



Electrospun, Oriented, Ferromagnetic Ni_{1-x}Fe_x Nanofibers

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OPEN ACCESS

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Specialty section:

This article was submitted to
Nanoscience,
a section of the journal
Frontiers in Chemistry

Received: 13 October 2019

Accepted: 15 January 2020

Published: 13 February 2020

Citation:

Bhugra VS, Hughson FR,
Williams GVM, Chong SV and Nann T
(2020) Electrospun, Oriented,
Ferromagnetic Ni_{1-x}Fe_x Nanofibers.
Front. Chem. 8:47.
doi: 10.3389/fchem.2020.00047

Electrospinning has been used to fabricate ferromagnetic Ni_{0.47}Fe_{0.53} nanofiber mats that were composed of individual, orientated Ni_{0.47}Fe_{0.53} nanofibers. The key steps were processing a polyvinylpyrrolidone nanofiber template containing ferric nitrate and nickel acetate metal precursors in Ar at 300°C and then 95% Ar: 5% H₂ at 600°C. The Ni_{0.47}Fe_{0.53} fibers were nanostructured and contained Ni_{0.47}Fe_{0.53} nanocrystals with average diameters of ~14 nm. The Ni_{0.47}Fe_{0.53} ferromagnetic mats had a high saturation magnetic moment per formula unit that was comparable to those reported in other studies of nanostructured Ni_{1-x}Fe_x. There is a small spin-disordered fraction that is typically seen in nanoscale ferromagnets and is likely to be caused by the surface of the nanofibers. There was an additional magnetic contribution that could possibly stem from a small Fe_{1-z}Ni_zO phase fraction surrounding the fibers. The coercivity was found to be enhanced when compared with the bulk material.

Keywords: magnetic properties, electrospinning, nanofibers, ferromagnetic, polyvinylpyrrolidone (PVP), nanomaterials

INTRODUCTION

There is extensive ongoing research into nanoscale magnetically ordered materials sparked by their potential applications in medicine (Pankhurst et al., 2003; Kim et al., 2007), magnetic sensors (Chen et al., 2011; Kennedy et al., 2014), RF components (Chinnasamy et al., 2015), magnetic fluids (Chen et al., 2011), and magnetic memory (Moser et al., 2002). Current research is driven by the fact that the properties of nanoscale ferromagnetic materials can be different from those seen in the bulk (Hendriksen et al., 1993; Batlle and Labarta, 2002; Goya et al., 2003; Upadhyay et al., 2016; Williams et al., 2019). For example, they have been reported to show enhanced coercivities that can be useful for magnetic memory storage devices (Barakat et al., 2009; Dong et al., 2014). The appearance of uncompensated moments or spin-disorder in the shell can lead to an exchange bias that can be used in magnetic random-access memory (RAM) (Katti, 2002; Nogués et al., 2005; Wu et al., 2013). It is also possible that the magnetocrystalline anisotropy can change particularly due to a large shell contribution that will affect the magnetic properties that include the saturation field (Batlle and Labarta, 2002; Goya et al., 2003; Demortière et al., 2011). They can even have an enhanced magnetostriction (Balaji et al., 2012), which is advantageous for composite multiferroics for applications that include magnetic sensors (Chong and Williams, 2019). If all three dimensions of the nanoscale materials are small enough, they can display superparamagnetism above a blocking temperature, T_B , where the thermal energy is greater than the magnetocrystalline anisotropy energy (Cullity and Graham, 2009). The hysteresis is negligible above T_B (Cullity and Graham, 2009),

which is advantageous for applications that include magnetic sensing where the appearance of hysteresis affects the repeatability and minimum detectable field.

Most research has focused on nanoscale materials where all three dimensions are small (Pankhurst et al., 2003; Chen et al., 2011; Kennedy et al., 2014). However, one-dimensional ferromagnetic materials in the shape of nanofibers are also particularly interesting. Their length and narrow diameters make them ideal for magnetic flux guiding applications on the nanoscale. It is also possible that they can have nanodomains extending across the diameter of the fiber and hence they could be used for nanoscale domain wall magnetic memory (Cisternas et al., 2017). Metallic and ferromagnetic nanofibers that display a degree of electronic spin polarization could also potentially be used in spin-tunneling junctions (Katti, 2002; Parkin et al., 2004; Williams et al., 2018) for magnetic RAM applications where their small size could lead to high density storage as well as lower power when compared with multilayer thin films. They could also be used for spin-tunneling nano-magnetic sensors where the long lengths in relation to the diameter can ensure flux guiding along the length of the fiber and hence magnetic sensing directionality provided that the magnetic permeability is high enough.

