



Transforming an IORT Linac Into a FLASH Research Machine: Procedure and Dosimetric Characterization

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Since Favaudon's paper of 2014, there has been an increasing interest in FLASH radiotherapy. The FLASH modality could represent a breakthrough in radiation oncology; nevertheless, it brings new scientific and technological challenges. Currently, one of the main limits the scientific community has to cope with is the lack of a common technological platform to experiment with. Considering this framework, the possibility of readapting existing linac platforms to produce a FLASH beam is particularly attractive and different attempts have been already made. The purpose of this article is to illustrate how it is possible to transform a dedicated Intra Operative Radio Therapy (IORT) mobile linac into a FLASH research machine. Compared to the modification required by a standard medical linac, such transformation is easier, does not affect the machine settings and can be rapidly performed by the final user. NOVAC 7 is an IORT linac which can reach a maximum dose-per-pulse up to 13 cGy/pulse (average dose rate 39 Gy/min); such dose rate can be significantly increased by modifying the collimation system.

Four different Source Surface Distance (SSD) can be obtained:

- Clinical reference configuration;
- Upper applicator only (SSD 50 cm);
- Monitor chambers housing only (SSD 7 cm);
- Dismounted monitor chambers (SSD 1.6 cm).

The fourth configuration allows reaching values of dose-per-pulse up to around 18 Gy/pulse and dose rates up to around 500 Gy/s, at a Pulse Repetition Frequency (PRF) of 30 Hz. The other three configurations can be obtained without using any tool and without changing NOVAC settings, until reaching a FLASH dose rate in the third configuration. For FLASH configurations, relative and absolute dosimetric characterization of the beam were performed using radiochromic films EBT3. NOVAC7 transformed in FLASH mode can be used both for dosimetric testing and characterization of detectors and for radiobiological studies on cells and organoids, offering a wide range of dose-per-pulse, from 3 cGy/pulse up to 18 Gy/pulse; dose rates correspondingly change from 3 cGy/s up to 540 Gy/s.

Keywords: FLASH, high dose-per-pulse, electron beam, Novac, IORT

INTRODUCTION

The FLASH effect in radiotherapy is a radiobiological effect characterized by a loss of radiobiological effectiveness (RBE) on healthy tissue and an unaltered therapeutic efficacy on tumor tissue; this effect is obtained by delivering the entire therapeutic dose in a very short time, <100 ms with a dose rate above 40 Gy/s [1, 2]. The experimental evidences of the FLASH effect were obtained *in vivo*, using 4–6 MeV energy electron beams, and the robustness of these results are validated by the fact that they were reproduced in various animal models (mice, rat, zebrafish, pig, cats), in various organs (lung, skin, gut, brain) and by various radiobiology researchers [1–8]. These evidences aroused considerable interest in the radiotherapy community, due to the possible clinical implications [9, 10].

However, for a possible clinical use of the FLASH effect, several issues must be addressed and understood. The radiobiological mechanism underlying the FLASH effect is still unknown: oxygen consumption has been proposed as a possible solution but other works highlighted how this mechanism, probably, cannot be considered the only one [11–15].

Additionally, the dependence of the FLASH effect on some parameters characterizing the radiation beam are not fully known: only the dependences on the average dose-rate and on the duration of the entire irradiation have been clearly observed so far. The role of dose-per-pulse, instantaneous dose-per-pulse (i.e., dose-per-pulse divided by pulse duration), pulse duration and frequency still remain to be entirely understood [1–3, 9].

Finally, the dosimetric problems related to the response of the on-line dosimeters to these dose-per-pulse values (saturation problems) are important and completely new to scientific community; this aspect causes difficulties in monitoring the stability of the beam output and in the accuracy of the dosimetric measurements. Even though many issues related to FLASH remain to be understood, such effect was observed only for average dose rates above 40 Gy/s. In the following, “FLASH dose rate” or “FLASH beam” will indicate beams with average dose rate above such threshold.

Because of this scenario, it would be important to increase the number of centers where a technology capable of delivering FLASH beam is available and where researchers can study the FLASH mechanisms. Up to now, all experimental data have been obtained either by using re-adapted standard medical linac [16] or using industrial machines [17–19].

This work shows how an accelerator dedicated to IORT (Novac7, SIT, Aprilia, Italy) can be set up to obtain different dose-per-pulse regimes and, consequently, dose-rates; in particular, two of these configurations allow to obtain FLASH beams. Such configurations were characterized by means of Gafchromic EBT3 films [20]. EBT3 were chosen because of their excellent spatial resolution, dose-rate (dose-per-pulse) independence [21–24] and energy independence for photon and electron beams above hundreds of keV [21, 25].

Radiochromic films [20] provide absolute measurements of absorbed dose to water after conversion of the film response by means of an accurate calibration procedure to be determined for

any specific radiochromic film dosimetry system, which consists in the combination of the film model and the densitometer, usually a flat-bed scanner, together.

Few studies investigated the dose-rate and dose-per-pulse dependence of radiochromic films [21, 22, 24]; however they all agree in reporting small or negligible dependence in their response with respect to both variables. At the dose-rate values of conventional clinical linacs, with a dose-per-pulse up to 1×10^{-3} Gy/pulse, Borca et al. [21] reported, for EBT3 radiochromic film, a dose-rate dependence in the range of 0.1–0.6 Gy/min within 1% for 6 and 15 MV photon beams. Jaccard et al. [22] reported, for conventional linac electron beams of 4, 8, and 12 MeV and EBT3, negligible variation in the range of 0.6–4.4 Gy/min.

