}} \sin(s \arctan t)}{e^{2\pi t} - 1} dt,$$

that extends analytically in the whole complex plane up to the simple pole at s = 1 (with residue 1). The determination of the value of  $\zeta'_{\rm R}(0)$  is a bit harder, but follows as well differentiating the

formula of Hermite (and recalling some properties of the Euler Gamma function).

A main class of sequences that admit a zeta regularization is given by the spectrum of a class of linear operators: more precisely, given a non negative self adjoint differential operator H on a compact manifold, one can prove that the zeta function associated to the spectrum of H has an analytic extension at s = 0, and (hopefully) compute the derivative. Following Hawking (4), this provides a first important application in quantum physics, to deal with gaussian path integrals. The partition function for a finite temperature quantum field theory on a ultrastatic spacetime with compact spazial section is constructed as follows. Let M be a compact Riemannian manifold of dimension n, and consider the product  $N = S_r^1 \times M$ , where  $S_r^1$  is the circle of radius  $r = \frac{\beta}{2\pi}$ , and  $\beta = \frac{1}{T}$  is the inverse of the temperature. Let L be some non-negative self adjoint operator (typically the Laplacian) acting on some functions space (we shall deal with scalar fields) defined on M and  $H = -\partial_{\mu}^2 + L$ . The canonical partition function at temperature T of this model may be formally written as

$$Z = \det^{-\frac{1}{2}}(\ell^2 H),$$

where  $\ell$  is some renormalization constant. The zeta function of H is defined by (where Sp<sub>+</sub>H denotes the positive part of the spectrum of H)

$$\zeta(s; H) = \sum_{\lambda \in \operatorname{Sp}_{+}H} \lambda^{-s},$$

when  $\operatorname{Re}(s) > s_0$  (for some suitable  $s_0$ ), the regularized functional determinant of *H* is defined by

$$\det_{\zeta} H = e^{-\frac{d}{ds}\zeta(s;H)}\Big|_{s=0},$$

and the partition function is then

$$\log Z = \frac{1}{2}\zeta'(0; H) - \frac{1}{2}\zeta(0; H)\log \ell^2.$$

If the underlying manifold is not compact, there could be a continuous spectrum. In such a case, the idea of *relative zeta function* can be used (8). This works as follows. Let  $(H, H_0)$  be a pair of non negative self adjoint linear operators such that the difference of the resolvents (or the difference of the heat operators) is of trace class (some further properties are necessary, see for example 16 or 1 for details). Then, the relative zeta function is defined using the Mellin transform

$$\zeta(s; H, H_0) = \frac{1}{\Gamma(s)} \int_0^\infty t^{s-1} \operatorname{Tr} \left( e^{-tH} - e^{-tH_0} \right) \mathrm{d}t,$$

and the partition function reads

$$\log Z = \frac{1}{2}\zeta'(0; H, H_0) - \frac{1}{2}\zeta(0; H, H_0)\log\ell^2.$$

While several rigorous approaches have been introduced and developed in order to investigate the analytic properties of the zeta

function, comparatively few attempts to deal with the zeta determinant can be found in the literature. We briefly recall here our technique (see 12–14), whose main characterization is the fact of using a "new" spectral function, called *Gamma function*, instead of the usual heat kernel.

Let  $S = \{a_n\}_{n=1}^{\infty}$  be a sequence of non vanishing complex numbers with unique accumulation point at infinity, ordered by increasing modules. Assume the exponent of convergence of *S* is finite  $(e(S) = \limsup_{n\to\infty} \frac{\log n}{\log |\lambda_n|} < \infty)$ . We define the Gamma function associated to *S* by the convergent Weierstrass canonical product

$$\frac{1}{\Gamma(-\lambda,S)} = \prod_{n=1}^{\infty} \left(1 + \frac{-\lambda}{a_n}\right) e^{\sum_{j=1}^{p} \frac{(-1)^j}{j} \frac{(-\lambda)^j}{d_n}}.$$

where *p* is the *genus* of *S*, i.e. the least integer *p* such that the series  $\sum_{n=1}^{\infty} a_n^{-p-1}$  converges (absolutely).

