



Early-Time Non-Equilibrium Pitch Angle Diffusion of Electrons by Whistler-Mode Hiss in a Plasmaspheric Plume Associated with BARREL Precipitation

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In August 2015, the Balloon Array for Radiation belt Relativistic Electron Losses (BARREL) observed precipitation of energetic (<200 keV) electrons magnetically conjugate to a region of dense cold plasma as measured by the twin Van Allen Probes spacecraft. The two spacecraft passed through the high density region during multiple orbits, showing that the structure was spatial and relatively stable over many hours. The region, identified as a plasmaspheric plume, was filled with intense hiss-like plasma waves. We use a quasi-linear diffusion model to investigate plume whistler-mode hiss waves as the cause of precipitation observed by BARREL. The model input parameters are based on the observed wave, plasma and energetic particle properties obtained from Van Allen Probes. Diffusion coefficients are found to be largest in the same energy range as the precipitation observed by BARREL, indicating that the plume hiss waves were responsible for the precipitation. The event-driven pitch angle diffusion simulation is also used to investigate the evolution of the electron phase space density (PSD) for different energies and assumed initial pitch angle distributions. The results show a complex temporal evolution of the phase space density, with periods of both growth and loss. The earliest dynamics, within the ~5 first minutes, can be controlled by a growth of the PSD near the loss cone (by a factor up to ~2, depending on the conditions, pitch angle, and energy), favored by the absence of a gradient at the loss cone and by the gradients of the initial pitch angle distribution. Global loss by 1-3 orders of magnitude (depending on the energy) occurs within the first ~100 min of wave-particle interaction. The prevalence of plasmaspheric plumes and detached plasma regions suggests whistler-mode hiss waves could be an important driver of electron loss even at high L-value (L ~6), outside of the main plasmasphere.

Keywords: radiation belt, wave-particle interaction, electron precipitation, plasmaspheric plume, whistler-mode hiss/chorus, quasi-linear diffusion

1 INTRODUCTION

On August 10, 2015, BARREL balloon payload 3 A observed bremsstrahlung x-rays (~5–150 keV) attributed to precipitation of energetic electrons with energies < 200 keV. The balloon was located over Kiruna, Sweden, and magnetically conjugate to the Van Allen Probes (RBSP) spacecraft as they entered a region of high plasma density. The high density region was spatially confined and persisted for multiple orbits indicating that it is a plasmaspheric plume or detached plasma region. Strong whistler-mode hiss waves were observed by both spacecraft in the high density region.

Plasmaspheric hiss is a major driver of radiation belt loss inside the plasmasphere, contributing to the creation of the slot region between the inner and outer radiation belts (Lyons and Thorne, 1973). These low-frequency (~50 Hz–2 kHz) electromagnetic whistler-mode waves are broadband and incoherent. Their power is also proportional to the density, with more power in dense regions (Malaspina et al., 2016; Malaspina et al., 2018). Recent RBSP observations and simulations demonstrate the importance of hiss waves; these waves generate major loss up to $L = 5.5$ from pitch angle diffusion when the plasmasphere is widely extended during quiet times (Ripoll et al., 2016; Ripoll et al., 2017; Ripoll et al., 2019). More information about hiss waves and the dynamics of the radiation belts can be found in the review of Ripoll et al. (2020a).

Outside the plasmasphere, at higher L-values, energetic electron precipitation has been primarily attributed to whistler mode chorus waves or electromagnetic ion cyclotron (EMIC) waves [e.g., Millan and Thorne (2007); Thorne (2010)]. Recent work suggests that hiss may also be active at higher L-values in dense plasmaspheric plumes and detached plasma regions. A plasmaspheric plume (also called a “tail” in the literature) is a narrow region of dense cold plasma that extends out from the main plasmasphere at Earth. The EUV instrument on the IMAGE spacecraft provided the first global observations of plasmaspheric plumes, their formation and evolution (Sandel et al., 2001; Goldstein, 2004). Plume evolution is driven by changes in the dawn-dusk convection electric field due to changes in the solar wind (Grebowsky, 1970; Lemaire, 2000). A sudden increase in the convection electric field causes the plasmopause boundary to move inward leaving behind plasma that $E \times B$ drifts towards the dayside, forming a dayside plume that is wide in magnetic local time (MLT). With a subsequent decrease in the convection electric field, the plume wraps around Earth and the MLT extent narrows. Spacecraft passing through a plume observe a localized increase in cold plasma density over a narrow range in MLT. In the absence of global imaging, the twin Van Allen Probes spacecraft, moving in similar orbits - one following the other, provide an opportunity to distinguish between temporal and spatial variations, helping to identify spatial density structures like plumes.

Properties of hiss waves in plumes have been statistically characterized in Shi et al. (2019). The occurrence rate of hiss in plumes has a clear dependence on MLT and geomagnetic activity, peaking near the dusk sector during active times. These

waves are more powerful in plumes than inside the plasmasphere, particularly during active times, and have low wave normal angles compared to the wide wave distribution function in the plasmasphere. Plume whistler-mode waves were frequently observed by Van Allen Probes and may be an efficient loss mechanism for radiation belt seed electrons (Zhang et al., 2019). Zhang et al. (2018) calculated loss rates due to observed hiss emissions in a nightside plume and found that such emissions may result in fast (few hours) loss of 10–100 keV electrons at L in [4.5, 5.5] and a slower decay of higher energy particles. Li et al. (2019) examined a particular event and found that plume whistler mode waves are an effective pitch-angle scattering mechanism, particularly for electrons with energies from tens to hundreds of keV.

In this paper, we investigate plume whistler-mode hiss waves as the primary cause of the energetic electron precipitation observed by BARREL balloon 3 A on 10 August, 2015. The observations are described in **Section 2.1**. Data from Van Allen Probes are used to precisely determine the properties of observed waves, trapped electrons, and cold plasma density. An event-driven quasi-linear diffusion model is then used (**Section 2.2**) to investigate wave-particle interactions between the plume whistler-mode waves and the trapped electrons, seeking to identify the theoretical energy of the precipitating electrons. We further investigate the time evolution of the electron phase space density.