Karsch et al. [23] reported, for a 20 MeV electron beam (5 ps pulse width) from the superconductive linear accelerator ELBE, a EBT radiochromic film dependence with respect to the dose-per-pulse of 2% up to about 2×10^{-2} Gy/pulse and within 5% up to about 7.5×10^{-2} Gy/pulse.

Jaccard et al. [22] also investigated the usability of Gafchromic EBT3 as reference dosimeters for an Oriatron eRT6 electron linac and concluded that EBT3 films are dose-per-pulse independent between about 4×10^{-3} Gy/pulse and 18 Gy/pulse.

Dosimetric and geometric properties of the beams obtained in the two FLASH NOVAC configurations have been evaluated in terms of dose to the build-up, dose at different depths and transversal dose profiles. These results can be useful to all NOVAC7 users, to design radiobiological experiments and/or study the response of the various dosimeters to FLASH dose-rate values. It is interesting to remind that, before the FLASH promises attracted the attention of scientific community, IORT linacs represented a challenge both from the dosimetric [26–31] and radiobiological point of view [32].

MATERIALS AND METHODS

Radiochromic Films Calibration

For calibration purposes, Gafchromic EBT3 film samples of $5 \times 5 \text{ cm}^2$ were irradiated using a Varian Clinac DHX-S (Varian Medical Systems, Palo Alto, CA), with a 6 MeV electron beam. The electron beam was previously calibrated in water, following the IAEA TRS 398 protocol guidelines [33] using a Roos reference chamber [34] at SSD = 100 cm and an applicator of $20 \times 20 \text{ cm}^2$. The films were irradiated in the same set-up except for the material, a Plastic Water[®] phantom, at the equivalent z_{ref} depth in Plastic Water[®], calculated as suggested by the IAEA TRS 398 protocol [33]. In order to obtain a calibration curve, the films were exposed, as described before, in a dose range from 2 to 20 Gy, with steps of 2 Gy. The post irradiation readings were made after 48 h. A black cardboard template was fitted into the scanner to ensure the reproducibility of the film positioning on the central location of the scan surface. The films were scanned after a 15 min warm-up time of the flatbed scanner and 3 empty scans to stabilize it. Films were acquired in transmission mode with all the image enhancement filters turned off, with a resolution of 127 dpi and at 48-bit RGB (Red, Green, Blue, 16 bits per channel). All of them were scanned

