



Influence of Detector Size and **Positioning on Near-Infrared Measurements and Iso-pathlength Point of Turbid Materials**

Hamootal Duadi, Idit Feder and Dror Fixler*

Faculty of Engineering and the Institute of Nanotechnology and Advanced Materials, Bar Ilan University, Ramat Gan, Israel

Measuring physical phenomena in an experimental system is commonly limited by the detector. When dealing with spatially defined behaviors, the critical parameter is the detector size. In this work, we examine near-infrared (NIR) measurements of turbid media using different size detectors at different positions. We examine cylindrical and semi-infinite scattering samples and measure their intensity distribution. An apparent crossing point between samples with different scatterings was previously discovered and named the isopathlength point (IPL). Monte Carlo simulations show the expected changes due to an increase in detector size or similarly as the detector's location is distanced from the turbid element. First, the simulations show that the intensity profile changes, as well as the apparent IPL. Next, we show the average optical pathlength, and as a result, the differential pathlength factor, are mostly influenced by the detector size in the range close to the source. Experimental measurements using different size detectors at different locations validate the influence of these parameters on the intensity profiles and apparent IPL point. These findings must be considered when assessing optical parameters based on multiple scattering models. In cases such as NIR assessment of tissue oxygenation, size and location may cause false results for absorption or optical path.

Keywords: Monte Carlo simulations, near-infrared measurements, turbid media, scattering measurements, lighttissue interactions

INTRODUCTION

Visible and near-infrared (NIR) light are commonly used to non-invasively study diffusive media, such as biological tissues [1, 2], wood [3], fruit [4], and more. The interaction of NIR light with human tissue allows the assessment of hemodynamics and oxidative metabolism for diagnosis purposes [5]. Quantitatively finding the relative concentration of chromophores within the tissue requires an estimation of the optical properties, namely the absorption coefficient (μ_a) and reduced scattering coefficient (μ'_s). This is achieved by studying the diffusion profile compared to physical models derived from the radiative transfer theory under the diffusion approximation [6]. However, these solutions are inaccurate close to the light source [7, 8]. Alternatively, Monte Carlo (MC) simulations model photon migration and predict diffusion profiles as a function of geometry and components [9].

Even though it is possible to estimate the optical properties in ideal systems, non-idealities in real measurement systems can cause variations [10-12]. From a theoretical standpoint, the accuracy of the measurement system depends strongly, but not solely, on how well the solution of the theoretical

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> *Correspondence: Dror Fixler

> Dror.Fixler@biu.ac.il

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problem fits the experimental setup. Hence, curve-shape fitting methods mostly rely on calibration by phantoms with known optical properties to improve accuracy [5, 13]. However, while solving the practical problem, efforts still need to be made to understand the differences between theoretical and practical systems [14, 15].

During the last decade, many technological advancements for sources and detectors have been achieved [14, 16–18]. The signal to noise ratio is maximized by either increasing the detector size or positioning the detector close to the light source, since the maximal laser power is defined by the safety limits for skin. Further understanding of the influence of detector size and positioning when measuring diffused light could bridge the gap between theory and practice.

In this work, we test the influence of these parameters in NIR measurements. Specifically, we examine the change in the scattered light profile of cylindrical and semi-infinite scattering media. Previously we have demonstrated theoretically [19], numerically [20], and experimentally, first by phantoms [21] and then in human fingers [22], a unique optical phenomenon. This is the existence of an angle with minimal impact of the reduced scattering coefficient, named the iso-pathlength (IPL) point, dependent only on geometry. It also exists in semi-infinite geometry, more relevant in clinical applications of NIR spectroscopy on the head, arm, leg, etc. [23]. This phenomenon was demonstrated in a range of scatterings that fulfills the multiple scattering condition [24]. We have also shown that even though anisotropy also varies between different tissues, it does not change the IPL point [20]. Furthermore, absorption does not change the position of this point; it only attenuates the intensity [24]. It is therefore the optimal point for absorption measurements since it minimizes changes in optical pathlength [25]. Despite the fact that theoretical and numerical simulations are highly agreed, a consistent difference exists over experimental results [19]. This study explains this difference by examining experimental detector conditions and exploring the influence of the detector's size or equivalent position on three measurables: the angular intensity distribution, i.e. full scattering profile (FSP); the IPL point; and the optical pathlength (OPL).

