

# Tracing the Anharmonicity and Superionic Phase Transition of Hydrous FeO<sub>2</sub>H

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The weak x-ray scattering of hydrogen (H) has brought major challenges to the characterization of superionic transitions in high-pressure ice, hydrides, and hydroxides. Combining first-principles molecular dynamics and simulated nuclear magnetic resonance (NMR) spectroscopy, we investigated the behavior of the hydroxyl bonding and structural transitions in the hydrous FeO<sub>2</sub>H between 300 and 2750 K and up to 130 GPa. Evidence show that an intermediate plastic state with regional H diffusion and anharmonic O-H vibration exists in between the ordinary solid and the superionic phase. The intermediate state features asymmetric hydrogen bonds and anharmonic vibrations, which are readily distinguished from the high-temperature superionic phase. Our work shows NMR is a more sensitive probe to detect H diffusion in superionic solids even in the extreme conditions of Earth's deep interiors.

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# INTRODUCTION

A substantial portion of H in the Universe may diffuse like a liquid in superionic phases (Cavazzoni et al., 1999; Liu et al., 2020; Cong Liu et al., 2019; Millot et al., 2019; Millot et al., 2018). For example, superionic ice is predicted to be a dominant component of giant icy planets like Uranus and Neptune (Cavazzoni et al., 1999). Superionic hydrous mineral phases are also found to be stable in Earth's deep lower mantle (Hou et al., 2021; Hu and Mao, 2021). The superionization is signified by the exceptionally high H diffusion rate and soaring ionic conductivity as the H atom moves freely in the host lattice and generates protonic currents. Although the conception of H in the superionic phase is postulated theoretically in the 1980s (Demontis et al., 1988), its optical properties, structural transition, and electrical conductivity are only determined by experiments very recently (Millot et al., 2019; Millot et al., 2018; Zhuang et al., 2022a). Properties like the evolution of H bonding, coreelectron interaction, and chemical environments under high pressure-temperature (P-T) conditions are still inconclusive.

The major challenge to study the superionic phase is the experimental probe. Traditional x-ray probes, such as x-ray diffraction (XRD) have been employed as routine tools in characterizing structures under pressure (Mao et al., 2016). However, H is the lightest element, which scatters a very limited amount of x-ray. Even synchrotron-based XRD is unable to measure H motion but only the solid crystalline lattice in the superionic phase (Millot et al., 2018). In short, the behavior of H, particularly as a function of temperature, is still difficult to be described by experiments. Pioneering works used shock compression and optical measurements to detect the onset of superionic ice, but the *P*-*T* points have to follow the shock Hugoniot curve (Coppari et al., 2021; Millot et al., 2019). It is also possible to conduct electrical conductivity measurement and laser heating in a diamond anvil

cell (Zhuang et al., 2022b) to mark the superionic transition. However, factors like external electrical field and temperature gradient make such analysis much more challenging than other systems (Futera et al., 2020; Hou et al., 2021).

Nuclear magnetic resonance (NMR) spectroscopy is a versatile means to study the structural environments of solids and liquids (Harris, 2004). Owing to the invention of the electromagnetic Lenz lens, it is possible to refine the NMR resonator in a modified diamond anvil cell (Meier et al., 2017; Meier et al., 2021). One important advantage of NMR is that it detects the chemical environments of H with much higher accuracy than XRD (Chen et al., 2020). Also, first-principles calculations can obtain reliable NMR parameters, which are predictable to phase transitions relating to H (Holmes et al., 2017; Meier et al., 2019). For example, Meier et al. 2019 investigated the hydrogen interactions of FeH up to 202 GPa using NMR and revealed the formation of a caged H framework around Fe atoms. Both their calculation and experiments showed that the phase transition induced by H motion can be distinguished by chemical shielding in different structures.

