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Development and characterisation of caesium lead halide perovskite nanocomposite scintillators for X-ray detection

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We present work on the development of mixed-halide perovskite $(CsPbCl_xBr_{(1-x)})$ nanocrystal scintillators for X-ray detection applications. The effect of the varying the halide composition on the resulting peak emission and light yield is discussed, with the CsPbBr₃ materials displaying the greatest light yield. These perovskite nanocrystals were successfully loaded into PMMA, an inert plastic, at 2% mass weighting and the responses of these composites were compared to that of their colloidal dispersions. The composites were also characterised in terms of the radioluminescent light yield and decay response, alongside their X-ray sensitivity, in which the PMMA-CsPbBr₃ composites again outperformed the materials containing Cl^- anions.

KEYWORDS

X-ray scintillator, perovskite nanocrystals, caesium lead halide, nanocomposite scintillators, scintillation decay characterisation

1 Introduction

Lead halide perovskites are a class of material following the structure of the ABX₃ perovskite, where the B-cation site is occupied by Pb²⁺ and the X-anion site is a range of halides, e.g., Cl⁻, Br⁻ and I⁻ [1]. Pb²⁺ is an ideal component for these materials as radiation detectors, as its high-Z number ensures a large X-ray cross section and quantum efficiency. A benefit of the material is the tunable peak emission wavelength that the ratio of the anion site halide can bring. By varying the relative amounts of Cl, Br and I the peak emission from lead halide perovskites can be varied from 400–700 nm [2,3] by altering the semiconductor's band gap. By matching the spectrum emission of the perovskite to an associated photosensor, the system can be optimised to increase the overall efficiency of the detector. Additionally, the change in halide has effects on the scintillation decay time of the material as the band gap increases [4, 5], allowing further customisation for the system.

Previous work has shown that the A site cation can be widely varied, affecting the stability and electrical conductivity of the perovskite [6]. Choices include organic compounds, such as $CH_3NH_3^+$ (MA⁺) [7], $CH(NH_2)_2^+$ (FA⁺) [5] and $C(NH_2)_3^+$ (GA⁺); and inorganic cations such as Cs^+ [8], K⁺ [9] and Rb⁺ [10]. This work deals with CsPbX₃ perovskites. Figure 1A shows the change in mass attenuation coefficients for a range of common scintillator materials compared to both bulk CsPbBr₃ and a PMMA-CsPbBr₃ composite.



Perovskites can be used as both semiconductors and scintillators. For semiconductor purposes, perovskites are typically synthesised *via* melt-growth methods, such as the Bridgman method or solution growth methods. The Bridgman method can grow centimetre scale single crystals to be used for single crystal gamma spectroscopy [7]. Large polycrystalline materials can be formed *via* solution growth and then milled and pressed into wafers, for use as X-ray imaging systems [11] and dosimeters. Perovskites can also be used as scintillation radiation detectors, either in bulk form or as nanocomposites. This work covers characterisation of nanocomposite scintillators containing perovskite nanoparticles.

Perovskite nanoparticles have a light yield comparable or exceeding that of industry standard scintillators, benefiting both from the high X-ray sensitivity of the perovskite and also the increased light yield resulting from quantum confinement effects in the nanoparticles. Due to their structure, X-ray excitation leads to an excited triplet state, which has a high emission and fast decay time [12]. This likely originates from the strong spin-orbit coupling in the conduction band of the perovskite [13, 14], which allows for a fast single photon emission. However, perovskite nanoparticles do have limitations. The self-absorption of emitted light is an issue due to the low Stokes-shift of the system or from Mie scattering of the nanoparticles themselves [15]. This can be alleviated by adjusting the cations in the perovskite structure, using energy transfer processes such as FRET to shift the emitted light outside the absorption range of the material, and by adjusting the perovskite structure to increase the photon recycling efficiency [16].