Electrospinning (Li et al., 2003; Graeser et al., 2007; Wu et al., 2007) and electrodeposition using an aluminum oxide or other template (Salem et al., 2012; Zhang et al., 2013; Meneses et al., 2018; Frolov et al., 2019) are two common methods to make ferromagnetic metallic nanofibers of materials that include Ni (Wu et al., 2007; Meneses et al., 2018), Fe (Graeser et al., 2007; Wu et al., 2007), Co (Graeser et al., 2007; Wu et al., 2007), and $\text{Ni}_{1-x}\text{Fe}_x$ (Salem et al., 2012; Zhang et al., 2013; Frolov et al., 2019). Electrospinning has the advantage, that it is, a relatively simpler preparation method and it is possible to fabricate larger nanofibre mats. The electrospinning process starts with preparing a polymer solution of the metal precursors, a polymer, and solvent. A high voltage is applied between the polymer solution and the collector to ensure that the electric field is large enough to result in a charged jet of the polymer solution that rapidly thins due to solvent evaporation and leads to solid fibers being deposited on the collector (Graeser et al., 2007). Continual movement of the collector leads to orientated nanofibers. Fe and Ni nanofibers have been made by electrodeposition, but we are not aware of any reports of $\text{Ni}_{1-x}\text{Fe}_x$ nanofibers being made by electrospinning. $\text{Ni}_{1-x}\text{Fe}_x$ is a well-known bimetallic material with a very large magnetic permeability and it can have a low magnetocrystalline anisotropy for $x \sim 0.25$ (Bozorth and Walker, 1953; Cullity and Graham, 2009). It also has a degree of electronic spin polarization (Žutić et al., 2004) that can lead to spin-dependent tunneling or spin-dependent scattering (Daughton et al., 1994; Inoue and Maekawa, 1996; Prakash et al., 2014), which is useful for magnetic sensors (Daughton et al., 1994). Electrospun $\text{Ni}_{1-x}\text{Fe}_x$ nanofibers that have a degree of orientation have a number of potential applications, that include very nanoflux guides that can be used, for example, for thin and compact wireless power transfer. By coating the nanofibers with a piezoelectric polymer (e.g., PVDF) it would be possible to create a magneto-electric composite as already

shown for Fe_3O_4 nanoparticle/PVDF nanocomposites (Chong and Williams, 2019) that can find application in electrically tunable miniature RF filters and antennas (Petrov et al., 2008; Guo-Min et al., 2009).

In this paper we report the successful synthesis of $\text{Ni}_{1-x}\text{Fe}_x$ mats containing orientated submicron fibers with $x = 0.53$ by the electrospinning process followed by heat treatment and reduction in a 95% Ar:5% H_2 atmosphere. This value of x was chosen because it is close to where the saturation magnetic moment is the highest while still maintaining the FCC crystal structure seen for $x < \sim 0.6$ (Li et al., 1997). We were able to show that our fibers were nanocrystalline and ferromagnetic and with a saturation magnetic moment that is comparable to that seen in other nanostructured $\text{Ni}_{1-x}\text{Fe}_x$ compounds with similar x values. We also showed that there is evidence for spin-disorder regions that may arise from the surface regions of the $\text{Ni}_{1-x}\text{Fe}_x$ fibers.

EXPERIMENTAL DETAILS

Materials

All the chemicals used in the reaction were of analytical grade. The polyvinylpyrrolidone (PVP, $M_w = 1,300,000$ —polymer for electrospinning), iron(III) nitrate nonahydrate ($\text{Fe}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$), >98%—precursor for Fe and nickel(II) acetate tetrahydrate ($\text{C}_4\text{H}_{14}\text{NiO}_8 \cdot 4\text{H}_2\text{O}$), 98%—precursor for Ni were obtained from Sigma-Aldrich. Methanol was obtained from Chem Supply.

Experimental Work

In a typical procedure, 2 g (5 mmol) of Fe nitrate and 1.5 g (6 mmol) of Ni acetate were mixed in 10 ml of methanol and 0.5 g polyvinylpyrrolidone (PVP). The whole mixture was stirred overnight at 300 rpm and loaded into a 5 mL plastic syringe fitted with a 21-gauge stainless steel needle. The needle was connected to a high-voltage supply (Gamma High Voltage Research, Ormand Beach, FL). A rotating drum was chosen as a collector, which acted as a counter electrode to collect highly charged fibers and to create aligned fibers. The rotating drum was wrapped with aluminum foil prior to fiber collection, which is advantageous for post electrospinning treatment of the fibers and it also avoids mechanically induced damage. The syringe was attached to a syringe pump (Harvard Apparatus) as shown in **Figure 1** and the solution was pumped out at a rate of 0.65 ml/h. The voltage for electrospinning was set to 17.5 kV and the distance between the electrodes was fixed to 10 cm. The sample was transferred into a tube furnace after electrospinning and heated in an argon atmosphere at a rate of $10^\circ\text{C}/\text{min}$ and then held for 1.5 h at 300°C . This was then followed by increasing the temperature of the furnace at a rate of $10^\circ\text{C}/\text{min}$ to 600°C in a 95% Ar:5% H_2 atmosphere and holding for 3 h. The furnace was then switched off and allowed to cool down to room temperature. The resultant sample was black and brittle.