2 MATERIALS AND METHODS

2.1 Data and Observations

In this study, we use data from BARREL Campaign 3 and the two Van Allen Probes spacecraft to characterize observed electron precipitation and *in situ* plasma and wave properties respectively. BARREL measures electron precipitation by detecting the bremsstrahlung x-rays produced as electrons collide with the neutral atmosphere [see e.g., Sample et al. (2020) for review]. Each BARREL payload carries a single 3” diameter by 3” tall sodium-iodide scintillator nominally sensitive to x-rays from 20 keV to 10 MeV (Millan et al., 2013). Several BARREL data products are available as described in Woodger et al. (2015). In the present study, we use the fast spectra (FSPC) with 50 ms x-ray count rates, available in six energy channels covering 3–1300 keV. Note that the nominal FSPC energy range is 25–1500 keV but the Level 2 data product used here is calibrated to account for temperature-dependent gain shifts. GPS timing and positioning are transmitted using the Iridium satellite network along with the science and other housekeeping data. The third BARREL campaign was carried out at the Esrange facility near Kiruna, Sweden in August 2015. A total of seven balloons were launched over 16 days. Here, we use data from BARREL 3 A which was launched on 10 August at 13:50 UT and terminated on 10 August at 21:18 UT.

The twin Van Allen Probes (here also referred to as RBSP-A and RBSP-B) were launched on 30 August 2012 into nearly-equatorial orbits with 618 km perigee and 30,414 km apogee, providing two passes through the outer radiation belts every

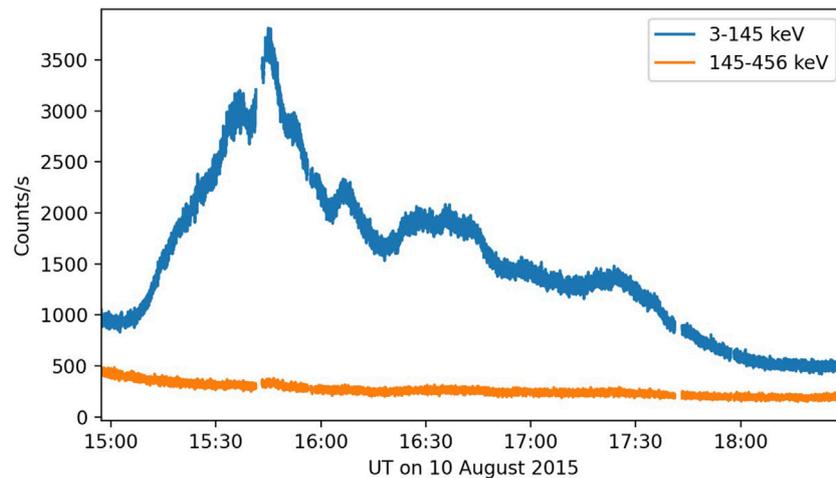


FIGURE 1 | BARREL Payload 3 A x-ray count rate versus time in two energy ranges: 3–145 keV (blue) and 145–456 keV (orange), observed on August 10, 2015. The 50 ms FSPC data have been summed to a 1 s cadence.

roughly 9 h. Each spacecraft carries a comprehensive suite of instrumentation to study particles and plasma waves throughout the inner magnetosphere (Mauk et al., 2013). In this study, we use data from the EMFISIS (Electric and Magnetic Field Instrument Suite and Integrated Science) (Kletzing et al., 2013), and MagEIS (Magnetic Electron Ion Spectrometer) instruments (Blake et al., 2013).

The EMFISIS waves instrument uses a tri-axial search coil magnetometer and electric field measurements from the EFW (Electric Field/Waves) instrument. A 6-channel waveform receiver (WFR) provides electric and magnetic field power from 10 Hz to 12 kHz. A single channel high frequency receiver (HFR) covers 10–500 kHz. For this study, we use Level 4 CDF files for cold plasma density, and the Level 2 WFR Spectral matrix product CDF files for wave observations, obtained from <https://emfisis.physics.uiowa.edu/data/index/>. The analysis used to determine the wave properties is described in more detail in **Section 2.2** below. For more information about EMFISIS, see Kletzing et al. (2013).

The MagEIS instrument consists of four magnetic spectrometers, one low-energy unit (20–240 keV), two medium-energy units (80–1200 keV), and a high-energy unit (800–4800 keV). For this study, we use Level 3 CDF files from: https://www.rbsep-ect.lanl.gov/data_pub/rbspa/mageis/level3/ which combine low, medium and high data into pitch-angle resolved electron flux. For more information about the MagEIS instrument design see Blake et al. (2013).

