

# Comparison of Intergroup Mass Transfer Coefficient Correlations in Two-Group IATE for Subcooled Boiling Flow

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The interfacial area concentration (IAC) is of vital importance in determining the interfacial transfer terms of mass, momentum, and energy between phases in the two-fluid model. The two-group (2-G) interfacial area transport equation (IATE) dynamically models the IAC evolution of large and small bubbles, and the intergroup transfer bookkeeps the mass transferring between two groups. Different intergroup mass transfer coefficient correlations are employed in the 2-G IATE model, which include one old correlation and three new correlations previously developed by the authors. The two-group IATE is benchmarked with subcooled boiling experimental dataset. The group-1 interfacial area concentration results show that the three modified correlations improve the physically incorrect prediction by the old correlation. With the modified correlations, the 2-G IATE is foundationally capable of predicting the group-1 interfacial area concentration in subcooled boiling flow.

Keywords: two-phase flow, boiling flow, IATE, two-group, intergroup transfer

# INTRODUCTION

The two-fluid model is now widely used to simulate two-phase flow systems like in RELAP5 (Thermal Hydraulics Group, 1998) and (TRACE (2010)). The model is formulated with two sets of conservation equations to separately consider the vapor phase and the liquid phase (Ishii and Hibiki, 2011) and the interface exists between the two phases. The interfacial area concentration,  $a_i$ , plays a key role in the interfacial drag closure and the interfacial heat transfer closures in the two-fluid model, as it determines the interfacial transfer terms of mass, momentum, and energy between phases (Khan et al., 2020) (Liu et al., 2021). **Figure 1** shows the diagram of the interface between the liquid and the vapor phase. The interfacial transfer terms can be written in the following form,

#### Interfacial transfer = $a_i \times \text{driving flux}$ (1)

Many researchers (Hibiki and Ishii, 2002) developed empirical interfacial area concentration correlations, including Ishii and Mishima (1980), Zeitoun et al. 1994, Zeitoun and Shoukri (1996), and Hibiki et al. (2006). The correlations can be divided into two kinds, namely geometry-based semi-empirical correlations and empirical correlations (Zhu et al., 2021). Different correlations are used in bubbly, slug, and the following regimes because they are

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dependent on the flow regime. As a result, numerical instabilities can occur at the transitions between the flow regimes. However, the interfacial area transport equation (IATE) is developed to dynamically model the interfacial area concentration across the flow regimes. The successful closure of interfacial area concentration in the two-fluid model relies on the interfacial area transport equation.

Kocamustafaogullari and Ishii (1995) developed the onegroup IATE based on the fluid particle number density transport equations analogous to Boltzmann's transport equation. The one-group IATE disregards the type of bubbles and treats them equally as spherical bubbles. Hibiki et al. (2003) provided the formulation of one-group IATE for subcooled boiling flow with the bubble layer thickness model. Brooks and Hibiki (2016) proposed a complete framework of the one-dimensional one-group (1-G) IATE model for phase change flow, then Zhu et al. (2021) further validate the 1-G IATE using extensive boiling flow datasets and improved the condensation closure in the model. At low void fractions, the one-group two-fluid model is sufficient for describing cases with uniform bubble size (Hibiki and Ishii, 2000a). However, for high void fraction cases, large bubbles in cap shape or slug shape come into existence as the consequence of bubble coalescence, bubble expansion, or heating from the wall. Under the high-void-fraction circumstance, the one-group IATE does not accurately predict  $a_i$  due to the sphericalshape and uniform-size assumption being applied to large bubbles (Zhu et al., 2019).

As the void fraction increases, the differences in shape, size, and number density between larger bubbles and small bubbles necessitate the separate consideration of bubbles in the twofluid model (Hibiki and Ishii, 2000b). In order to account for the different mechanisms for heat and mass transfer of bubbles with different shapes and sizes, the two-group (2-G) IATE model is



introduced in the work of Ishii and Kim (2004). Spherical and distorted bubbles are categorized as group-1 bubbles while cap and slug bubbles are categorized as group-2 bubbles in the two-group model (Hibiki and Ishii, 2000b). The separation of bubble groups thereby requires an additional set of conservation equations for the gas phase. Brooks et al. (2014) proposed the one-dimensional two-group IATE for subcooled boiling flow considering the uniform phase distribution pattern across the channel. Group-1 bubbles are averaged within a bubble layer while the group-2 bubbles span all the way across the channel.