Characterization

X-Ray diffraction (XRD) was performed on a Panalytical X-Ray diffractometer using Cu-K_α radiation with an operating voltage

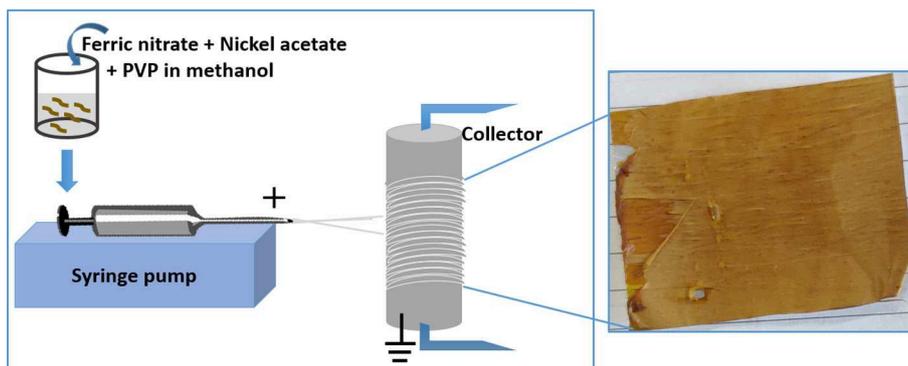


FIGURE 1 | Electrospinning setup for making PVP/Fe nitrate/Ni acetate nanofibers.

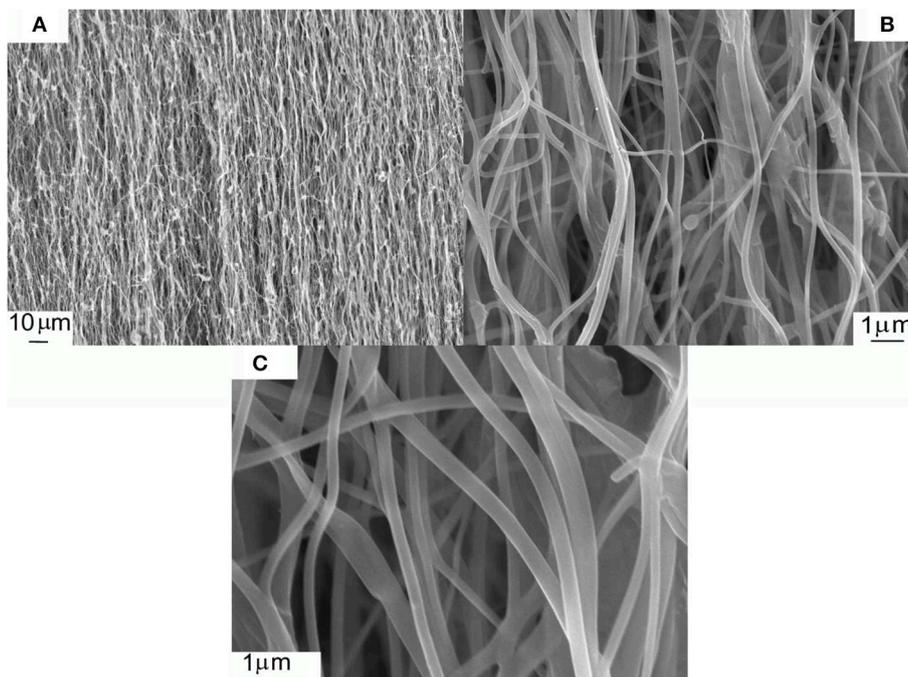


FIGURE 2 | SEM micrographs of PVP/Fe nitrate/Ni acetate electrospun nanofibers. (A–C) show different magnifications of the same sample.

of 45 kV and current of 40 mA, at room temperature, to study the crystal structure of the electrospun fibers. Part of the sample was ground before the XRD measurement so that a powder pattern could be obtained for phase identification. Scanning electron microscopy (SEM) measurements were performed on a Jeol SEM-6500 to study the morphologies of the fibers. Energy dispersive X-ray spectroscopic (EDXS) analysis and scanning transmission electron microscopy (STEM) mapping of the sample were carried out using the same equipment for SEM imaging to study the elemental composition of the fibers. Magnetic measurements were made using a magnetic property measurement system (MPMS) from Quantum Design. The sample was sealed in a gelatine capsule and placed inside a straw and then loaded into the MPMS.

RESULTS AND ANALYSIS

It can be seen in the SEM micrographs in **Figure 2A** that electrospinning has resulted in orientated PVP/Fe nitrate/Ni acetate fibers. A higher magnification micrograph in **Figure 2B** shows that there are some nanofibers that were not in the general orientated direction and that there is also some fusing of nanofibers. **Figure 2C** shows that the nanofiber diameters range from ~ 200 to ~ 250 nm.