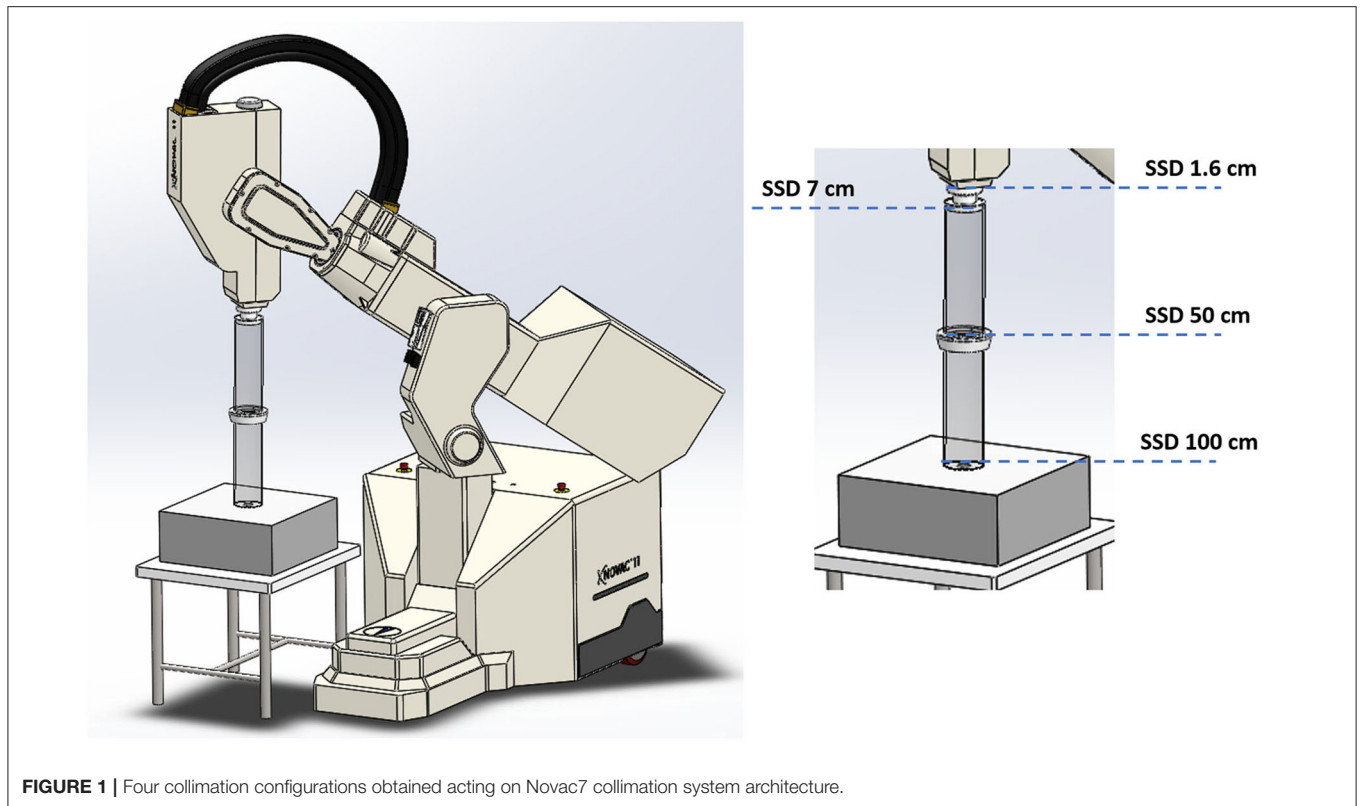


FIGURE 1 | Four collimation configurations obtained acting on Novac7 collimation system architecture.

in portrait orientation, i.e., the side of the $5 \times 5 \text{ cm}^2$ film sample corresponding to the long edge of the original film was positioned along the scanning direction. The images were saved in TIFF format.

A 2D Wiener filter was applied to both pre- and post-irradiation images, as suggested in [35]. The aforementioned protocol was used to obtain for each film sample the average net optical density \overline{netOD} , which is the difference between the irradiated and unirradiated optical density, over five $6 \times 6 \text{ mm}^2$ ROI (Region of Interest) positioned around the center of the radiation field. For each film sample the absorbed dose to water, D , measured with the ionization chamber, was plotted vs. the corresponding average \overline{netOD} for the Red channel. The calibration curve was determined by fitting the experimental data through the following equation:

$$D = a \cdot \overline{netOD} + b \cdot \overline{netOD}^n \quad (1)$$

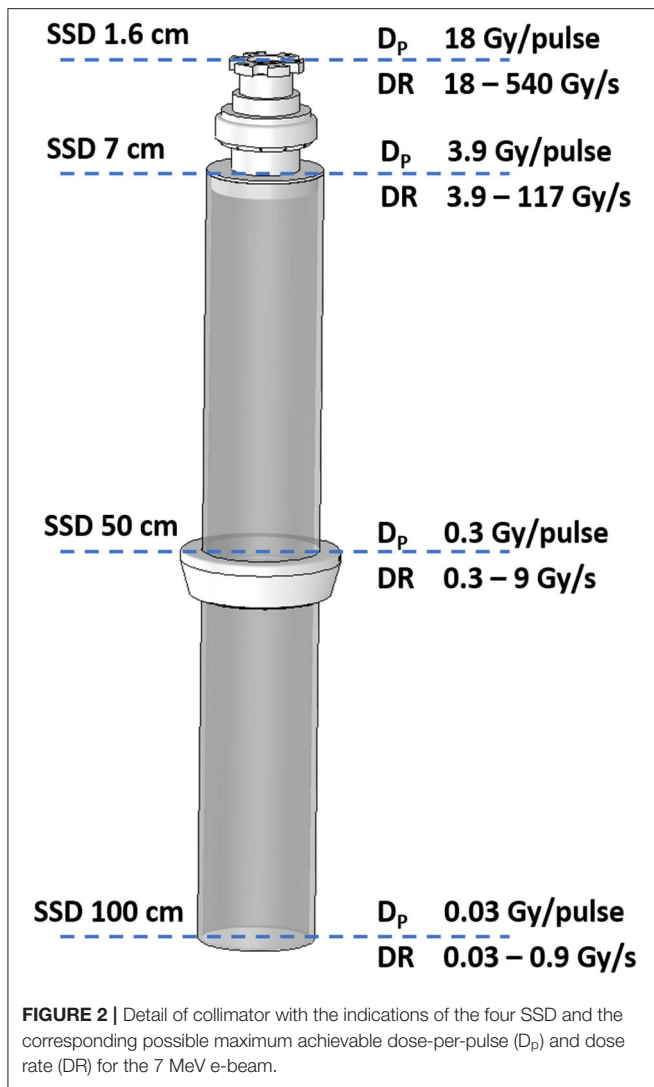
where a , b , and n are the fitting parameters. All the analysis was performed using home-made scripts in MATLAB R2018a environment (MathWorks, Natick, MA, USA). The dose curves fitting procedure was performed in two steps, following the method described in [35], which suggested to fix the parameter n after the first procedure to reduce the fitting uncertainty. Hence, once fixed n , a second fitting procedure was carried out, obtaining new values for the parameters a and b with their corresponding uncertainties.

NOVAC7

FLASH irradiations were performed using the IORT NOVAC7 (SIT, Aprilia, Italy) accelerator [36].

NOVAC7 provides four nominal electron energies (3, 5, 7, and 9 MeV) and the electron beam collimation system is purely passive; NOVAC7 does not use any scattering foil for beam broadening. The collimation system consists of polymethylmethacrylate (PMMA) cylindrical applicators that can be directly attached to the radiant head. The applicator is made of two parts: an upper part called applicator holder or upper applicator—directly mounted to the radiant head—and the terminal part called terminal applicator, which is connected to the upper applicator by means of a ring nut. The PMMA wall of the applicator is 5 mm thick, the internal diameter ranges from 4 to 10 cm and the very end of the terminal can be flat or beveled. The length of the applicators determines the SSD, which is 100 cm for the reference applicator with a diameter of 10 and 80 cm for the others. Thanks to this relatively simple architecture, it is possible to obtain several collimation configurations (Figure 1). Every configuration lead to a different SSD and, consequently, to a different resulting dose-per-pulse. The measurements were performed using the nominal energy of 7 MeV, which is the most used in the clinical practice and the closest to the electron energies for which the experimental FLASH effects were highlighted.