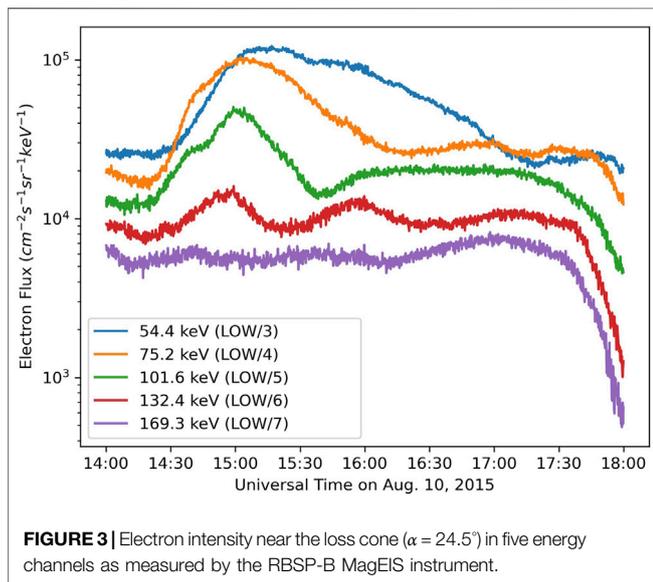
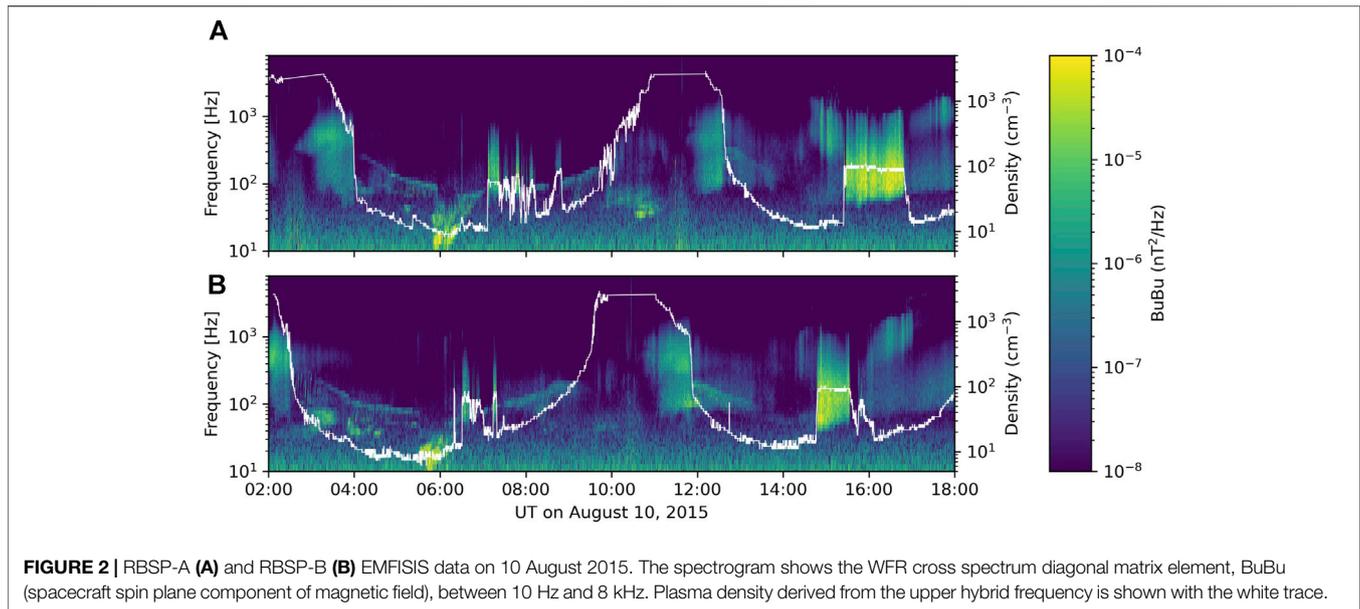
BARREL 3 A began to detect an increase in the 3–145 keV x-ray count rate (obtained by summing the FPSC 1a, 1b, and 1c energy channels) at 1505 UT (**Figure 1**, blue trace) just after the balloon reached a float altitude of 37 km. The energy of the precipitating electrons producing the x-rays was primarily below 145 keV as can be seen by the lack of count rate increase in the FPSC2 (145–456 keV) channel (**Figure 1**, orange trace). Precipitation was observed for about 3 h. Note that precipitation may have begun prior to 1505 UT while the

balloon was ascending and thus not sensitive to these low energy x-rays. At the onset of the observed precipitation, BARREL 3 A was located at L ~6, MLT ~1740 (T89 magnetic field model with Kp = 2), just east of the RBSP-B magnetic footprint. BARREL is nearly fixed geographically and thus moves in MLT as earth rotates.

Figure 2 shows EMFISIS data for just over two RBSP orbits on 10 August 2015. The two spacecraft were in similar orbits with RBSP-B (bottom panel) leading RBSP-A (top panel) by about 45 min. The white trace shows plasma density derived from the upper hybrid frequency (Kurth et al., 2015). At 1150 UT, RBSP-B exits the plasmasphere, as indicated by a steep decrease in plasma density. At 1445 UT, it encounters a localized region of increased density ($\sim 80 \text{ cm}^{-3}$). About 40 min later, at 1525 UT, RBSP-A encounters a similar localized increase in density. Both spacecraft encountered the high density region near apogee (L = 5.9) and in the same magnetic local time range (~ 16 –17 MLT), indicating that the density increase was a stable spatial structure which we identify as a plasmaspheric plume or detached plasma region. The high density region was also present on the previous orbit but with more spatial variation (**Figure 2**).

Both RBSP spacecraft detected wave power between roughly 50 Hz–1 kHz, identified as whistler-mode hiss waves, throughout the high density plume region (**Figure 2**, color spectrogram). There was substantial variability of the observed wave amplitude, with values ranging from ~ 50 –200 pT. For comparison, mean hiss amplitudes are between 10 and 30 pT depending on geomagnetic conditions (e.g., (Malaspina et al., 2018)). Large amplitude whistler waves similar to those shown here are also reported in the statistics of Shi et al. (2019). The wave analysis and properties are discussed in more detail in **Section 2.2** below.

We also examined electron data from the RBSP MagEIS instrument (**Figure 3**). A small substorm injection between 54–132 keV was observed starting at 1430 UT, just before RBSP-B entered the plume region. The injection shows some energy dispersion at both RBSP spacecraft. Data from the LANL



geosynchronous spacecraft also show evidence of a substorm injection, with the spacecraft closer to midnight (e.g., LANL-02 A) observing a dispersionless injection at around 1330 UT (not shown). This is consistent with the RBSP observations of a dispersed injection a bit later, since the RBSP spacecraft were in the afternoon sector. **Figure 4** shows example MagEIS pitch angle and energy distributions with superposed fits that are used in the modeling analysis described below.

2.2 Analysis and Model Description

To investigate whether the observed plume hiss waves are the cause of the precipitation observed by BARREL, we use an event-driven quasi-linear diffusion model. Wave properties (mean frequency, frequency cut-offs, wave normal angle, and wave

amplitude) over the time interval 1440–1700 UT on 10 August were generated from the EMFISIS data. We first compute the characteristic magnetic field amplitudes of whistler-mode waves using the trace of the magnetic-field power-spectral matrix calculated onboard RBSP-A from 3D measurements of the EMFISIS search coil antennas. The data were further processed into the Level 4 WNA data set (Kletzing et al., 2013) which serves as a basis for our analysis and additionally provides wave-normal angles. We select waves which propagate in the whistler mode and are right-hand polarized with ellipticity above 0.2 (Santolík and Gurnett, 2002). We restrict the analysis to a frequency range typical for plasmaspheric hiss from 50 Hz to 2 kHz (Santolík et al., 2001; Li W. et al., 2013; Thorne et al., 1973). The total power spectral density is therefore integrated over this frequency band. Mean frequency, mean frequency width, and wave normal angle are defined by equations (2)–(4) in (Ripoll et al., 2017) (not repeated here for brevity).