METHODS

Monte Carlo Simulations

First, MC simulations are used to find the influence of detector size on the intensity profile and OPL. The medium is characterized by the reduced scattering coefficient and anisotropy factor g. All photons enter the medium at (Z = -D/2, X = 0), and change their direction after a distance of dr according to the scattering coefficient with a probability of 1–exp $(-\mu'_s dr)$. If the photon was scattered, its direction is varied by an angle of arccos(g) in either the polar or azimuthal angle. Once the photon reencounters the air-tissue interface, it exits the tissue. This MC simulation was previously validated by diffusion theory [19] and experimental results [26]. Our first MC simulation defines the air-tissue interface by a circle equation with a diameter D centered on the XZ surface, while Y is the longitudinal axis of the cylinder (**Figure 1A**). Photons accumulate in the FSP, which is the angular intensity distribution. In this study, the FSPs were collected with an angular accuracy $d\theta = 2^{\circ}$ and from a 1 mm thick slice in the Y axis. The number of steps it took for the photon to exit the cylinder was then used to calculate the OPL. MC simulations for varying cylinder tissue diameter D (between 6 and 14 mm) of different scatterings ($\mu'_s = 10$, 16 and 26 cm⁻¹) were carried out for different detector sizes. Taking into account these different diameters, we decided to set the maximum angular accuracy to 2° and look at detector sizes between one to 20 times larger than the length defined by this angular accuracy.

We then simulated a scenario where the collection of photons takes place some distance from the air-tissue interface since this also effectively increases the detection area. Since most tissues are a half-infinite medium, we also carried out simulations of this kind. Each simulation launched 2¹⁸ photons, and for simplicity, there was no absorption ($\mu_a = 0$).

Experimental Setup

To validate our findings, an experimental system for FSP measurements was established [21]. A continuous wave (CW) He-Ne laser with a wavelength of 632.8 nm, and 0.8 mW power (Thorlabs, Japan) was used as an input source. A rail was placed on a rotation stage and a post was mounted on a rail carrier to allow adjustments of the distance between the sample and the detector. Different detectors were placed on the post holder. Samples were aligned to the center of the rotation stage and measured in two ways. First, a silicon photodetector (PD, SM05PD2B, Thorlabs), with an active area of 13 mm², was placed close to the sample (not touching). Next, a multimode fiber, with a core diameter of 1.5 mm (M107L02, Thorlabs) was placed on the post at a similar distance.

The cylindrical samples were prepared from Intralipid (IL) in different concentrations as a scattering component and 1% agarose for solidification [27]. After heating the agarose, the IL was added and poured into glass cylindrical tubes with a diameter of 13 mm. The liquid was cooled under vacuum to avoid bubbles.

RESULTS

MC Study of the Influence on Intensity Profile

In the following results, the solid, dashed, and dotted lines of the intensity profiles correspond to reduced scattering coefficients of $\mu_s' = 10,16$ and 26 cm^{-1} . First, we examined the influence of detector size (defined in **Figure 1A**). For example, two selected diameters, 10 mm (**Figure 1B**) and 12 mm (**Figure 1C**), show a loss in accuracy. A small detector (black lines in **Figures 1B,C**) shows different IPL points for each diameter (171°, for D = 10 mm, and 173° for D = 12 mm). However, when increasing the detector size by a factor of two or more, there is a change in the FSPs and the apparent IPL point seems the same. For example, when increasing the detector by a factor of eight (red lines in **Figures 1B,C**) the apparent IPL points for both diameters are identical (157° for both diameters) and shift to a lower angle



(moving away from the light source). This emphasizes that this is not just an attenuation or linear shift, but rather a loss of spatial information and accuracy. To fully understand this, we looked at the change in apparent IPL point as the detector size increases for different diameters (D = 6,7,8,10,12 and 14 mm are red, yellow, green, cyan, blue and violet lines in **Figure 1D**). Each diameter presents a different slope of decrease so that as the detector size increases, the range of apparent IPL points decreases as well.

The influence of the location of the detector in respect to the air-tissue interface was also examined (**Figure 2A**). While moving away from the interface of an 8 mm diameter cylinder (**Figure 2B**) the IPL point moves from 168° (black lines) to 166° at a distance of 0.5 mm from the surface (blue lines), and 165° at a distance of 1 mm from the surface (red lines).

Since many clinical applications involve large tissue dimensions, MC simulations of a semi-infinite medium were examined (**Figure 2C**). As in the cylindrical geometry, the IPL point moves away from the source (**Figure 2D**). Given a step size of 30 µm, the IPL point measured for a 30 µm detector is r = 0.22 mm (black lines), compared to a 360 µm detector where the IPL point is at r = 0.28 mm (blue lines). The change in the apparent IPL point of the semi-infinite geometry (**Figure 2E**) has a similar linear behavior as the cylindrical geometry (**Figure 1D**). Note that a decrease in angle is an increase in distance from the

light source (according to the definition of the angle in Figure 1A).

MC Study of the Influence on Optical Pathlength

Next, we looked at the change in OPL (**Figure 3A**) and consequently differential pathlength factor (DPF) (**Figure 3B**), defined as DPF = OPL/($D \sin((\pi - \theta)/2)$). Note, that $D \sin(\pi - \theta)/2$ is the length of the cord between the angle θ and the light source point of entry. As before, we see evidence of a significant change in OPL, and as a direct outcome DPF at the shorter distances (r) from the light source (corresponding to high angles), while the influence reduces in farther distances. The high sensitivity of NIR measurements in close range to the light source has been previously discussed [14, 28–30]. A summary of the simulated effect of shift and accuracy of IPL point, as well as change in DPF, is presented (**Figure 3C**). For example, an increase of factor four in detector size - will result in a 3° ± 1°, change in IPL point and a change of factor 0.87 in OPL and DPF.