It is well worth underlying that, in general, the average dose-rate DR generated by a pulsed electron beam is directly



- (2) Upper applicator connected to the monitor chambers housing (SSD 50 cm);
- (3) Monitor chambers housing only (SSD 7 cm);
- (4) Dismounted monitor chambers (SSD 1.6 cm).

NOVAC 7 monitor chambers behavior remains unaffected in configurations 1, 2, and 3 (short term stability better than 0.5%); in the fourth configuration instead, NOVAC can be operated only setting the number of pulses to be delivered. In **Figure 2** such configurations, together with the possible maximum achievable dose-rates, are detailed. For all four configurations, the maximum dose-per-pulse value on the central axis of the beam in equivalent water phantom was measured, while for the last two configurations, being the only ones reaching FLASH values (called FLASH1 and FLASH2, corresponding to SSD of 7 and 1.6 cm, respectively), depth-dose measurements and dose profiles were also performed. The experimental setup used to characterize the beam in these two operating modes is shown in **Figure 3**. In order to characterize the beam in terms of dose-per-pulse and depth-dose curve, the radiochromic films were inserted perpendicularly to the electron-beam axis between Plastic Water® slabs at different depths for the first (1) and second (2) FLASH regimen, as reported in **Table 1**. Then, the accelerator head was put in contact with the first slab (see **Figure 3**, on the right).

Given the short distance between the first Plastic Water® layer and the beam exit window, it was possible to center the films manually with great accuracy. Moreover, the instantaneous darkening after the irradiation provided the possibility of an immediate check of the correct positioning.

The dimensions of films ($5 \times 5 \text{ cm}^2$) were suitable to include all the useful beam considering its broadening in depth. All the irradiations were performed using the nominal energy 7 MeV and delivering, for each point of measure, a total dose between 10 and 20 Gy. The total number of pulses delivered was changed according to the specific set-up, ranging from 400 pulses in clinical configuration down to just one pulse in FLASH 2 mode.

The radiochromic films irradiated with FLASH beams were read after 48 h through the same reading procedure adopted for their calibration.

To determine the amount of dose delivered in each image pixel of the films, the netOD_i for the i -th pixel was calculated [according to Equation (1)], then it was converted in dose using the calibration curve. Thus, the distribution of dose in the transverse plane was obtained at each depth for each FLASH condition.

A series of dose profiles were extracted from the above-mentioned dose maps for each measurement depth. The central region of each map was considered and the dose values of eight consecutive horizontal/vertical lines were averaged pixel by pixel to obtain the final dose profiles along the horizontal/vertical direction (thereinafter x and y).

To provide an estimation of the beam size, the Full Width at Half Maximum (FWHM) of the curves was used. All the profiles were fitted with a Gaussian function (Gaussian fits were performed by using Matlab R2018a fit tools), as shown in **Figures 6, 7**.

proportional to the dose-per-pulse D_p

$$DR = PRF \cdot D_p \tag{2}$$

where PRF is the pulse repetition frequency. The instantaneous dose rate IDR (or dose rate within pulse) is obtained dividing the dose-per-pulse D_p by the pulse length Δt . For the NOVAC Δt is about $2.5 \mu\text{s}$, its IDR can be easily calculated as

$$IDR \cong 4 \cdot 10^5 D_p \tag{3}$$

Due to the relatively low PRF (5 Hz in Clinical mode, up to 30 Hz in Service mode), dose rate is not extremely high; nevertheless, the dose-per-pulse can reach very high values.

Measurements

The different setups were obtained as follows:

- (1) Clinical reference configuration: (SSD 100 cm);



FIGURE 3 | Schematic view of the experimental set-up. On the left, the set-up with the accelerator head. On the right, the configuration during the irradiation (notice the proximity between beam window exit and the Plastic Water® slabs).

TABLE 1 | Film positions at different depth for the two FLASH configurations.

FLASH beam 1 depths [mm]	0	6	10	15	20	25	30	
FLASH beam 2 depths [mm]	0	6	11	13	15	17	22	32

The beam profiles were exploited to choose the ROI size for dose calculation as a compromise between two conflicting requirements: a larger ROI size allows to reduce the statistical error, while a smaller ROI increases the dose distribution uniformity. To preserve such advantages by minimizing the above-mentioned trade-off, ROI sizes variable with depth were adopted. In particular, ROI size was chosen in order to guarantee dose values fluctuation lower than 2% inside. The average dose delivered inside these ROIs is used to calculate the corresponding dose-per-pulse value.

Depth-dose curve was obtained from the depth-dose distribution and the depth corresponding to the 50% of the maximum absorbed dose, R_{50} , was evaluated.

RESULTS

Dose-Response Curves for EBT3 Calibration

The dose values plotted as function of the netOD are shown in **Figure 4**, together with the fitting curve calculated with the parametrization given in Equation (1). The corresponding fitting parameters are $a = 14.07 \pm 0.04$ Gy, $b = 52 \pm 2$ Gy, and $n = 3.45$.