The transparency of the composite must also be optimised, balancing a sufficiently high mass loading of nanoparticles to ensure a high quantum efficiency with the impact of the mass loading on the internal scattering arising from the nanoparticles. This can be resolved by ensuring an even distribution of nanoparticles in the composites during synthesis, and by limiting the material thickness to a few millimetres. Finally, perovskite nanoparticles are limited by their environmental stability. Lead halide perovskites are extremely hygroscopic and often have a short shelf life without the addition of stabilising agents such as ligands. By loading the nanoparticles into PMMA they are protected from the ambient environment and show a good level of stability.

Perovskite nanocrystals can be synthesised via a range of methods, including solid-state mechanosynthesis [17], precipitation [18] and hot injection [14]. These can produce materials in the size range of 5-100 nm, the size range in which quantum confinement effects can occur. CsPbX₃ nanoparticles have been recently shown to act as a scintillator, with a light yield reported to be 17,000 photons per MeV [19]. Composite perovskite scintillators are an active area of research, with recent work on FAPbX₃-PMMA composite scintillators reported as a viable direct X-ray imaging device [4]. In this work we report the production and analysis of 1 mm thick nanocomposite materials loaded with commercially available CsPbX3 nanoparticles, with a high X-ray stopping power and transmission. A range of bluegreen emitting nanoparticles were purchased, characterised and loaded into PMMA disks and their optical properties characterised, including their peak emission wavelength and time response to both laser and X-ray excitation.

2 Methods

2.1 Materials and composite fabrication

Dispersions of three different CsPbX₃ nanocrystals were purchased from the supplier Quantum Solutions. The three different compositions have peak emission wavelengths of 450, 480 and 510 nm, and the nanoparticles have a nominal diameter of ~10 nm. The dispersions consisted of the CsPbX₃, where X = Cl or Br, nanocrystals treated with oleic acid and oleylamine, dispersed in toluene (10 mg/mL) [20].

The PMMA composites were fabricated by mixing 1.5 g of PMMA (Aldrich) with a dispersion containing 30 mg of the dispersed nanocrystals in a vial with 3 mL of toluene. In preparation for the casting of the nanocrystals into the plastic, approximately 70% of the toluene was allowed to evaporate in a nitrogen-atmosphere glove box to increase the concentration of the dispersion to 17 mg/mL. The mixture was then heated at 56°C for 1 h on a hot plate, and stirred using a magnetic stirrer. The resulting mixture was then sonicated for

30 min to remove bubbles from the mixture and poured into 2 mm deep, 38 mm diameter moulds and left to cast over 2 days, covered to allow the toluene to recirculate and prevent the surface from distorting. The resulting composite disks were then removed from the moulds.

2.2 Characterisation methods

Transmission electron microscopy (TEM) imaging was performed with a Talos F200i TEM (Thermo Fisher Scientific). Dynamic light scattering (DLS) measurements were taken using a Zetasizer Nano S (Malvern Instruments), with the dispersion contained in square glass cuvettes. Photoluminescence (PL) spectra were taken with a 405 nm laser, recorded using a QE 6500 spectrometer (Ocean Insight), with a 420 nm long pass filter covering the collecting lens. For low temperature measurements, dropcast samples of the dispersions were mounted to a cryostat stage connected to a heating plate and liquid nitrogen supply, with the temperature set using a Temperature Measurement Control Unit TIC 304-MA (CryoVac). Absorption spectra were obtained using a UV-2401PC Spectrophotometer (Shimadzu). Time-resolved PL (TRPL) decay spectra were measured using a PicoQuant fluorescence lifetime spectrometer (PicoQuant), illuminated by a 405 nm laser pulsed with a DDS Function Generator (Aim-TTi). A 420 nm long pass filter covered the collecting lens. Scanning electron microscopy (SEM) imaging and energy dispersive X-ray spectroscopy (EDX) measurements were taken with a Thermo Scientific Apreo 2 SEM (Thermo Fisher Scientific).