The average wave properties are shown in **Table 1** for RBSP-A and RBSP-B. Although we use only RBSP-A values for the quasi-linear computations described below, the wave properties measured by RBSP-B are also shown in **Table 1** to illustrate the stability and small variations of the wave system within the high density plume as measured by both spacecraft. We note that the waves have relatively small wave normal angles (last column in **Table 1**) in accordance with the statistical behavior found by Shi et al. (2019). The median frequency is also quite low, in agreement with observations of low frequency hiss (Malaspina et al., 2017), which favors interactions with higher energy electrons. Cold plasma density, derived from the upper hybrid frequency as measured by EMFISIS (Kurth et al., 2015), is reported in **Table 2**, along with the normalized wave parameters.

Diffusion coefficients are computed with the CEVA code originally developed by Réveillé et al. (2001). In this code, bounce averaged diffusion coefficients were computed following the method and equations of Lyons et al. (1972),

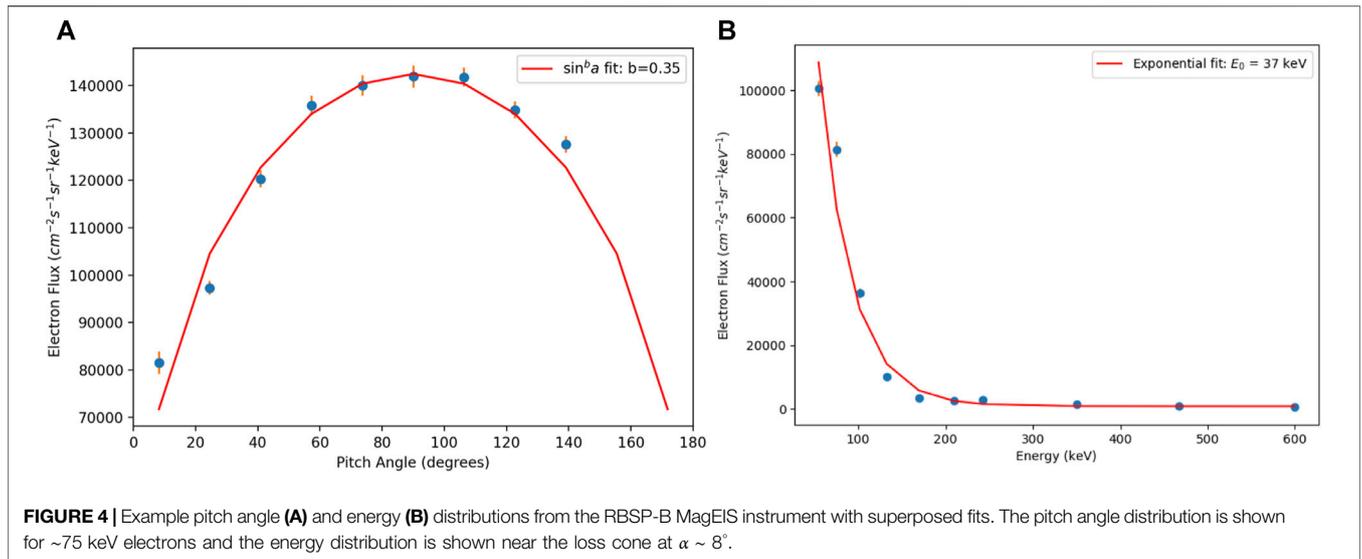


TABLE 1 | Average wave parameters obtained from both Van Allen Probe spacecraft during the event interval. θ is the wave normal angle.

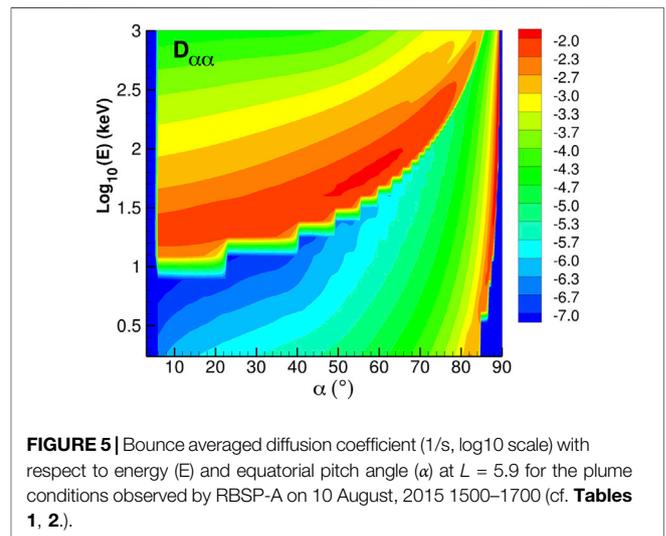
Spacecraft	Time	Amplitude	f_{median}	f_{min}	f_{max}	θ
RBSP	UT	(pT)	(Hz)	(Hz)	(Hz)	(degrees)
A	1500–1700	140	154	76	313	11
B	1440–1540	92	150	63	400	12

TABLE 2 | Average plasma and normalized wave parameters during the event interval. ω_{med} is the median wave angular frequency calculated using f_{median} from Table 1.

Spacecraft	n_e	ω_{pe}	ω_{ce}	ω_{med}/ω_{ce}	$\omega_{med}\omega_{ce}/\omega_{pe}^2$
RBSP	(cm ⁻³)	$\times 10^5 s^{-1}$	$\times 10^3 s^{-1}$	$\times 10^{-1}$	$\times 10^{-5}$
A	79.4	5.03	4.26	0.97	1.63
B	79.4	5.03	4.26	0.94	1.59