Experimental Results

The experimental system (Figure 4A) was used to measure the FSPs of the different scatterings (solid squares, dashed circles, and



FIGURE 2 | Influence of detector position and size. (A) In case the detector is distant from the surface of a cylindrical sample, the detector collects from a larger angular range (gray area), hence the detector acts as if it was larger than its physical size. (B) FSPs of 8 mm diameter cylinder adjacent to the air-tissue interface (black lines) in comparison to 0.5 mm (blue lines) and 1 mm (red lines) away from the interface. The original iso-pathlength point (IPL) point of 168° (black lines) moves to 166° at a distance of 0.5 mm (blue lines) and 165° at a distance of 1 mm (red lines). (C) MC simulation of the semi-infinite geometry demonstrate a similar behavior of the intensity as a function of distance (r), i.e., diffusion reflection profile. (D) While increasing the detector size, the IPL point moves from its original location (0.22 mm) further away from the input; for a 12-fold increase in detector size, the new IPL is at 0.28 mm (blue lines). (E) The apparent IPL point linearly increases with the detector size in the semi-infinite case.



dotted triangles in **Figure 4B** represent $\mu'_s = 10$, 16, and 26°cm⁻¹). An IPL point appeared at 115° (blue lines) when a PD was used to collect photons from the sample. When using the fiber (black lines) at the same distance from the sample as the PD, the apparent IPL point was much closer to the laser input, at 160°. The higher angle is closer to the expected IPL point in the ideal theoretical case (~174°), since the fiber collects from a smaller angular range. However, since the fiber is not touching the surface, it is effectively acting as a larger detector (**Figure 2A**). The fiber was then moved on the

rail so that it was further away from the sample surface (2 cm). The apparent IPL point, after moving the fiber (green lines in **Figure 4B**) was smaller (135° instead of 160°) than the previous position (black lines).

Let us consider the fiber from the experiment, which has a core diameter of 1.5 mm and a NA of 0.5. After moving the position of the fiber 2 cm away from its original position - the apparent IPL point moved 25°. Given the NA of the fiber (0.5 corresponds to an angle of 30°), a 2 cm distance translates to a factor of 4.85 of collected angular range compared to the core



diameter. This translates to a factor of 24 increase in the detection area. The experimental data shows a shift of $25^{\circ} \pm 5^{\circ}$, which is consistent with the values from the simulation (Table in **Figure 3C**).

DISCUSSION

In this work, we examined the influence of the detector size and positioning on measurements of cylindrical and semi-infinite scattering samples. We performed a MC study of the averaging effect increasing the detector size compared to the angular accuracy in different cylinder diameters. The increased detection range changed the FSPs and shifted the apparent IPL point further away from the light source. This also occurred when moving the detector away from the surface of the media, since this similarly increases the collection area. The increase in the detector indicates a loss in angular or special resolution. Furthermore, we simulated the influence of detector size on apparent OPL. Since the OPL is directly related to the DPF, both demonstrated a similar behavior. At closer distances to the light source, r, the OPL, and DPF decreased with the increase in detector size. This type of behavior, of the higher sensitivity of experimental measurements of OPL in the close range from the light source, is supported by previous reports [14, 28-30]. To further validate these findings, we performed CW measurements of the FSPs of cylindrical scattering phantoms. First, we measured the same phantoms by two detectors: by a PD and using a fiber. Since these detectors have different geometries, the apparent IPL point shifted from a theoretical one according to size. While the PD is larger and hence has improved signal to noise ratio and dynamic range, the fiber presents an IPL closer to the theoretical one. Nonetheless, using the larger detector, while considering the influence it has on the diffused light seems optimal in our opinion, giving the best of all worlds. Note that in the MC simulations the detector size was increased in one axis only, hence the experimental change in detector size in two dimensions had a stronger effect. As simulated, for a similar reason, while using the same detector and moving it further away from the sample surface, the apparent IPL point shifted further from the light source.

CONCLUSION

The findings of this work help the transition of NIR methods from research to clinical settings. Clinical applications of photoplethysmography, NIR diffuse optical imaging, and spectroscopy are increasingly adopted in the medical community as multimodal imaging has become more popular in recent years [10, 11, 14]. However, the diversity of responses of different medical instruments and the need for calibration complicates their use. This study highlights factors that should be considered in such instruments to allow quantitative, reliable, and sensitive physiological measurements.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article, further inquiries can be directed to the corresponding author.

AUTHOR CONTRIBUTIONS

This research conceptualization was formed by DF and HD. DF was responsible for the project administration, funding acquisition and research supervision methodology. The software, investigation, and formal analysis were done by HD The experiments were performed by IF and final validation was conducted by HD. The writing of the original draft was performed by HD, and DF was responsible for the review and editing for improving the paper.

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Conflict of Interest: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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