Radioluminescence (RL) measurements were obtained with a QE 6500 spectrometer (Ocean Insight), connected to an integrating sphere (Labsphere) to ensure consistency of measurement geometry. The samples were irradiated by an X-ray tube (Hamamatsu) set to 80 kV and 100 µA. Time-resolved radioluminescence measurements were taken with samples coupled to a 8 mm diameter PMT (Hamamatsu H11901), housed within a copper Faraday cage. The samples were irradiated by a fourcapacitor Scandiflash X-ray tube, with X-ray bursts 4 µs apart. The tube voltage was set to 75 kVp, and the current had a 2 kA maximum. X-ray sensitivity data was taken in a dark box, with the samples mounted to a 90° mirror stand, with the light collected with a 2 inch PMT (ET Enterprises). The X-ray source was a Mini-X X-ray tube, with the voltage set to 40 kV and the dose varied by adjusting the current between 10-200 µA in 10 µA steps.

3 Characterisation

3.1 Perovskite nanocrystal dispersions

The dispersions were reported to emit at 450, 480 and 510 nm, and to have a size of ~10 nm by the manufacturer. The dispersions consisted of the CsPbX₃ nanocrystals, where X = Cl or Br, treated with oleic acid and oleylamine, dispersed in toluene. A photo of the three different CsPbX₃ nanocrystal dispersions scintillating under UV illumination can be seen in Figure 1B. TEM images of CsPbBr₃ nanocrystals are shown in Figure 1C. The TEM imaging shows a well-ordered cubic morphology of the nanoparticles, and a size distribution of approximately 10 nm. Once the materials had been received, the chemical compositions of the three nanoparticle dispersion were determined *via* EDX measurements taken with a SEM, utilising a circular backscatter (CBS) detector. The measured ratios of chlorine to bromine were found to be 1:1, 1:2 and 0:3 for the 450, 480 and 510 nm dispersions respectively. Therefore, it was determined that their structures were CsPbCl_{1.5}Br_{1.5}, CsPbClBr₂ and CsPbBr₃. The EDX spectra for samples of all three materials can be seen in Figure 2A, with the relevant peaks labelled, and the atomic percentages of the elements in Table 1. Si and Au peaks are labelled due to their presence in the measurement, with the samples placed on a silicon wafer and covered in a 3 nm gold layer.

Dynamic light scattering (DLS) measurements were performed to determine the size of the nanoparticles. Average sizes were 8 ± 1 nm for CsPbCl_{1.5}Br_{1.5}, 9 ± 1 nm for CsPbClBr₂ and 10 ± 1 nm for CsPbBr₃, consistent with the TEM imaging. The DLS spectra can be seen in Figure 2B, showing the relevant size distributions. As the bromine ratio increases, so too does the average particle size. For nanoparticles, as the size of the material decreases, the bandgap increases as the crystal lattice becomes smaller [21] and quantum confinement effects occur [22]. However, in these samples the halide present has a far greater impact on the bandgap, and therefore the emission peak of the samples is dominated by the amount of Br present, as can be seen in the photoluminescence spectra.

Normalised photoluminescence spectra of the nanocrystal dispersions can be seen in Figure 2C, illuminated by a 405 nm laser. The peak emission wavelengths for the dispersions were $CsPbCl_{1.5}Br_{1.5} = 450 \pm 2 \text{ nm}, CsPbClBr_2 = 475 \pm 2 \text{ nm}$ and $CsPbBr_3 = 512 \pm 2 nm$, agreeing with the reported emission peak wavelengths from the supplier. CsPbBr3 NC thin films have been reported to have peak photoluminescence emissions of 508 nm for a thin film of 2 nm nanoparticles [23] and 550 nm for a thin film of 9 nm radii nanoparticles [19]. The range of peak emission wavelengths measured for these materials shows the impact of the size of the nanoparticles. In our samples the peak intensities of the dispersions increased with the amount of Br present. Taking the CsPbCl_{1.5}Br_{1.5} peak intensity as 100%, the CsPbClBr₂ had an intensity of 175% and the CsPbBr₃ 200% as shown in Supplementary Figure S1A The decreased light yield for lead halide perovskites with chlorine present has also been seen in studies of organic FA lead halide nanoparticles [4].

Absorption spectra are also shown in Figure 2C, with absorption edges at CsPbCl_{1.5}Br_{1.5} = 446.0 \pm .5 nm, CsPbClBr₂ = 474.5 \pm .5 nm and CsPbBr₃ = 510.5 \pm .5 nm. It can be determined that the resulting Stokes-shift of the dispersions is slight, on the order of 1 nm, indicating that self-absorption may be occurring at the higher energies of the emission.