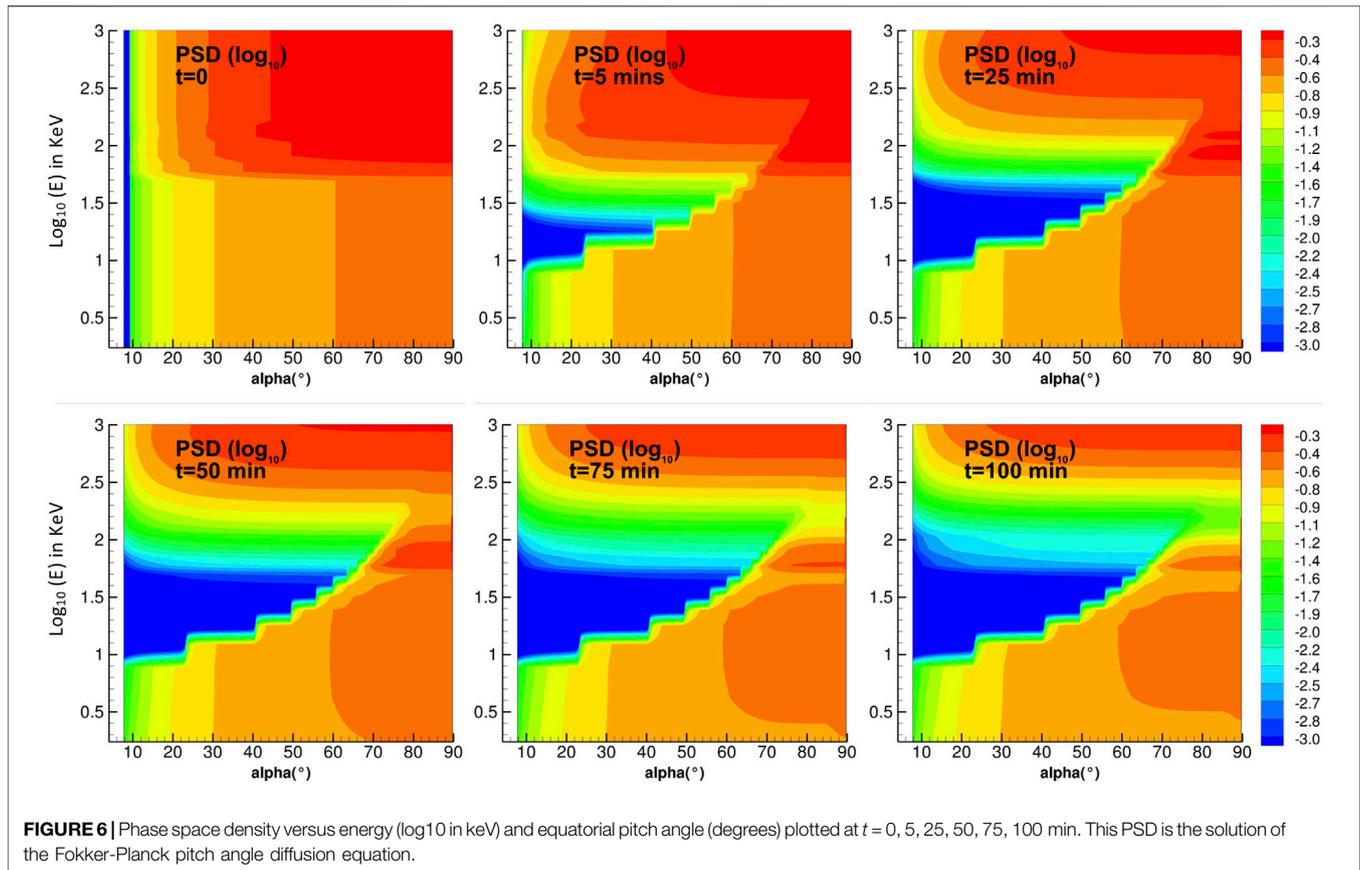
which account for a sum over all harmonics ($-n \dots, 0, \dots, n$), a wave normal integration, and bounce averaging between the mirror points. Equations 2 to 8 of Lyons et al. (1972) define the diffusion coefficients. The limit of low frequency ($\omega_{med}/\omega_{ce} < 1$) and high-density ($\omega_{med}\omega_{ce}/\omega_{pe}^2 \ll 1$) that are assumed in these computations are satisfied (cf. Table 2, last two columns). A more synthetic and modern expressions of the diffusion coefficients is available through Equations 8, 9 in Ripoll and Mourenas (2013) using the notations of Albert (2005). Verification by comparison with diffusion coefficients computed with the codes from the US AFRL and BAS (e.g., Albert (1994, 2008); Meredith et al. (2007)) have been performed in Ripoll and Mourenas (2013). Validation studies include Ripoll et al. (2016, 2017, 2019), Ripoll et al. (2020b), Ripoll et al. (2020c), and Loridan et al. (2019).

Diffusion coefficients were calculated using the average wave and plasma parameters from RBSP-A. They are based on the observed size of the high density region (~1 h MLT) and there is



no MLT-averaging (e.g., Spasojevic et al. (2015)) that accounts for the wave variations throughout the drift path (cf. more comments below). The inclusion of the available ambient measurements produces so-called event-driven diffusion coefficients which are today the most sophisticated method for reproducing radiation belt observations (see also Thorne et al. (2013); Tu et al. (2014); Ripoll et al. (2020b); Pierrard et al. (2021)). In doing so, we attempt to quantify the accuracy of this type of modeling. The calculated diffusion coefficients are discussed in Section 3 below.

We also investigate the evolution of the electron phase space density (PSD) using the calculated diffusion coefficients in a pitch angle diffusion based Fokker-Planck equation (e.g. Ripoll et al. (2017)). The initial electron distribution was taken to be dependent on pitch angle following the RBSP-B observations. A convolution with the energy spectrum can be performed after the computation is made (in the absence of energy