FLASH Beams Characterization

Figure 5 shows the relative dose distributions measured in the two operating modes: FLASH beam 1 on the left and FLASH

beam 2 on the right. The dose distribution measured on the film is not uniform in the radial direction.

The profiles measured at each depth are presented in **Figure 6** (FLASH beam 1) and in **Figure 7** (FLASH beam 2). As confirmed by R^2 values reported in the figures, the profiles are well-approximated by a Gaussian curve except for the FLASH beam 2 profiles close to phantom surface ($R^2 < 0.99$).

The beam profiles along the x and y axes are reported in **Figures 6, 7** as function of the depths. The beam becomes significantly narrower when removing the monitor chamber, providing the smallest FWHM (**Figure 7**). Furthermore, the beam size increases with increasing depth in both irradiation modes, although the dependency from the measurement depth is not the same for the two cases (**Figure 8**).

The different width of the beam in the two configurations is due to both the different distance and the presence/absence of the two monitor chambers. A NOVAC7 monitor chamber basically consists of two aluminum electrodes, each 0.02 mm thick behaving as a thin scattering element [28].

The depth-dose distributions measured in the Plastic Water® phantom are shown in **Figure 9** for the two FLASH regimens. The maximum dose-per-pulse values and their relative errors obtained from **Figure 9** are presented in **Table 2**. The highest dose-per-pulse value is reached without monitor chamber (FLASH beam 2). In **Table 2**, the R_{50} calculated from the dose deposition curves are presented.

The uncertainty associated to dose measurements is 3% at SSD 100 cm, where the dose is measured by means of ionization chamber [26–28, 31] and 5% for all other points, where EBT3 are used. EBT3 tends to underestimate dose deposition when beam energy is below few tenths of keV [25]; therefore a higher uncertainty affects surface dose measurements. However, due to the very limited range of such electrons, this effect can be

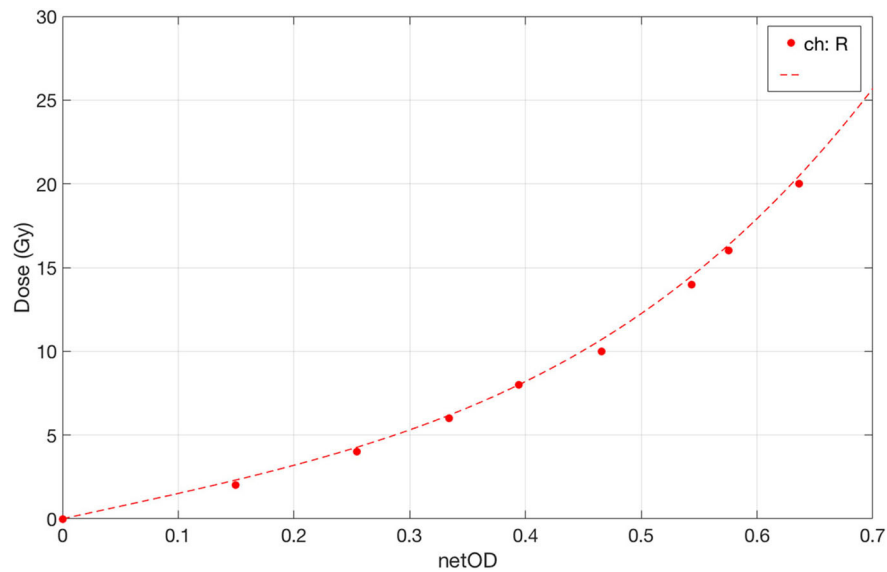


FIGURE 4 | Dose-response curve of EBT3 under conventional dose-rate irradiation with 6 MeV electrons in the range 0–20 Gy. The fit curve is plotted as dotted line.

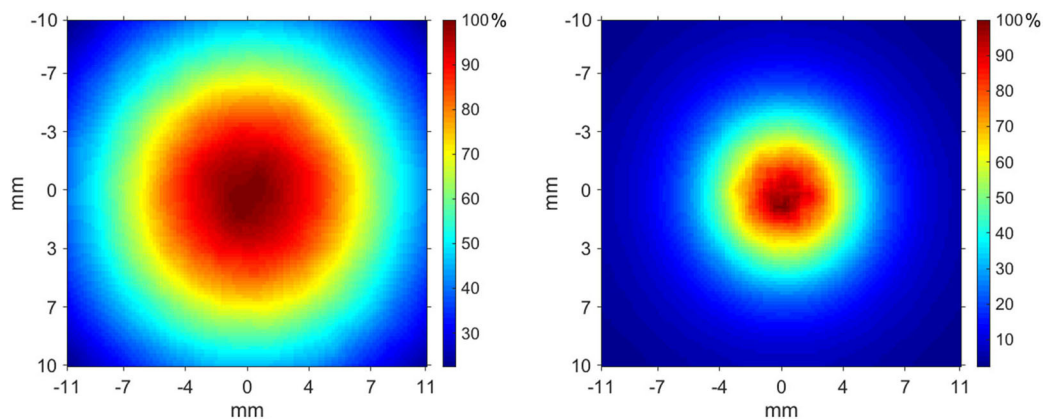


FIGURE 5 | Transverse dose distribution measured at 6 mm in Plastic Water® slabs—FLASH beam 1 (left), FLASH beam 2 (right).

considered negligible beyond 2–3 mm, and, consequently, it does not affect build-up measurements reported in **Table 2**.