To investigate the effect of temperature on the emission spectra, photoluminescence spectra were taken in the range 293–77 K with drop cast samples of all three dispersions. The various spectra for the CsPbCl₁Br₂ drop cast spectra can be seen in Figure 2D. It can be seen that as the temperature decreases, the intensity of the peak increases and the peak wavelength shifts to a lower energy. Upon cooling the bandgap decreases in these materials, unlike for traditional semiconductors, causing the red shift in the peak emission. Due to the low exciton binding energy, there is also a significant increase in light intensity upon cooling. Figure 2E shows that for all three dispersions the FWHM of the spectra decreases with temperature, due to a reduction in thermal broadening [24]. Sub-structure in the photoluminescence spectra due to exciton emission is not observed in



(A) EDX spectra of samples of the three CsPbCl_x Br_{1-x} materials. (B) Size distributions of the dispersions measured using DLS. (C) Absorption and PL spectra of the dispersions. (D) PL emissions for the CsPbClBr₂ dispersion over the temperature range 293-77K. (E) The FWHM against inverse temperature for the dispersions over the temperature range 293-77K.

our data acquired at 77K, which would require temperatures of around 10K to be visible [25, 26].

3.2 Nanocomposite disks

To investigate the response of the nanoparticles in a solid matrix, and to obtain reliable radiation measurements that are difficult to perform with dispersions, plastic nanocomposite disks were produced. The plastic used was the optically-inactive and chemically inert PMMA, and the nanoparticles were loaded at 2% weight. This amount was chosen to balance the requirement for a high mass attenuation coefficient for photoelectric absorption with the need to minimise incoherent scattering, which would otherwise have a detrimental effect on the optical transmission of the nanocomposite materials. A photo of typical 1 mm thick, 38 mm diameter disks can be seen in Figure 3A. Figure 3B shows a scanning electron microscope image taken using the circular back scatter detector at 1,500× zoom of a PMMA disk loaded with CsPbBr₃ nanocrystals. The image is of a cross section of the composite disk. The nanocrystals appear as the bright spots in the image due to their high-Z compared to the plastic matrix and are evenly distributed throughout the sample. The cross

TABLE 1 Atomic percentages of Cs, Pb, Cl and Br for the three CsPbX_3 samples, measured with EDX.

Sample	Cs (%)	Pb (%)	CI (%)	Br (%)
CsPbCl _{1.5} Br _{1.5}	0.1	0.2	0.4	0.5
CsPbClBr ₂	0.3	0.4	0.5	1.1
CsPbBr ₃	0.4	0.4	0	1.7

section can be seen to contain many small voids. The loaded composites contain microbubbles resulting from the evaporation of the toluene introduced with the dispersion during the casting process. As discussed in the methods section, the plastic mixtures are treated *via* sonication and vacuum cycling to reduce the size of these bubbles, to diminish their impact on the transmission of the sample.

Normalised photoluminescence spectra of the composite disks are shown in Figure 3C. The PMMA composites have peak emission wavelengths of CsPbCl_{1.5}Br_{1.5} = 457 \pm 2 nm, CsPbClBr₂ = 482 \pm 2 nm and CsPbBr₃ = 517 \pm 2 nm. All three composite disks have a peak wavelength emission that is ~ 6 nm higher than their respective dispersions, indicating the some degree of nanoparticle agglomeration has occurred in the PMMA. Non-normalised spectra