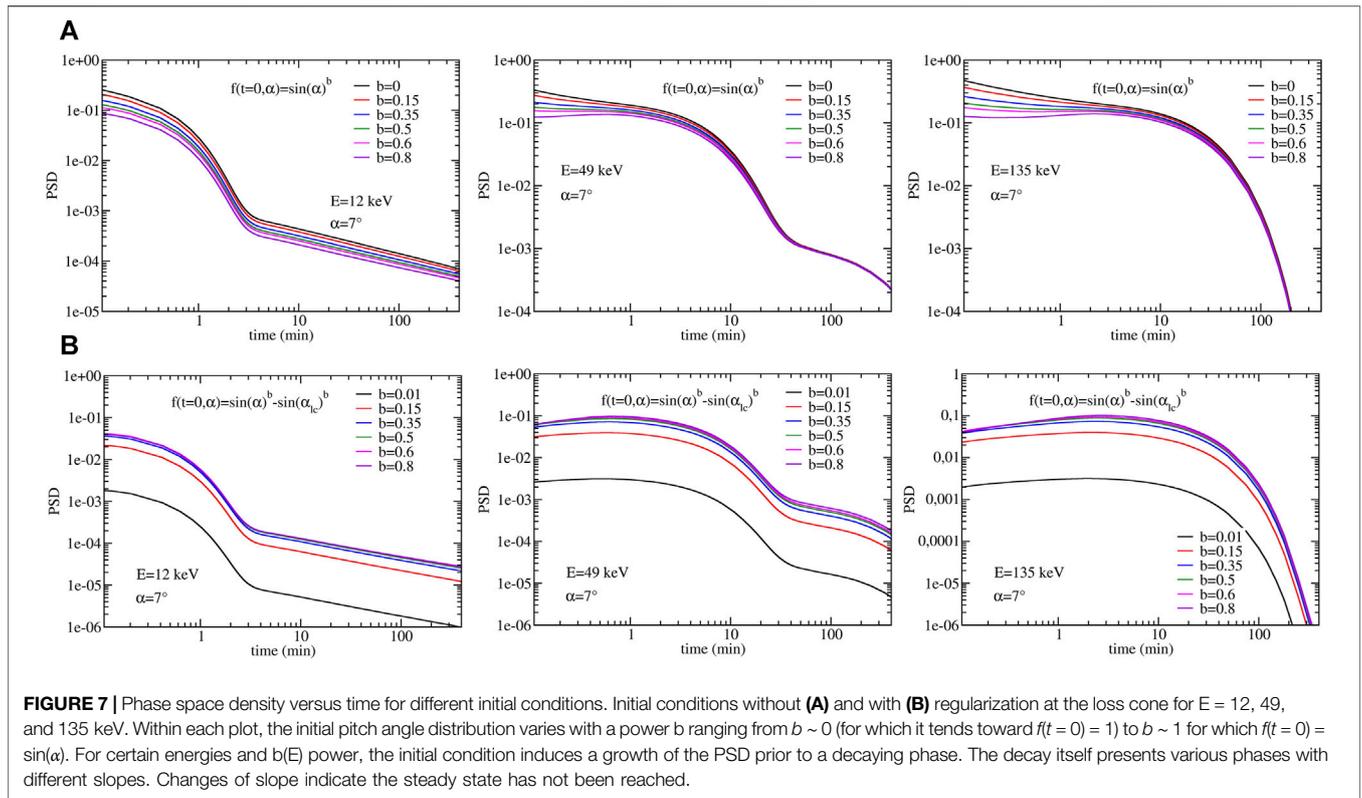
diffusion). For each MagEIS energy channel, the pitch angle distribution was fit with a $\sin^b(\alpha)$ function with a power $b(E)$ depending on energy (example shown for 75 keV in left panel of **Figure 4**). We find $b(E) = [0.15, 0.35, 0.52, 0.8, 0.6]$ for $E = [54, 75, 102, 132, 169]$ keV. The initial condition is regularized in a continuous manner within the loss cone using the expression $f(t = 0, E) = \sin^{b(E)}(\alpha) - \sin^{b(E)}(\alpha_{lc})$ (Li Z. et al. (2013)).

3 RESULTS

Figure 5 shows the bounce averaged diffusion coefficients calculated using the input parameters described above. At large L-shell (e.g., in the outer belt or in plasmaspheric plumes, $L \sim 6$), hiss waves will predominantly act on low energy electrons. Here, the diffusion coefficient for small pitch angle particles is largest between ~ 10 – 150 keV, consistent with the energy of x-rays observed by BARREL. As mentioned previously, the diffusion coefficient is not drift-averaged and is representative of the wave particle interactions only within the plume. The waves were observed only within the plume region, extending ~ 1 h of MLT, thus drift-bounce-averaged diffusion coefficients would be a factor $1/24$ lower in amplitude. Here, due to the intense wave amplitude and the absence of MLT-averaging, diffusion coefficients reach very high rates with $D_{\alpha\alpha} \sim 10^{-2} \text{ s}^{-1}$. Since BARREL was located in the northern hemisphere, it

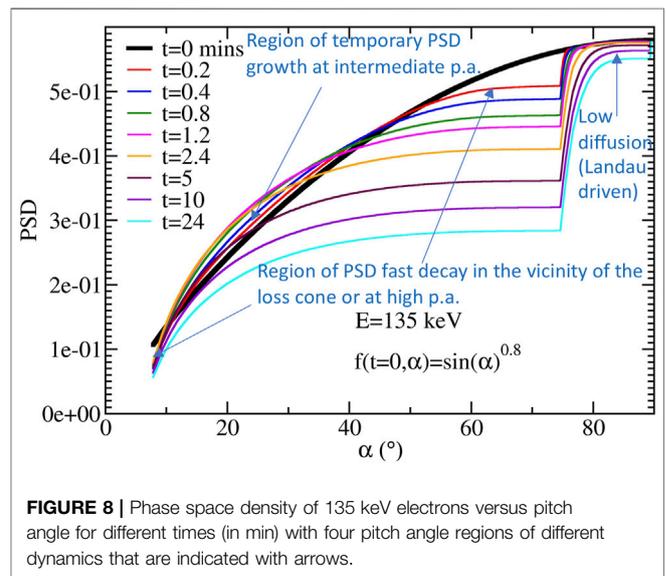
observed predominantly local precipitation. This local bounce loss comes from the immediate and intense (low pitch angle) precipitation within the plume (thus proportional to the local wave amplitude squared) and is not the result of wave particle interactions all along the drift path (that would be proportional to the MLT-averaged squared amplitude, a factor $1/24$ lower in this case).

Figure 6 shows the evolution of the phase space density for electrons as a function of energy, E , and equatorial pitch angle, α , from $t = 0$ to $t = 100$ min of interaction with hiss waves in the high-density region. The initial condition (top left panel) is $f(t = 0, E) = \sin^{b(E)}(\alpha) - \sin^{b(E)}(\alpha_{lc})$. Sudden loss of 10–30 keV electrons occurs almost immediately, within the first 5 min, in the vicinity of the loss cone (**Figure 6**, top middle). Loss of 100 keV electrons by one order of magnitude takes 25 min of interaction. After 100 min (**Figure 6**, bottom right), the phase space density has vanished by ~ 2 (resp. 1) orders of magnitude for all pitch angles below 70° for E in $[20, 70]$ keV (resp $[70, 200]$ keV). Note that electrons, particularly of higher energy, may interact with the plume multiple times as they gradient-curvature drift around Earth in order to reach a total of 100 min of interaction. 10 keV electrons with small pitch angle take ~ 40 min to drift across the plume, thus the rapid loss of low energy electrons within the first 5 min occurs upon their first encounter with the plume. The predicted range of energies of the main loss are also in agreement with the range of energy found by Shi et al. (2019). Some loss of