The sample morphology of the electrospun fibers dramatically changed after processing at 600°C in 95% Ar/5% H_2 as can be seen in **Figure 3A**. It is apparent that the fibers seen before processing have coalesced to form larger oriented submicron fibers with diameters ranging from ~ 600 to ~ 900 nm. These

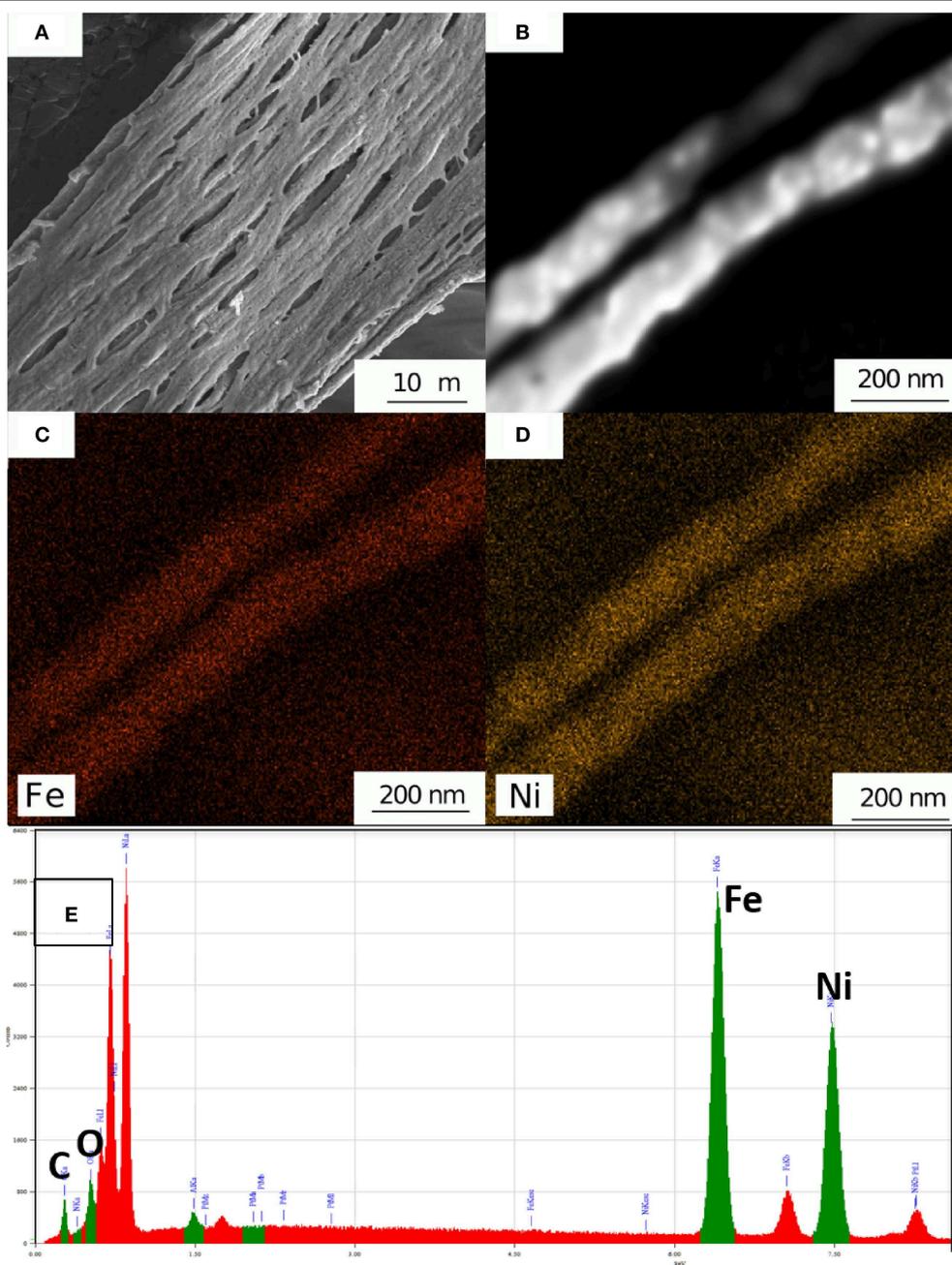
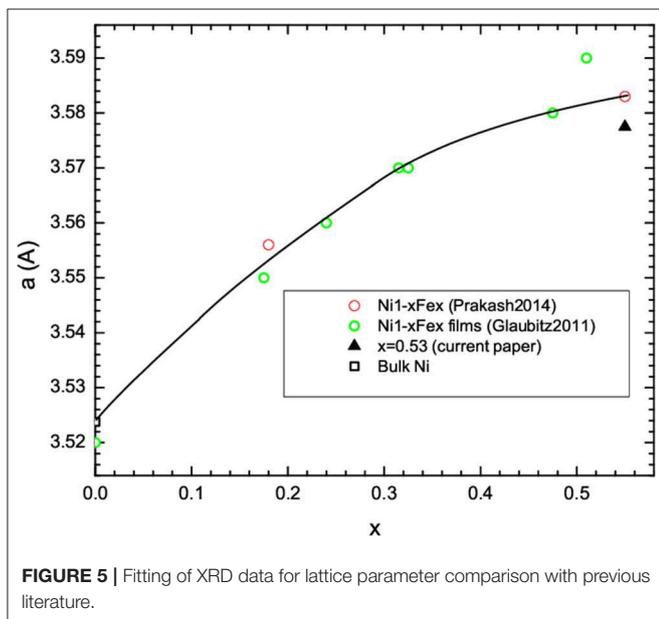
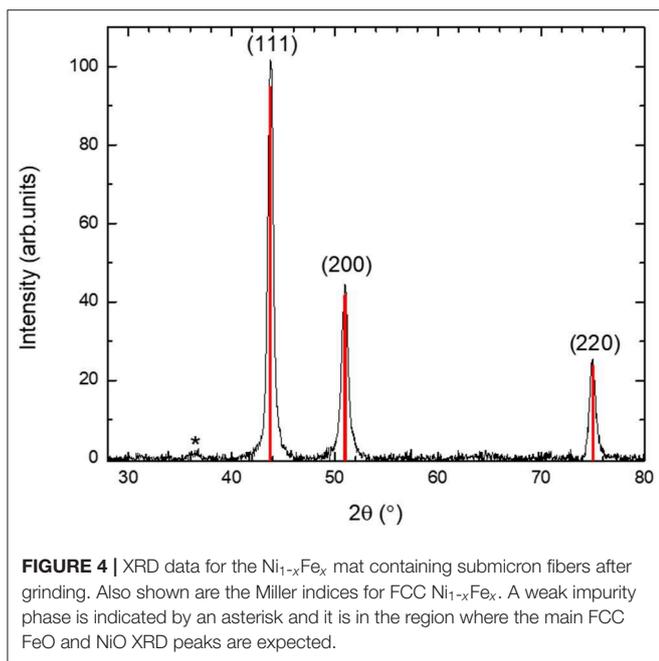