Goethite, or  $\alpha$ -FeO<sub>2</sub>H, is a key component of the banded iron formations (Bekker et al., 2010), and it exhibits a set of structural phase transitions in Earth's deep interiors (Lu and Chen, 2018; Thompson et al., 2020). FeOOH is thus a potential carrier of water throughout the mantle (Zhuang et al., 2022a). The pyrite-type FeO<sub>2</sub>H is stable at subduction slabs P-T conditions at 1800 km depth or deeper (Hu and Liu, 2021). Under warmer geothermal conditions, for example, at 81 GPa and 2100 K, a recent study has shown FeO<sub>2</sub>H decomposes to mixed valence iron oxides and oxygen-rich fluids (Huang and Hu, 2022; Koemets et al., 2021a). The chemical stability of FeO<sub>2</sub>H may extend to deeper regions above the core-mantle boundary by incorporating Mg (Hu et al., 2021) or inertia elements like He or Xe into the pyrite-type lattice (Peng et al., 2020; Zhang et al., 2018). At the lowermost mantle, the phase transition and crystal chemistry of FeO<sub>2</sub>H are crucial to the interaction between the mantle and the core and may contribute to the light element budget in the outer core (Mao et al., 2017; Nishi et al., 2017). It is, therefore, necessary to study the high P-T structural transformation of FeO<sub>2</sub>H at relevant lower mantle conditions.

In this work, we focused on the temperature-induced transition of pyrite-type FeO<sub>2</sub>H (Hu et al., 2016; Nishi et al., 2017), which is found to enter the conductive superionic phase at above 80 GPa and ~1800 K (Hou et al., 2021; Hu and Mao, 2021). The temperature resolved NMR spectra and their corresponding structures are calculated by first-principles calculation and molecular dynamics (MD). Prior to the superionic state, FeO<sub>2</sub>H exhibits an intermediate phase with broadened NMR peak and anharmonic motion and regional diffusion of H. The hierarchical changes of NMR peaks can be used to precisely locate the transition boundaries of the superionic phase in the experiment.

## COMPUTATIONAL METHODS

Although H in the hydrous  $FeO_2H$  is reported to be nonstoichiometric (Hu et al., 2017; Tang et al., 2021), we considered the full occupation of H in this work to simplify the structural model. The previous spectroscopic experiment also reported that the presence of H has a minor effect on the crystal chemistry of the system of FeO2-FeO2H (Liu et al., 2019). Our system consisted of 32 FeO<sub>2</sub>H units in a  $2 \times 2 \times 2$ supercell. The simulations are based on density functional theory in the PBE approximation (Perdew et al., 1996). We used the projector augmented wave (PAW) method (Kresse and Joubert, 1999) as implemented in Vienna's ab initio simulation package (Kresse and Furthmüller, 1996). The core radii are Fe: 1.302 Å (3s23p2), O: 0.741 Å (2s<sup>2</sup>2p<sup>4</sup>), and H: 0.370 Å (1s<sup>1</sup>). The basis-set energy cutoff is set as 700 eV, which is found sufficient to converge the total energy and pressure to within 5 meV/atom and  $\pm 0.5$  GPa. Born-Oppenhiemer molecular dynamics calculations are performed in the canonical ensemble using the Nose-Hoover thermostat (Hoover, 1985) and allowed to run for 20-30 picoseconds (ps) with 1-femtosecond (fs) time step. Since H is extremely mobile under high P-T conditions, we also run the test at 0.5 fs with longer trajectories at selected conditions. The structural features obtained from shorter MD timestep are consistent with the one using 1 fs MD timestep. We, therefore, used 1 fs timestep throughout the simulation. The Brillouin zone is sampled only at the Gamma point for MD simulation. It generally takes 5 ps to heat the system to the target temperature. After reaching the temperature, we adjusted lattice parameters to ensure that the stress is pseudo-hydrostatic along all axes, and the system is equilibrated at desired pressures (with ±2 GPa uncertainty). Thermal equilibrium between ions and electrons is judged by the Mermin functional (Mermin, 1965). The adjusting process takes another 5-10 ps. After fully equilibrated, the last 10 ps run is used for data production.