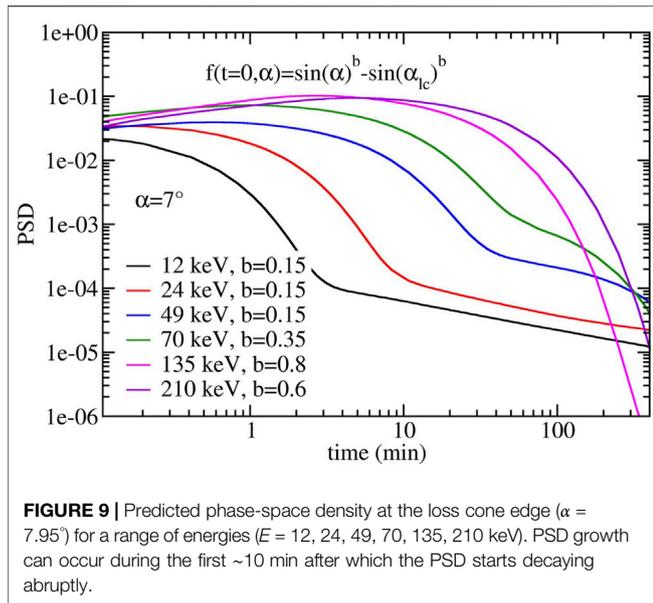


high energy electrons up to 1 MeV within the vicinity of the loss cone are also computed.

We also study the influence of the initial conditions on the loss and on the dynamics at early times, when pitch angle diffusion has not yet reached an equilibrium state. **Figure 7** shows the evolution of the PSD in the vicinity of the loss cone (at $\alpha = 7.95^\circ$, rounded to 7° in the figure’s label) during the first 100 min of interaction with hiss waves. Comparing the top and bottom rows of **Figure 7** illustrates the important role of the loss cone regularization. Comparing lines within each figure allows to see the influence of the $b(E)$ power index, with two limits bounding the variations. The low limit $b(E) \rightarrow 0$ (i.e. $b = 0.01$) is equivalent to the widely used and standard $f(t = 0, \alpha) = 1$ initial condition, either dropping abruptly from one for $\alpha \neq \alpha_{lc}$ to 0 for $\alpha = \alpha_{lc}$ if there is no regularization (top row) or going continuously from one to 0 if regularized (bottom row). The second limit of $b(E) \rightarrow 1$ is a PSD approaching $\sin(\alpha)$ and with significant gradient at intermediate pitch angle. **Figure 7** provides interesting new insight into the dynamics and how they are modeled. First, as expected, the larger the gradient at the loss cone, the more intense will be the loss in the vicinity of the loss cone and at the earliest times (generally below 1 min for our case, depending on energy). The regularization, which can be seen as a continuous process to fill out the loss cone, limits the gradient and thus the intensity of the earliest loss. In turn, large $b(E)$ power index leads to intermediate pitch angle diffusion that will make the overall loss-gain budget at fixed pitch angle to lean in favor of a temporary growth of the PSD level during the earliest times. This is better seen in the presence of the regularization (since loss



is even milder). On the contrary, without regularization the loss can dominate for certain low energy and $b(E)$, and temporary growth can no longer exist. Most past parameter studies in the literature have been done with $f(t = 0, \alpha) = 1$ and $f(t = 0, \alpha_{lc}) = 0$ (e.g., Shprits et al. (2008)) so that PSD growth could not be seen. This is because the infinite gradient at the loss cone creates intense sudden loss, and the absence of other gradients with respect to pitch angle make the appearance of

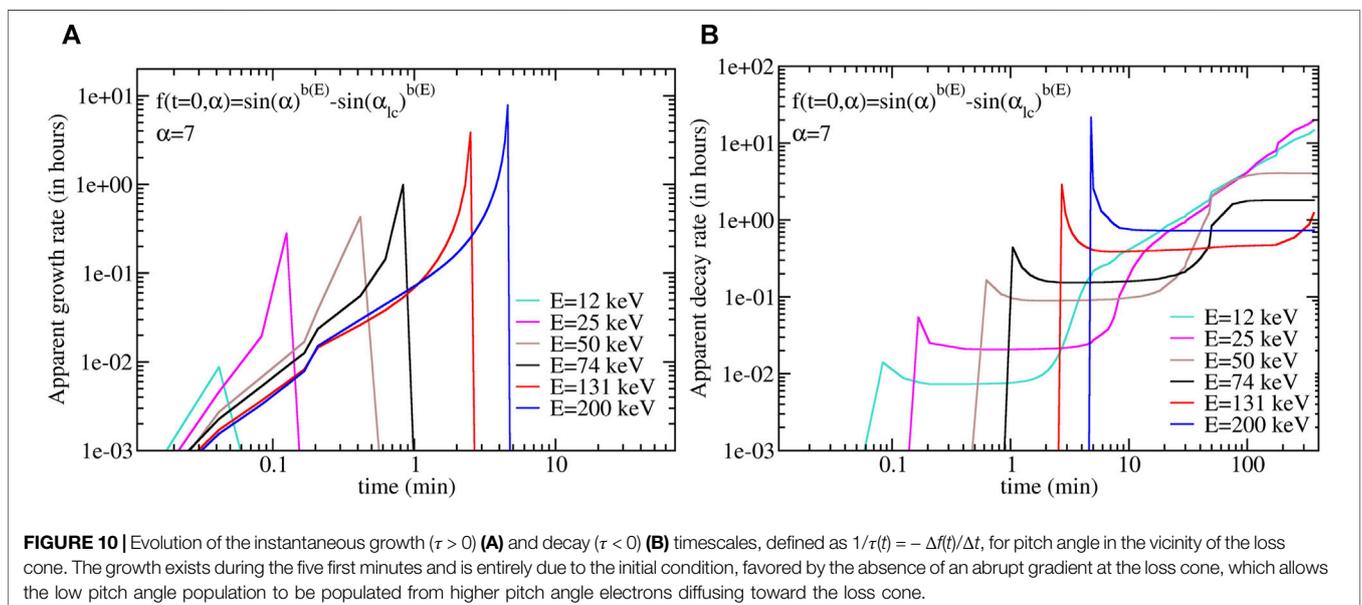


temporary PSD growth impossible. This explains why this particular behavior has not been discussed in the past (to the authors' knowledge).