DISCUSSION AND CONCLUSION

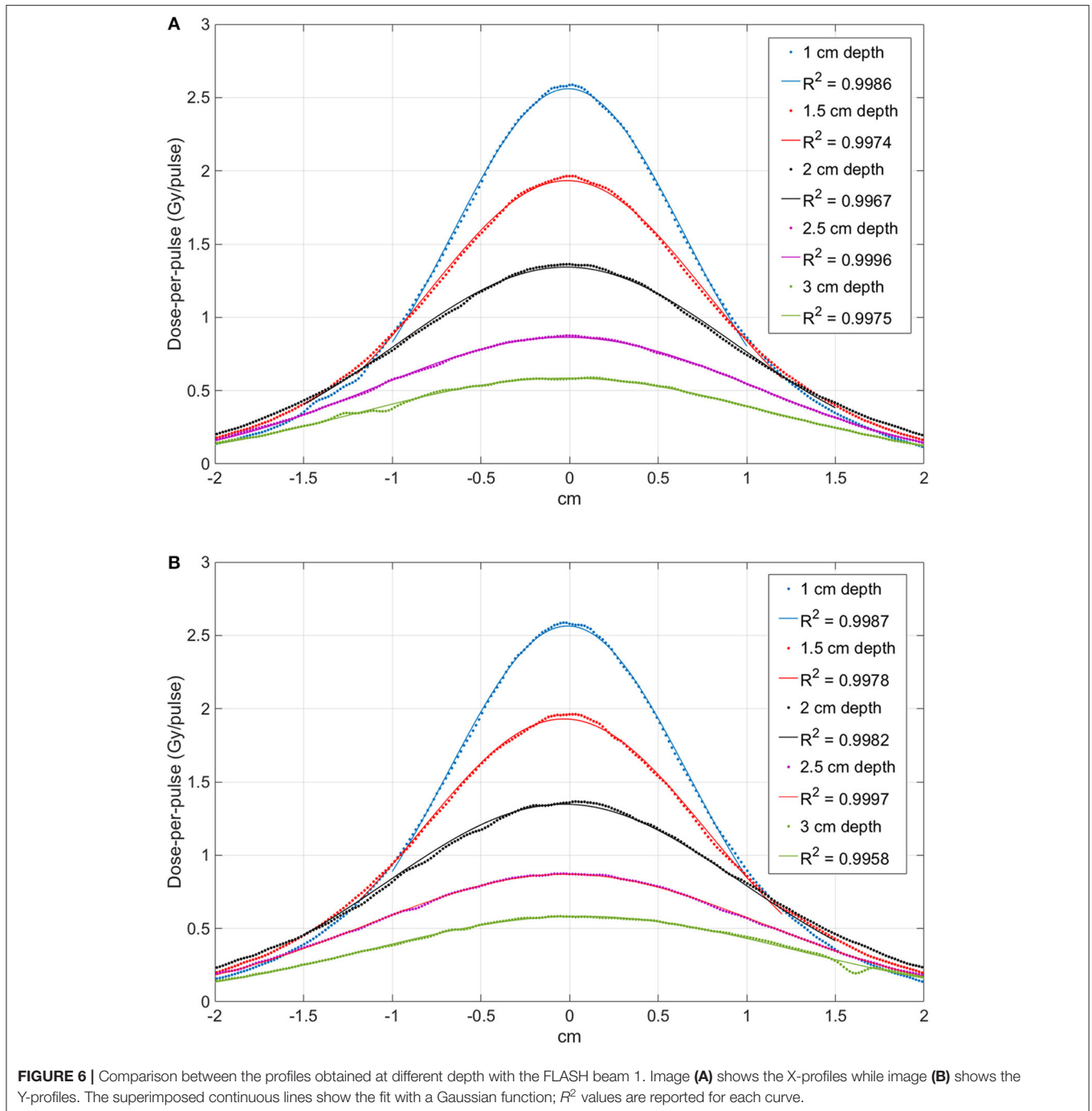
Currently, the two centers that published most of the FLASH results, Marie Curie Institute in Paris and University of Lausanne, adopted Kinetron and Oriatron, respectively, linacs originally designed for industrial use [1–3, 5, 7–9, 17–19].

Another possible solution is described in [16], where a procedure for modifying a standard clinical linac in order to get a FLASH beam is illustrated.

Several research groups are working to build a dedicated FLASH machine; not only electron based machines [37] are considered, but also linacs for X-ray [38] or proton accelerators [39]. The concept of PHASER [38] is particularly interesting,

but its feasibility remains extremely challenging both from the clinical and the technological point of view. Clinical difficulties related to “FLASH IMRT” are discussed in [40]; furthermore, the generation of a X-ray beam capable of reaching FLASH dose rate requires at least four times the electron current needed for linac working in electron mode [refer to NIST data [41] for Bremsstrahlung efficiency]. On the other hand, the effective implementation of FLASH with proton is feasible, even though it also poses several technological issues. However, the maximum dose rates achievable are significantly lower respect to electron based linacs, and many issues, in particular those related to real time beam monitoring, remain unsolved (Jolly et al. unpublished).

In this context, the possibility of expanding the number of researchers who can experiment with a FLASH beam may represent a crucial element for speeding up and validating the understanding of all phenomena involved.



This work described a procedure for transforming NOVAC, an IORT linac, into a FLASH machine: two out of four configurations identified reach the FLASH region (dose rate >40 Gy/s).