NMR shielding is calculated by using the gauge including the projector augmented wave method (Bonhomme et al., 2012; Pickard and Mauri, 2001). Increased cut-off energy of 900 eV and a *k*-point spacing of 0.3 Å<sup>-1</sup> are used. A denser *k*-point mesh with a spacing of 0.15 Å<sup>-1</sup> is also tested, but it does not generate qualitative different results. In our simulation, we extracted snapshots from an MD trajectory and average NMR shielding from those configurations (Dumez and Pickard, 2009). The snapshots are taken every 200 fs, corresponding to a total number of 50 configurations from the production run. The relatively longer spacing between snapshots is used to ensure they are time uncorrelated (Allen and Tildesley, 1987).

## TEMPERATURE-RESOLVED FIRST-PRINCIPLES MD

Following our previous work (Hou et al., 2021), we conducted high *P*-*T* first-principles MD at 70, 100, and 130 GPa, and temperatures up to 2750 K. We first monitored the meansquared displacements of Fe, O, and H atoms, which have been commonly used to determine the solid-superionic transition. Similar to our previous work, the phase boundary between the ordered solid and the superionic phase followed a positive dP/dT slope, and H atoms exhibited cross-boundary diffusion when temperatures are above 2000 K (Hou et al., 2021). For the superionic state, we took snapshots along the



FIGURE 1 | Evolution of NMR spectroscopy and their corresponding structures. (A). Calculated NMR spectroscopy at 100 GPa and up to 2750 K. (B–E) Selected structures at 300, 1250, 1750, and 2,500 K. Hydrogen diffusion increases substantially at higher temperatures.



trajectory which clearly showed the crystal lattice of superionic FeO<sub>2</sub>H is filled with fast-moving protons (**Figure 1**). The chemical shift  $\delta$ , which is related to the isotropic NMR shielding  $\sigma_{iso}$  by an arbitrary reference  $\sigma_{ref}$  was derived at each temperature:

$$\delta = \sigma_{\rm ref} - \sigma_{\rm iso}$$

In contrast to the superionic transition determined from H diffusion, the shape of the NMR peak is substantially broadened at much higher temperatures. For example, the full width at half maximum at 1000 K, 100 GPa increased by a factor of four compared to the one at 300 K (**Figures 1A**, **2**). This is approximately 1,000 K before it enters the previous known



superionic state (Hou et al., 2021). It is worth noting that at three individual pressures of 70, 100, and 130 GPa, such temperatureinduced broadening all occurred between 650 and 1,000 K (Figure 2), and the transition turns out to be insensitive to a pressure similar to ice (Bove et al., 2013). From the Clausius-Clapeyron relation, the sharp slope  $dP/dT = \Delta S/\Delta V$ indicated that the transition induces negligible volume change, which is the signature of a second-order phase transition. Within the P-T regions of this work, they formed a nearly horizontal boundary between 650 and 1,000 K. Below this temperature all the H atoms are symmetrically bonded with two O atoms. Such a symmetric bond corresponds to the single, sharp NMR peak. When the center of the peak is almost unchanged as the temperature is raised to 1,000 K and above, the broadened NMR peak is regarded as an assembly of multiple split peaks with comparable chemical shielding. The shape of the broad peak is similar to that found in glass (Youngman, 2018), suggesting the type of H bonding has become variant.

## LOCAL ENVIRONMENT OF H

We are then motivated to construct the local chemical environment of H in variable temperatures. Angular distribution of O-H-O is a direct probe to study the evolution of H bonds with temperature. In the case of symmetric H bonding, the angle formed by O-H-O bonds should be  $180^{\circ}$ . However, H bonds in FeO<sub>2</sub>H connect two FeO<sub>6</sub> polyhedra and thus are inevitably affected by the crystal field. At 100 GPa and 300 K (**Figures 3A, D**), peak values are centered around  $170^{\circ}$ , which is a little bit off the ideal  $180^{\circ}$ . In a quasi-harmonic model,



the thermal vibration of atoms caused the distribution of angles to follow the normal distribution (**Figure 3D**). The normal distribution is valid until the temperature raised above 1,000 K, where the second peak of angular distribution

occurred around 110° (**Figure 3C**). This means the vibration of the O-H bond is no longer described by the classic harmonic vibration model. When H atoms gain enough kinetic energy to overcome the harmonic potential well, they are still constrained by neighboring O atoms and move regionally in an anharmonic potential well (**Figures 1C**, **4**). Unlike the superionic state in which H atoms perform cross-boundary diffusion, this intermediate state can be regarded as an anharmonic state or plastic state described by the motion of protonic H freely moving around their center of mass. This plastic state has been well documented in helium ammonia compounds under high pressure and high temperature (Liu et al., 2020).