FIGURE 3 | (A) SEM micrograph after high temperature processing in an Argon/H₂ atmosphere leading to a Ni_{1-x}Fe_x mat containing submicron fibers. **(B)** A SEM micrograph of two Ni_{1-x}Fe_x fibers and STEM elemental maps of **(C)** Fe and **(D)** Ni. EDX spectra of Ni_{1-x}Fe_x fibers.

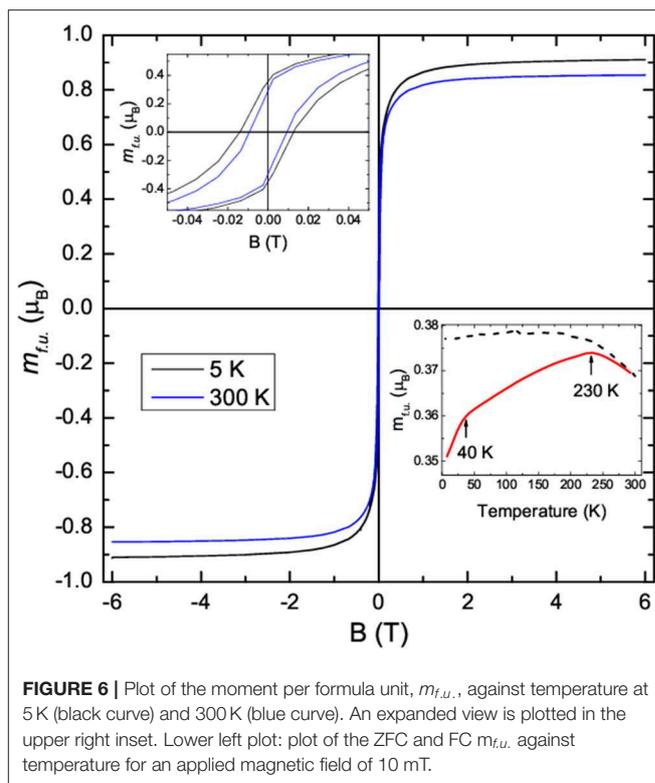
fibers have fused to create a mat with significant voids. There are a few small branching nanofibers with diameters in the ~300 nm range. A small area (**Figure 3B**) was selected for scanning transmission electron microscopy (STEM) mapping and the STEM maps are shown in **Figure 3C** for Fe and **Figure 3D** for Ni are shown in **Figure 3B**. It can be seen that Fe and Ni are uniformly distributed within the fiber. Energy dispersive X-ray spectroscopy (EDXS) showed that the fibers were Ni_{1-x}Fe_x

with $x = 0.53 (\pm 0.03)$, which is the value expected from the initial precursor ratio (**Figure 3E**). The carbon peak was small and shows that there was removal of significant fraction of the polymer. The low carbon signal was likely from the SEM sample handling and processing (for example, the adhesive on the tape holding the sample on SEM stub).