The geometric and dosimetric characterization of the beams was obtained through the use of Gafchromic EBT3 radiochromic films; EBT3 films were chosen because of their excellent spatial resolution, energy independence above hundreds of keV [25] and dose-per-pulse independence [21, 22, 24]. Good levels of

accuracy in measuring absolute dose could also be reached provided that a rigorous protocol is established [35].

The difference between the depth deposition curves in the different configurations, as shown in **Figure 9**, can be explained by the different electron spectra. The beam exiting the accelerating waveguide has a small but significant low energy tail [28]; such spectral component is entirely absorbed and filtered along the beam optic. In fact, the low energy components have a high spatial divergence and are either absorbed or scattered

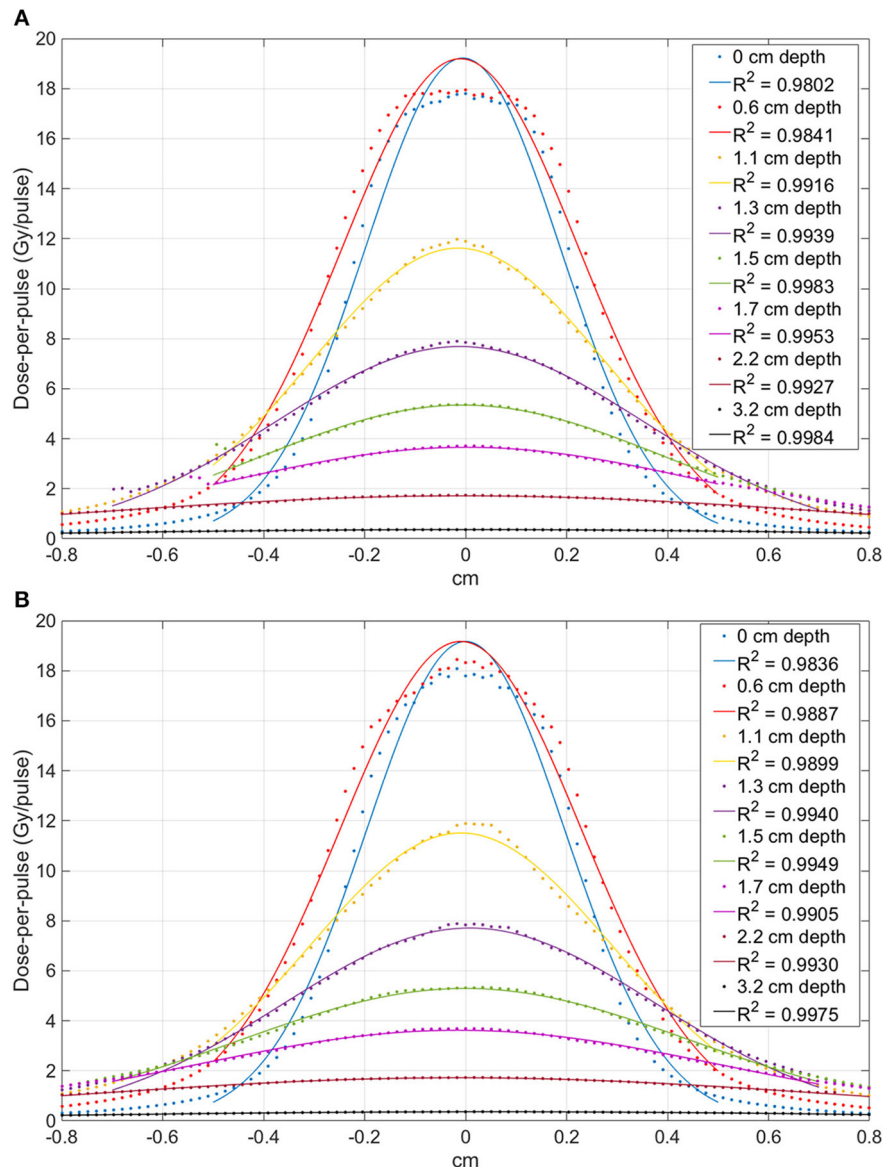


FIGURE 7 | Comparison between the profiles obtained at different depth with the FLASH beam 2. Image **(A)** shows the X-profiles while image **(B)** shows the Y-profiles. The superimposed continuous lines show the fit with a Gaussian function; R^2 values are reported for each curve.

away. Nevertheless, when measuring the beam at SSD 1.6 cm such component is still present. At SSD 7 cm (after the two monitor chambers) a significant fraction of such low energy electrons has been already absorbed or scattered away and the functional shape of the curve changes accordingly, with an increase of the parameter R_{50} . The advantage of this approach respect to the methods discussed by Lempart et al. [16] consists in its reproducibility and simplicity (no tool is needed and the modification is entirely and easily reversible). Fields sizes achievable are smaller (0.5 vs. 4 cm FWHM) but dose-per-pulse is higher (18 vs. 5 Gy/pulse).

The solution of transforming NOVAC7 IORT linac into a FLASH research machine is straightforward and gives to

its users the possibility of investigating mainly the detectors response to the new challenging dose-per-pulse region; any detector with transverse dimensions compatible with the beams produced can be tested (for example, all the dosimetry diodes and small plane parallel chamber such as PTW Adv. Markus).