We have the following formulas, relating the Gamma function to the other relevant spectral functions:

$$f(t,S) - 1 = -\frac{t}{2\pi i} \int_{\Lambda_{\theta,c}} e^{-\lambda t} \log \Gamma(-\lambda,S) d\lambda, \qquad (1)$$

$$\zeta(s,S) = \frac{s}{\Gamma(s)} \int_0^\infty t^{s-1} \frac{1}{2\pi i} \int_{\Lambda_{\theta,c}} \frac{e^{-\lambda t}}{-\lambda} \log \Gamma(-\lambda,S) d\lambda dt,$$
(2)

where  $\Lambda_{\theta,c} = \left\{ z \in \mathbb{C} | | \arg(z-c)| = \frac{\theta}{2} \right\}$ , oriented counter clockwise, and

$$f(t,S) = \sum_{n=1}^{\infty} \mathrm{e}^{-ta_n} + 1,$$

is the *heat function*, i.e. the trace of the heat operator.

In the classical approach the heat function is used to study the analytic properties of the zeta function using the Mellin transform

$$\zeta(s,S) = \frac{1}{\Gamma(s)} \int_0^\infty t^{s-1} (f(t,S) - 1) \mathrm{d}t$$

In fact, by this formula, explicit knowledge of the asymptotic expansion of the heat function for small t gives the position of the poles of the zeta function and the values of the residues.

In our approach, the main spectral function is the logarithmic Gamma function. It is clear, by the formula in Equation (1), that the asymptotic expansion of the heat function for small *t* follows from the one of the logarithmic Gamma function for large  $\lambda$ . Therefore, all information contained in the heat function are contained in the Gamma function (13).

However the main improvement obtained using the Gamma function is due to the fact that the coefficients in the asymptotic expansion of the logarithm of the Gamma function for large  $\lambda$  give the derivative of the zeta function [using the formula in Equation (2)], while this information is not available from the coefficients of the expansion of the heat function. In other word,

the expansion of the Gamma function is a global invariant, as the derivative of the zeta function, while the expansion of the heat function, as well known, is local, namely can be determined using local quantities constructed using the geometry of the manifold. This is the content of the next theorem (see 13 or the first section of 15 for details and proofs).

**Theorem 0.1.** If S is a regular sequence of spectral type of order  $\alpha_N \leq 0$ , then the associated zeta function is regular at s = 0, and near s = 0

$$\zeta(s, S) = -a_{0, 1} - a_{0, 0}s + O(s^2),$$

where the  $a_{0,1}$  and  $a_{0,0}$  are the coefficients of the logarithmic term and of the constant term in the asymptotic expansion of the logarithmic Gamma function for large  $\lambda$ , namely

$$\log \Gamma(-\lambda, S) = \sum_{j=0}^{N} a_{\alpha_{j},0} (-\lambda)^{\alpha_{j}} + \sum_{j=0}^{N} a_{\alpha_{j},1} (-\lambda)^{\alpha_{j}} \log(-\lambda) + O((-\lambda)^{\alpha_{N}}).$$

In the example of the Riemann zeta function considered above, the Gamma function is

$$\log \Gamma(-\lambda, S) = -\log \prod_{n=1}^{\infty} \left(1 + \frac{-\lambda}{n}\right) e^{-\frac{-\lambda}{n}}$$
$$= \gamma(-\lambda) + \ln \Gamma(-\lambda + 1)$$

that has the following expansion for large  $\lambda$ 

$$\log \Gamma(-\lambda, S) = (-\lambda) \ln(-\lambda) + (\gamma - 1)(-\lambda) + \frac{1}{2} \ln(-\lambda) + \frac{1}{2}$$
$$\ln 2\pi + \sum_{j=1}^{\infty} \frac{B_{2j}}{2j(2j-1)} (-\lambda)^{1-2j}.$$

Application of Theorem 0.1 gives

$$\zeta_R(0) = -\frac{1}{2}, \qquad \zeta'_R(0) = -\frac{1}{2}\log 2\pi.$$

We conclude with a remark on references. It is impossible to mention all the relevant references in this field. We decided to cite explicitly only our works (that contains the details on the approach that we briefly outlined here), and some particular references that are the undiscussed references usually associated to some particular subject. However, we want to add here some further references that are either interesting works on some particular problem or that contain themselves a good list of references on some aspect of the field. For rigorous treatment of regularized products we suggest (5–7, 17). For a good account of applications to physics see (2). For heat kernel methods and functional determinant see (3, 10). Of course these suggestions are far from being exhaustive.

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