The results from XRD measurements after the Ni_{0.47}Fe_{0.53} nanofibers synthesis are shown in **Figure 4**. The main phase



can be fitted to FCC $\text{Ni}_{1-x}\text{Fe}_x$, which was expected for $x < 0.6$ (Li et al., 1997). Fitting of the XRD data shows that the lattice parameter was $a = 3.5775 \text{ \AA}$. By comparison with data in the literature (Glaubit et al., 2011; Prakash et al., 2014), we found that this lattice parameter is in the range expected for $x = 0.53$ (Figure 5). Since XRD is an averaging probe over the whole sample, this, when combined with the EDXS and STEM map data, shows that the x content did not significantly vary across the sample. A very small and broad peak was found at 2θ of $\sim 36.5^\circ$. This may be due to a small fraction of $\text{Fe}_{1-z}\text{Ni}_z\text{O}$ because FCC NiO and FeO have (111) peaks in this region.



Broadening of the XRD peaks in Figure 4 was observed, which arises in materials with small nanocrystals and/or strain. For this reason, the XRD peaks were fitted to the pseudo Voigt functions to obtain the linewidth. A Williamson-Hall analysis (Birkholz et al., 2006) showed that there is negligible strain and the $\text{Ni}_{0.47}\text{Fe}_{0.53}$ fibers were composed of nanocrystals with average diameters of $\sim 14 \text{ nm}$.

Magnetic measurements were performed on the as-made $\text{Ni}_{0.47}\text{Fe}_{0.53}$ fiber mats. The moment per formula unit, $m_{f,u}$, was plotted as a function of applied magnetic field, B , in Figure 6 at 5 and 300 K. $m_{f,u}$ was plotted in Bohr magneton units, μ_B . $m_{f,u}$ was calculated as the moment/(no. of moles \times Avogadro's no.). The data clearly shows the presence of ferromagnetic order that is expected for $\text{Ni}_{1-x}\text{Fe}_x$. The high field magnetic moment, $m_{f,u,sat}$, was large and it was $0.91 \mu_B$ at 5 K. This was less than typical bulk samples, where $m_{f,u,sat} = 1.78 \mu_B$ for $\text{Ni}_{1-x}\text{Fe}_x$ with $x = 0.53$ (Li et al., 1997). It is comparable to that reported for $\sim 2.7 \text{ nm}$ diameter $\text{Ni}_{1-x}\text{Fe}_x$ nanoparticles with $x = 0.47$ and made by dual ion beam implantation where $m_{f,u,sat} = 1.0 \mu_B$ (Williams et al., 2019). Nanoscale magnetic materials typically have lower $m_{f,u,sat}$ than the bulk (Vitta et al., 2008; Demortière et al., 2011; Upadhyay et al., 2016; Williams et al., 2018).

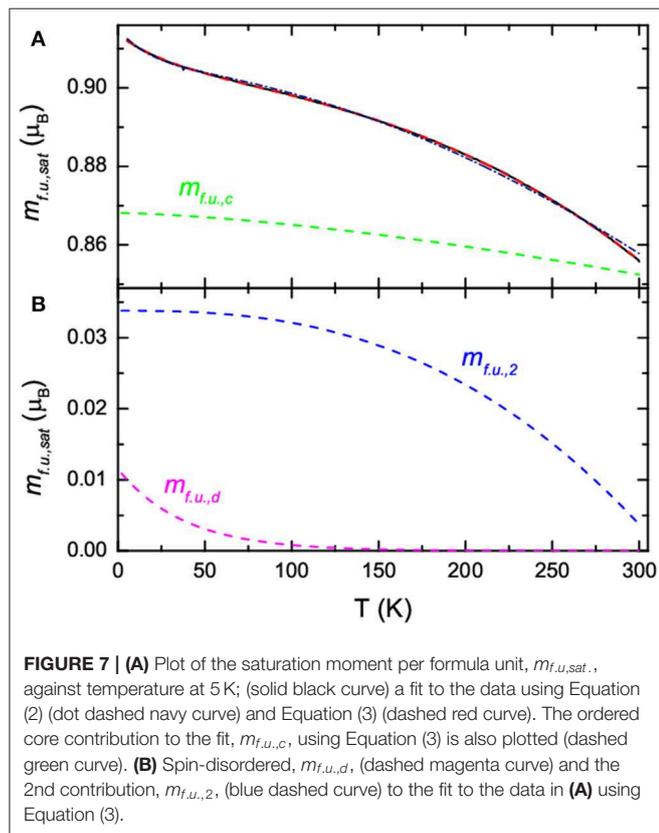
The upper left inset to Figure 6 shows that the coercive field was enhanced when compared with the bulk compound where the coercive field and the saturation field are small. The enhanced coercivity is likely to be due a high density of domain wall pinning sites at the nanograin boundaries. A larger coercivity was observed at low temperatures because domain wall motion is thermally activated as well as occurring via the

applied magnetic field. The zero-field-cooled (ZFC) and field-cooled (FC) $m_{f.u.}$ is plotted in the lower right inset to **Figure 6**. The ZFC data were taken after cooling in zero applied field to low temperatures. A magnetic field of 10 mT was then applied and magnetic measurements were taken while warming up to 300 K followed by field-cooling down to low temperatures. Magnetic hysteresis is still evident at the highest temperature and indicates that the nanocrystals are too large for superparamagnetism to occur.