This aspect is extremely important because passive dosimeters like TLDs, alanine pellets, Fricke gels could be considered suitable with respect to dose-per-pulse (dose-rate) independence also at FLASH regimen, but they all lack of spatial resolution and they do not provide on-line dosimetric information, while most of the active read out dosimeters are apparently affected by significant saturation problems [37].

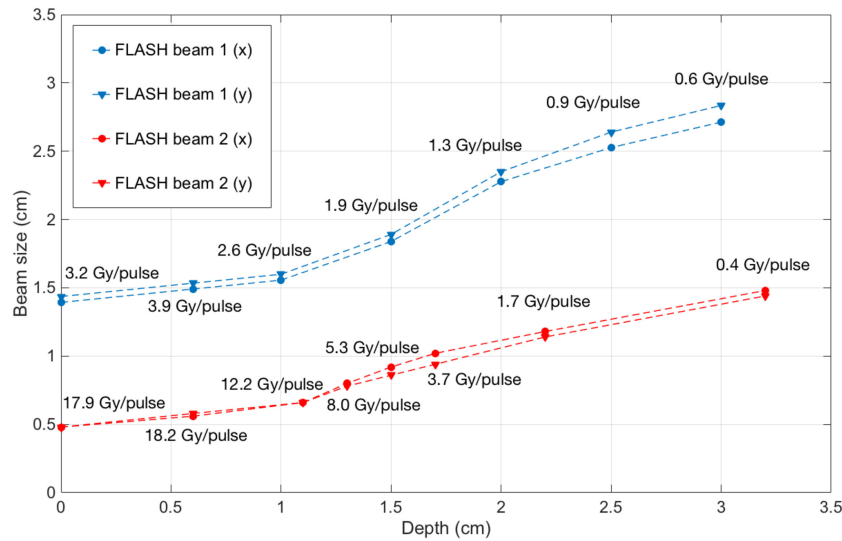


FIGURE 8 | Beam FWHM measured at different depths inside solid water with FLASH beam 1 (blue) and FLASH beam 2 (red), with the corresponding dose-per-pulse.

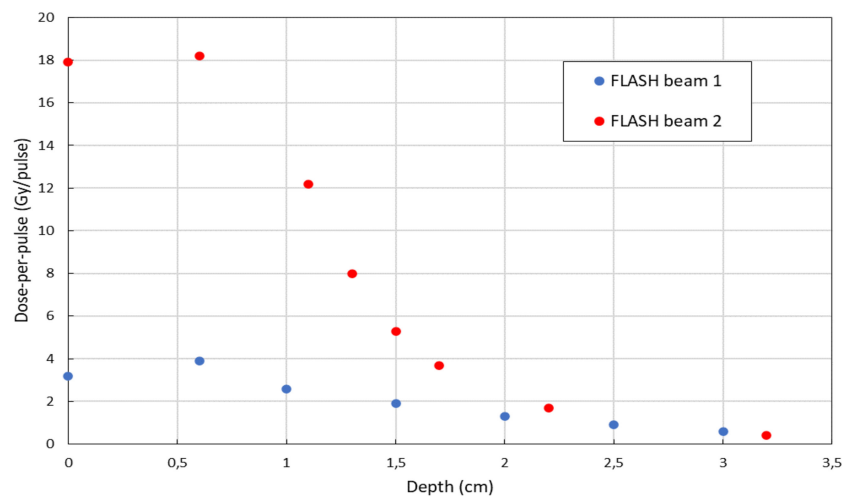


FIGURE 9 | Depth-dose distribution of electrons measured in a plastic phantom with EBT3 in FLASH beam 1 mode (blue) and FLASH beam 2 mode (red).

TABLE 2 | Maximum dose-per-pulse values and R_{50} obtained for the two irradiation modalities.

	Maximum dose-per-pulse [Gy/pulse]	R_{50} [cm]
CLINICAL BEAM—SSD 100 cm	0.030 ± 0.001	2.6 ± 0.1
SSD 50 cm	0.300 ± 0.015	2.6 ± 0.1
FLASH beam 1—SSD 7 cm	3.9 ± 0.2	1.7 ± 0.1
FLASH beam 2—SSD 1.6 cm	18.2 ± 0.9	1.2 ± 0.1

All the active detector commonly used in radiation therapy dosimetry have a signal collection time shorter than pulses repetition time (from 2 to 10 ms for a typical PRF of a conventional linac ranging from 400 to 100 Hz); consequently,

the saturation effect is influenced exclusively by the dose-per-pulse. Even in clinical configuration, NOVAC7 accelerator for IORT, with dose-per-pulse ranging from 3 to 13 cGy/pulse, represents a critical situation in the use of ionization chambers and several solutions have been already proposed to overcome this drawback [26–31].

Nevertheless, due to the small dimensions of the fields where FLASH dose rates are achievable, the only biological experiments that can be performed with NOVAC are cells plate or organoids [42] irradiation.

The possibility of transforming NOVAC into a FLASH device may lead to an increase in the number of researchers who can work with a FLASH beam, investigating and resolving the numerous dosimetric issues in order to set up rigorous radiobiological experiments and clinical trials.

DATA AVAILABILITY STATEMENT

The raw data supporting the conclusions of this article will be made available by the authors, without undue reservation.

AUTHOR CONTRIBUTIONS

GF and FD designed the study. FD, GF, and SL have written the paper. Measurements and data analysis have been performed by all the authors.

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