are shown in Supplementary Figure S1B. Absorption spectra are also shown in Figure 3C, with absorption edges at CsPbCl_{1.5}Br_{1.5} = 439.0 \pm .5 nm, CsPbClBr₂ = 464.3 \pm .5 nm and CsPbBr₃ = 505.7 \pm .5 nm. These have also been shifted compared with the nanocrystal dispersions, with the absorption edge for all the composite disks at a lower wavelength. This results in larger Stokes shifts for all three composites: $CsPbCl_{1.5}Br_{1.5} = 18 \pm 2 \text{ nm}$, $CsPbClBr_2 = 17 \pm 2 \text{ nm}$ and $CsPbBr_3 = 11 \pm 2$ nm. This increase in the Stokes shifts of the nanoparticles in PMMA, compared to the dispersions, will tend to reduce the degree of self-absorption of the scintillation light in the plastic disks. The disks were illuminated by a 405 nm laser pulsed at 2.5 MHz to obtain photoluminescence decay spectra. The decay spectra are shown in Figure 3D, and were fitted with a biexponential fit, from which the average decay times were extracted. All the samples had a short decay time component of 9 ns. The long component of the decay of the PMMA composites were determined to be: CsPbCl_{1.5}Br_{1.5} = 26.7 \pm .2 ns; CsPbClBr₂ = 32.2 \pm .3 ns and CsPbBr₃ = $28.2 \pm .4$ ns, and the average decay times were: $CsPbCl_{1.5}Br_{1.5} = 11.4 \pm .1$ ns; $CsPbClBr_2 = 12.2 \pm .1$ ns and $CsPbBr_3 =$ $9.3 \pm .1$ ns. These values and their relative weightings can be seen in Table 2. Previous reports of decay times from CsPbBr₃ nanocrystal materials have not observed this weak slow time component, with values of the slow component of 2.9 ns for a thin film of 9 nm



FIGURE 3

(A) 3.8 cm diameter plastic disks loaded with $CsPbX_3$ nanocrystals. (B) A 1,500x zoom SEM image of a cross section of a PMMA disk loaded with a mixed halide dispersion. The white particles in the image are the perovskite nanocrystals spread across the thickness of the disk. (C) Absorption and emission spectra for the composite disks. (D) Time decay profiles of the disks under illumination by a 405 nm laser pulsed at 2.5 MHz.

PMMA composite	τ_1 (ns)	τ_1 weighting (%)	τ_2 (ns)	τ_2 weighting (%)	τ _{av} (ns)
CsPbCl _{1.5} Br _{1.5}	9.0 ± .1	85	26.7 ± .2	15	11.4 ± .1
$CsPbCl_1Br_2$	9.0 ± .1	87	32.2 ± .3	13	12.2 ± .1
CsPbBr ₃	9.0 ± .1	96	28.2 ± .4	4	9.3 ± .1

TABLE 2 Decay times and their relative weightings for the time decay curves of the CsPbX₃ composites in PMMA, illuminated by a pulsed 405 nm laser and fit with biexponential curves.



nanocrystals [19] and 7.8 ns for a drop cast film of 15 nm nanocrystals [8], while $CsPbCl_3$ have been reported to have slow component decay times of 5.4 ns for a drop cast film of 12 nm [8].

X-ray measurements were carried out to characterise the X-ray response of the nanocomposite scintillators. Radioluminescence spectra of the composite disks were taken under illumination with an 80 kV 100 μ A X-ray source, utilising an integrating sphere, the spectra are shown in Figure 4A. The PMMA composites have peak emissions of CsPbCl_{1.5}Br_{1.5} = 458 \pm 2 nm, CsPbClBr₂ = 482 \pm 2 nm and $CsPbBr_3 = 520 \pm 2$ nm. These values match their peak emissions from photoluminescence emission, showing that the radioluminescent optical emission due to X-rays is consistent with that observed for photoluminescence. CsPbBr3 nanocrystal thin films have been reported to have peak radioluminescence emissions at 532 nm under X-ray illumination at a dose rate of 278 $\mu \text{Gy}\,\text{s}^{-1}$ at 50 kV for a thin film of 9.6 nm nanocrystals [27]. The PMMA disk loaded with the CsPbCl₁₅Br₁₅ dispersion was considerably dimmer than the CsPbClBr₂ and CsPbBr₃ composites, as with the PL emissions from both the original dispersions and the composites. Taking the CsPbCl_{1.5}Br_{1.5} integrated intensity as 100%, the CsPbClBr₂ had an integrated intensity of 340% and the CsPbBr₃ 415%, as shown in Supplementary Figure S2A. This again shows the trend of increased light yield with increasing percentage of Br present in the perovskite nanocrystals.