Two features can be seen in the ZFC $m_{f.u.}$ data. The first is a peak at ~ 230 K and the second is a faster decrease below ~ 40 K. The ~ 230 K peak was likely to arise from a small antiferromagnetic phase fraction. It is too small to be purely from NiO, which has a Neel temperature of 525 K. It is also higher than that seen in bulk FeO where the Neel temperature is 198 K. Thus, it may have been arisen from a $\text{Fe}_{1-z}\text{Ni}_z\text{O}$ phase that is seen in the XRD data and that would be expected to have a Neel temperature between 198 and 525 K. The more rapid decrease below 40 K may be related to a small spin-disorder region that is discussed below.

The saturation magnetic moment per formula unit, $m_{f.u.,sat}$, obtained from the moment per formula unit at 6 T, is plotted in **Figure 7A** as a function of temperature. The temperature dependence of $m_{f.u.,sat}$ for bulk $\text{Ni}_{1-x}\text{Fe}_x$ is known to follow the Bloch function (Bloch, 1930; Ashcroft and Mermin, 1976) that can be written as,

$$m_{f.u.,sat}(T) = m_{f.u.,sat}(0) [1 - \beta \times T^n] \quad (1)$$



where n is the exponent, T is the temperature, $m_{f.u.,sat}(0)$ is $m_{f.u.,sat}$ at 0 K, and β is the Bloch constant. $n = 3/2$ is expected for $\text{Ni}_{1-x}\text{Fe}_x$ and for simple ferromagnetic materials (Ashcroft and Mermin, 1976). β is proportional to $1/D^{3/2}$ where D is the spin stiffness that in turn is proportional to the exchange energy (Srivastava and Aiyar, 1987; Demortière et al., 2011).

It is clear in **Figure 7A** that $m_{f.u.,sat}$ does not follow Equation (1) due to the upturn at lower temperatures. This type of behavior has been reported before in ferromagnetic nanoparticles (Vázquez-Vázquez et al., 2011; Larumbe et al., 2012; Williams et al., 2019). It has been shown that $m_{f.u.,sat}$ can be fitted to (Vázquez-Vázquez et al., 2011; Larumbe et al., 2012),

$$m_{f.u.,sat}(T) = m_{f.u.,c}(0) \times [1 - \beta \times T^n] + m_{f.u.,d}(0) \times \exp(-T/T_f) \quad (2)$$

where the first term is the Bloch function and the second term is a phenomenological term to account for spin-disorder in the shell region. $m_{f.u.,c}(0)$ is the saturation moment per formula unit in the core at 0 K, $m_{f.u.,d}(0)$ is the spin-disordered saturation moment per formula unit at 0 K, and T_f is the characteristic spin-freezing temperature.

We showed in **Figure 7A** that Equation (2) provides an approximate fit to the data (navy dot dashed curve) with $m_{f.u.,c}(0) = 0.905 \mu_B$, $\beta = 2.0 \times 10^{-6} \text{ K}^{-1.78}$, $n = 1.78$, $m_{f.u.,d}(0) = 0.0099 \mu_B$, and $T_f = 17$ K. However, the value of n is larger than that seen in the bulk or $\text{Ni}_{1-x}\text{Fe}_x$ nanoparticles with diameters >35 nm (Vitta et al., 2008). It is possible that the larger n and approximate fit to the data is due to a small contribution from another magnetically ordered phase. This could be $\text{Fe}_{1-z}\text{Ni}_z\text{O}$ that can have a measurable magnetic moment due to the moments not being colinear. In this case, an extra term is added to Equation (2) leading to

$$m_{f.u.,sat}(T) = m_{f.u.,c}(0) \times [1 - \beta \times T^n] + m_{f.u.,d}(0) \times \exp\left(\frac{-T}{T_f}\right) + m_{f.u.,2}(0) \times [1 - \beta_2 \times T^{n_2}] \quad (3)$$

where $m_{f.u.,2}(0)$ is the saturation moment per formula unit from the second smaller phase fraction at 0 K, β_2 is the temperature prefactor, and n_2 is the temperature exponent.