The temporary growth of the PSD can be best analyzed and understood looking at **Figure 8** in which we see a temporary growth happening for times below 5 min at 135 keV for equatorial pitch angle between $\sim 12^\circ$ and $\sim 40^\circ$. The growth is due to the diffusion coefficient whose maximum (first cyclotron resonance) is located above 40° for 135 keV (cf. **Figure 5**). That diffusion turns out to be faster than the diffusion of the lower pitch angles below 40° and outside the loss cone vicinity region. High pitch angle electrons (40° – 75°) thus diffuse faster and fill out lower pitch angles ($\sim 12^\circ$ – 40°)

causing the flux at low pitch angles to rise progressively during the first 3–5 min. That type of diffusion explains the various growths of PSD we see happening in time in **Figure 7**. Comparison of **Figure 8** with **Figure 7** also shows that the growth cannot happen at $\sim 7^\circ$ for this energy and initial condition for which the absence of regularization in this example increases the gradient at the loss cone and favors the loss.

Figure 9 shows the evolution of the PSD using the specific energy-dependent $b(E)$ observed by RBSP-B. In addition, we plot in **Figure 10** the respective instantaneous growth ($\tau > 0$) and loss rates ($\tau < 0$), previously discussed. We define these rates as $1/\tau = |\Delta f(t)/\Delta t|$ in units of hours. The considerable variation of τ with time (including a change of sign) shows that steady state, which defines the general electron lifetime (Lyons et al., 1972), is not reached within the first 100 min for most of the energies below 200 keV. At 74 keV, we see, for instance, a fast loss rate of ~ 12 min persisting during ~ 30 min, which signifies that the flux decreases by a factor e^1 every 12 min during 30 min, i.e. a factor of $e^{(30/12)} = 12$ after 30 min. The constant loss rate during that time is controlled by the largest gradients, which flatten as diffusion occurs. After 30 min, the loss rate of 74 keV starts to increase and reaches a plateau value at ~ 80 min, which is the steady state value. This is in agreement with the theoretical lifetime from direct steady state computation for 74 keV electron that is 90 min for this problem. Therefore, during the first 30 min that followed the growth phase, the loss is 7.5 times faster than the lifetime would predict. After 100 min, the 49 and 74 keV electron flux are reduced by two orders of magnitude and the 100–200 keV electron flux is reduced by ~ 1 order of magnitude. The initial electron energy distribution observed by RBSP is exponential and also strongly weighted towards lower energies (**Figure 4**), thus the precipitation observed by BARREL was also dominated by these energies.



4 DISCUSSION

An event-based study was conducted to investigate wave-particle interactions in the radiation belts. A fortuitous magnetic conjunction between the RBSP spacecraft and BARREL balloon payload 3A provided a rare opportunity to study pitch-angle scattering by plume whistler-mode hiss waves. A quasi-linear pitch angle diffusion model was developed using wave and plasma parameters observed by RBSP as they passed through a stable region of dense cold plasma ($n \sim 80 \text{ cm}^{-3}$) near $L = 6$. The model results show that interaction with the plume hiss waves leads to rapid loss of electrons over the same energy range observed by BARREL. Moreover, a significant decrease in the trapped electron flux can occur within 2 h for electrons of pitch angle up to 70° for E in [50, 100] keV.

We also examined the effects of the initial pitch-angle distribution that is assumed when evolving the phase space density. Many previous studies assume a completely empty initial loss cone distribution which leads to an artificially rapid loss rate due to the very steep gradient at the loss cone. Regularization of the loss cone distribution reduces this effect and also reveals how diffusion from larger pitch angles can actually cause an initial increase in flux near the loss cone. This effect is very sensitive to the initial pitch angle distribution assumed for a given energy. These results illustrate the importance of modeling the dynamic evolution of the PSD; the assumption of steady-state loss lifetimes is not accurate for event-based studies such as this one, since the event duration is comparable to the time to reach a steady state. We incorporated the observed pitch angle distribution into the dynamic model of PSD. The present study did not include the effects of the energy-dependent drift time for electrons, rather examined the evolution of the PSD as a function of total interaction time with the waves. The lowest energy (~ 10 keV) electrons spend ~ 40 min within the wave region (the plume) while higher energy (~ 100 keV) electrons drift across the plume more quickly and may encounter the plume several times as they drift around Earth. In order to more accurately model the energy spectrum observed by BARREL, these effects must be included. The development of the model presented here will allow for such future work.

Recent work from the Van Allen Probes mission has revealed the prevalence of whistler-mode hiss waves in plasmaspheric plumes, and several studies have suggested that plume hiss waves could lead to strong loss of electrons. The event examined here occurred during a substorm injection of electrons into the radiation belts as observed by the RBSP MagEIS instruments. Thus it was not possible to separate the effects of acceleration and loss in order to quantify the overall impact of this plume on the radiation belts. However, the loss rate indicated by this study suggests that scattering by plume whistler-mode hiss waves,

which is currently not included in radiation belt models, could be an important loss mechanism, particularly for radiation belt seed electrons. In particular, the plume investigated in this study was stable and persisted for several RBSP orbits. Electrons gradient-curvature drift around the earth many times per day and thus experience multiple encounters with such high-density plumes and the hiss waves they support. Future work should investigate the overall impact of plumes as a loss mechanism for radiation belt electrons and the seed particles that generate them.

DATA AVAILABILITY STATEMENT

Publicly available datasets were analyzed in this study. This data can be found here: <https://cdaweb.gsfc.nasa.gov/pub/data/barrel/>, <https://emfisis.physics.uiowa.edu/Flight/>, and <https://rbsp-ect.newmexicoconsortium.org/science/DataDirectories.php>.

AUTHOR CONTRIBUTIONS

RM and J-FR together developed the study concept. RM carried out the basic data analysis and produced figures from BARREL and Van Allen Probes to characterize the observations. OS provided further analysis of the Van Allen Probes wave data to obtain wave parameters used in the modeling. WK processed EMFISIS data to provide cold plasma density. J-FR developed the quasi-linear diffusion model and produced the model result figures. RM and J-FR together wrote the first draft of the manuscript. All authors contributed to manuscript revision, read, and approved the submitted version.

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