To compare to the decay times from time-resolved photoluminescence emission, time resolved radioluminescence was performed using a pulsed X-ray set. This X-ray set has a pulse width of approximately 35 ns, so any decay mechanisms faster than this time would be unavailable to resolve, however this setup would allow any slower decay components to be investigated. The averaged decay curves for the CsPbBr₃ composite disk is shown in Figure 4B. The remaining spectra can be found in Supplementary Figure S3. In these spectra it can be seen that a short, ~30 ns day is observed consistent with the pulse width of the X-ray source, along with a longer component at a lower intensity. The long decay times were determined to be $CsPbCl_{1.5}Br_{1.5}$ = 191 ± 6 ns, $CsPbClBr_2$ = 206 ± 6 ns and $CsPbBr_3 = 123 \pm 3$ ns. These values and their relative weightings can be seen in Table 3. Previous work on drop cast films of CsPbBr3 nanocrystals has not observed this weaker, longlived component when exposed to gamma sources, with a decay time of 44.6 ns reported under exposure to a pulsed ¹³⁷Cs source [2].

The light response of the composite disks at different X-ray doses was measured with an X-ray tube, set at 40 kV and the current varied from 200 to 20 μ A. The response data can be seen in Figure 4C. Consistent with the radioluminescence measurement, the PMMA composite disk loaded with the CsPbCl_{1.5}Br_{1.5} nanocrystal dispersion was considerably dimmer than the CsPbClBr₂ and CsPbBr₃ composites. It can be seen that the composite disks have linear responses to X-ray radiation in this energy range, an important feature for X-ray detectors. The gradient of this response can be used to determine the sensitivity of the scintillators. Taking the CsPbCl_{1.5}Br_{1.5}

PMMA composite	τ_1 (ns)	τ_1 weighting (%)	τ ₂ (ns)	τ_2 weighting (%)	τ _{av} (ns)
CsPbCl _{1.5} Br _{1.5}	26.3 ± .1	97	191 ± 6	3	30.6 ± .2
CsPbClBr ₂	29.0 ± .1	97	206 ± 6	3	33.9 ± .2
CsPbBr ₃	23.4 ± .2	96	123 ± 2	4	27.7 ± .2

TABLE 3 Decay times and their relative weightings for the time decay curves of the CsPbX₃ composites in PMMA, illuminated by a pulsed 100 kV 2 kA X-ray and fit with a biexponential curve.

sensitivity as 100%, the CsPbClBr₂ had an sensitivity of 350% and the CsPbBr₃ 413%. These relative values are similar to that of the relative radioluminescence integrated areas, showing that relative light the yields of the composites are independent of the photosensor used to measure the scintillated light.

4 Conclusion

In summary, the optical properties of commercial $CsPbX_3$ nanoparticles were investigated, and their response to X-ray irradiation was studied. The nanocrystals were loaded into 1 mm thick PMMA disks to function as X-ray detectors, with the aim of producing a material that displayed the high radioluminescent light yield of perovskite nanocrystals whilst ensuring an even distribution of nanoparticles within the composites.

The nanocomposite materials present analogous peak wavelength emissions to their colloidal dispersions, and measurements showing the decay times of both laser and X-ray stimulated emission display consistent decay characteristics. The presence of chlorine in the perovskites can be seen to be linked to a decreased X-ray sensitivity and emission intensity, with the samples containing a higher ratio of bromine exhibiting superior performance.

Going forward, it is therefore recommended that the CsPbBr₃ nanoparticles are utilised when creating nanocomposites for X-ray detectors. These are ideal for use as beamline detectors and are being developed as direct X-ray imaging systems.

Data availability statement

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

Author contributions

CC, MW, PS, and SR contributed to the conceptualisation of the study. The composites were fabricated by JO and IB. The characterisation was performed by JO, with assistance from IB, JG, MM, and SR. Analysis was performed by JO, with assistance from IB, CC, JG, and PS. All authors commented on the manuscript and aided in the reviewing process.

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Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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Supplementary material

The Supplementary Material for this article can be found online at: https://www.frontiersin.org/articles/10.3389/fphy.2022.1046589/ full#supplementary-material

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