We showed in **Figure 7** that Equation (3) provides an excellent fit to the data (red dashed curve) with $m_{f.u.,c}(0) = 0.868 \mu_B$, $\beta = 3.5 \times 10^{-6} \text{ K}^{-1.5}$, $n = 1.5$, $m_{f.u.,d}(0) = 0.011 \mu_B$, $T_f = 38$ K, $m_{f.u.,2}(0) = 0.034 \mu_B$, $\beta_2 = 3.0 \times 10^{-6} \text{ K}^{-2.6}$, and $n_2 = 2.6$. In this case n is the same as that found in the bulk and in $\text{Ni}_{1-x}\text{Fe}_x$ nanoparticles (Vitta et al., 2008; Williams et al., 2019), and β is in the range found for bulk $\text{Ni}_{1-x}\text{Fe}_x$ (Zhang et al., 1998; Vitta et al., 2008). The value of T_f is where a faster decrease in the ZFC $m_{f.u.}$ was observed. **Figures 7A,B** also show the contribution to the total fit from the three terms in Equation (3). The dominant contribution is from the ferromagnetic $\text{Ni}_{1-x}\text{Fe}_x$ core (green dashed curve in **Figure 7A**). The spin-disordered contribution is small.

The low value of $m_{f.u.,d}(0)$ suggests that it is unlikely that the spin-disordered contribution is from the shell regions between the $\text{Ni}_{1-x}\text{Fe}_x$ nanocrystals. A simple estimate of the

shell thickness can be obtained using a model developed for nanoparticles where the average spin-disordered shell thickness, t_{shell} can be estimated from $t_{shell} = (1 - \gamma^{1/3}) \times r$ where $\gamma = m_{f.u.,c}(0) / [m_{f.u.,c}(0) + m_{f.u.,d}(0)]$ and r is the average radius of the nanocrystals (Prakash et al., 2016). We find that t_{shell} is only 0.03 nm for an average r of 7 nm, which is unrealistically small. It is possible that the spin-disordered region surrounds the $Ni_{1-x}Fe_x$ fibers. In this case the simple model gives,

$$t_{shell} = (1 - \gamma^{1/2}) \times r \quad (4)$$

where t_{shell} is now thickness of the spin-disordered region surrounding the fibers and r is the radius of the fibers. Using an average fiber diameter of 750 nm, this gives $t_{shell} = 2.3$ nm, which is plausible.

The temperature dependence of the additional contribution to the fit to the data in **Figure 7A** was plotted in **Figure 7B** (blue dashed curve). This contribution is small and close to zero at 300 K. This suggests a magnetic phase with a magnetic ordering temperature at or slightly below 300 K. It may be that it is due to $Fe_{1-z}Ni_zO$ surrounding the fibers with a range of magnetic ordering temperatures up to ~ 300 K. In this case an estimate of the average thickness of the oxidized region surrounding the fibers can be obtained from Equation (4). Using an average fiber diameter of 750 nm, we find that $t_{shell} \sim 7$ nm.

CONCLUSIONS

In conclusion, ferromagnetic $Ni_{0.47}Fe_{0.53}$ mats containing orientated $Ni_{0.47}Fe_{0.53}$ nanofibers were fabricated by electrospinning solutions of ferric nitrate and nickel acetate metal precursors in polyvinylpyrrolidone (PVP) polymer solutions followed by thermal processing. The final synthesis step at 600°C in 95% Ar:5% H_2 led to the complete removal of the polymer and only $Ni_{0.47}Fe_{0.53}$ fibers and a small fraction of $Fe_{1-z}Ni_zO$ remained. There were a range of $Ni_{0.47}Fe_{0.53}$ nanofiber diameters that varied from ~ 600 to ~ 900 nm where

they had fused to form the $Ni_{0.47}Fe_{0.53}$ mats. The fibers were nanostructured and contained nanocrystals with average diameters of ~ 14 nm as estimated from the XRD data. Magnetic measurements clearly showed the presence of ferromagnetic order. The saturation magnetic moment per formula was less than that found in the bulk but it was similar to that reported in other studies of nanostructured $Ni_{1-x}Fe_x$ with similar x values. Modeling of the temperature dependence of the saturation moment showed that there was some spin-disorder, probably from a thin layer in the surface region of the fibers, with a characteristic spin-freezing temperature of 38 K. There was also an additional minor magnetically ordered phase with a magnetic ordering temperature below room temperature that may be from a thin antiferromagnetic $Fe_{1-z}Ni_zO$ layer surrounding the nanofibers. The coercivity was enhanced when compared with the bulk, which is probably due to the small nanocrystals that lead to a high density of domain wall pinning centers.

DATA AVAILABILITY STATEMENT

The datasets generated for this study are available on request to the corresponding author.

AUTHOR CONTRIBUTIONS

FH undertook initial experiments to electrospin magnetic nanofibers. VB undertook the electrospinning, SEM, and XRD measurements. GW did the XRD and magnetic analysis and interpretation. SC did the magnetic measurements and assisted in the magnetic interpretation. TN conceived and helped to oversee the research. All authors were involved in the data interpretation.

FUNDING

This study was funded by the New Zealand Ministry of Business, Innovation and Enterprise (UOAX1711).

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Conflict of Interest: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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