The peak near 110° in the angular distribution continues to grow and becomes significant in the superionic state. It is worth pointing out that we set a threshold of 2.0 Å in O-H bond length for searching H bonding, which is usually large enough for a compressed system. Considering the interstitial O···O distance is ~2.1 Å at 100 GPa and 2,500 K, the cutoff angle in the searching radar is ~60°. We can possibly trace the motion of H in the superionic phase by fitting the angular distribution into two Gaussian peaks, which are located at 110° and 155°, respectively (**Figure 3B**). In a system with rapidly moving H, very few protons are concentrated at the free-standing sites of the lattice, which had small O-H-O angles (e.g. <90°) and non-interacted O-H bonds. However, they preferred to stay in the original space in between two FeO<sub>6</sub> polyhedra while moving to neighboring sites by hopping. In this scenario, the H atom is still weakly interacted with at least one O atom.

#### **DISCUSSION AND CONCLUSION**

Routine high-pressure probes like XRD and electrical conductivity measurement can identify the onset of the superionic state caused by H diffusion. However, they fail to directly detect the local chemical environment of H, which is critical in describing the high P-T phase transitions. Neutron scattering would be an alternative approach; however, it is too expensive to conduct neutron experiments at the sensitive pressure of the aforementioned superionic phases (Hattori et al., 2019). In short, for its high sensitivity, NMR spectroscopy would be a perfect tool for measuring the H site and describing H motion under high-pressure conditions.

Our simulation is performed on the stoichiometric system of  $FeO_2H$ , in which H takes full occupation. This is an approximation to the experimental observation that H becomes highly mobile under high-temperature conditions (Tang et al., 2021). In laboratory experiments, the content of H is reported to be fully stoichiometric (Nishi et al., 2017), nearly stoichiometric (Yuan et al., 2018), or partially dehydrated

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(Koemets et al., 2021b), depending on the *P*-*V* and sample conditions (e.g., coating and pressure medium). The variant H contents have become a signature of dense hydrous materials which modulate the mineral structural transition and local redox condition. For example, the superionic transition temperature of FeO<sub>2</sub>H is lowered (e.g., by ~150 K at 100 GPa) after losing a portion of H (Hou et al., 2021). Such nonstoichiometry is likely to have similar effects on the transition boundary of the plastic state. The onset temperature of the plastic state might be lower than our prediction in the experiment.

Our computational work has not only verified the superionic transition previously determined by experiment but also identified a harmonic to plastic state transition in hydrous FeO<sub>2</sub>H. The results are consistent with Raman spectroscopy (Hou et al., 2021), which shows hydroxyl mode softening even at ambient temperature and high pressure. While superionic transition in FeO<sub>2</sub>H and probably other hydroxides usually occur at high temperatures, they may have precursor structures with fragile H bonding. The transition from ordered solid, via plastic state to the superionic phase, is in line with the archetypal ice water. Compared with plastic ice (Bove et al., 2013), hydrous minerals like FeO<sub>2</sub>H may also feature higher shear viscosity under moderate heating. Our simulation suggested the set of transitions is readily captured by the change in the NMR spectrum. The use of NMR spectroscopy will be more sensitive than x-ray probes in characterizing potential superionic materials. We proposed the plastic and superionic phase transition would be a universal phenomenon in the highpressure ice, hydrides, and hydroxides.

## DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/supplementary material; further inquiries can be directed to the corresponding author.

# **AUTHOR CONTRIBUTIONS**

QH conceived the project and carried out the theoretical simulation. QH and MT performed the data analysis and wrote the manuscript.

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