Intergroup transfer happens between group-1 bubbles and group-2 bubbles: small bubbles can grow into large ones contributed from the coalescence, expansion, and nucleation, while the large bubbles can shrink into small ones as a result of the disintegration, condensation, and pressurization. Figure 2 shows the transfer between small and big bubbles. Intergroup transfer of the interfacial area concentration between groups is important for two-group IATE as it tracks and partitions of the gas phase into the bubble groups. This bookkeeping of bubble groups is facilitated by the intergroup mass transfer coefficient. The intergroup mass transfer correlation was first developed by Sun (2001). Zhu et al. (2020) provided a thorough review of the intergroup mass transfer coefficient and proposed three modified correlations for the term. Two analytical intergroup mass transfer coefficient correlations are also developed based on the bubble size distribution functions, including the Nukiyama-Tanasawa distribution Nukiyama and Tanasawa (1939) and the Rosin-Rammler distribution Rosin and Rammler (1933). Also, one empirical intergroup mass transfer coefficient correlation is obtained from the experimental data.

This paper focuses on the benchmark of the 2-G IATE with subcooled boiling experimental data using different intergroup mass transfer coefficient correlations. The modeling of two-group IATE and the various correlations of intergroup mass transfer coefficient are presented in **section 2**. The prediction results of interfacial area concentration and the sink/source terms are presented in **section 3**, accompanied by the discussion about the performance of the old and three modified correlations.

#### TABLE 1 | Correlations for the group-1 intergroup mass transfer coefficient.

Model	$\begin{aligned} \chi_1 &= 4.44 \times 10^{-3} \left( \frac{D_{am,1}}{D_c} \right)^{0.36} \alpha_1^{-1.35} \\ \chi_{1,NT} &= \frac{\frac{3}{4} D_c^{r_3} \exp(-2D_c^r)}{1 - (2D_c^{r_2} + 2D_c^{r_1} + 1) \exp(-2D_c^r)} \end{aligned}$	
Sun et al. (2004)		
Analytical correlation with Nukiyama-Tanasawa distribution (Zhu et al., 2020)		
Analytical correlation with Rosin-Rammler distribution (Zhu et al., 2020)	$\chi_{1,RR} = \frac{\frac{2}{2T} D_{c}^{-2} \exp[-\frac{(T_{c}^{-2})^{2}}{1-\exp(-(T_{c}^{-2})^{2})}]}{1-\exp(-(T_{c}^{-2})^{2})}$	
Empirical Correlation (Zhu et al., 2020)	$\chi_1 = \frac{b}{3} \frac{D_c^{sp+1} \exp[-bD_c^*]}{D_c^{sp}(bD_c^*)^{-p} (G(p+1)-G(p+1,bD_c^*))}$	
	where, $q = 1$ , $p = 0.25$ , $b = 0.780$ , $a = 0.80$	

Run#	Inlet pressure (kPa)	Heat flux (kW/m <sup>2</sup> )	Inlet liquid velocity (m/s)	Inlet subcooling (°C)	Max void fraction
2	273.0	90	1.25	5.95	0.52
3	248.7	56	0.68	6.47	0.55
4	489.9	90	1.32	5.12	0.31
5	900.8	90	1.33	4.04	0.24

## MODELING

Brooks et al. (2014) proposed the one-dimensional two-group IATE for subcooled boiling flow with the consideration of bubble layer thickness model; however, only intergroup mass transfer from group-1 bubbles to group-2 bubbles was modeled. Kumar and Brooks (2018) developed the model to include the intergroup mass transfer from group-2 to group-1. Therefore, incorporating the intergroup mass transfer from both group-1 and group-2 bubble sides, the change in group  $a_i$  in the one-dimensional, steady-state two-group IATE model for subcooled boiling flow is,

$$\begin{split} \Delta \langle a_{i,1} \rangle_B &= \frac{\Delta z}{\langle \langle v_{g,1} \rangle \rangle} \Big\{ \langle \phi_{WN,1} \rangle + \langle \phi_{Exp,1} \rangle_B + \sum_j \langle \phi_{j,1} \rangle_B - \langle \phi_{\Delta m_{12}} \rangle \\ &+ \langle \phi_{CO,1} \rangle_B + \langle \phi_{Conv,1} \rangle_B \Big\} \end{split}$$
(2)

and

$$\Delta \langle a_{i,2} \rangle = \frac{\Delta z}{\langle \langle v_{g,2} \rangle \rangle} \left\{ \langle \phi_{Exp,2} \rangle + \sum_{j} \langle \phi_{j,2} \rangle + \langle \phi_{\Delta \dot{m}_{12}} \rangle + \langle \phi_{Conv,2} \rangle \right\}$$
(3)

where subscripts g, 1, 2, and B denote vapor, group-1, group-2, and averaging over the bubble layer, respectively. The brackets <> and <<>> represent cross-sectional area average and void-weighted area average. The bracket <><sub>B</sub> represent the bubble layer average, which is calculated by the improved flat bubble layer thickness model in Zhu et al. (2022). The terms  $a_i$ , z, v, and  $\phi$  are interfacial area concentration, axial coordinate, velocity, and interfacial area concentration source/sink rate, respectively. There are nine source/sink terms in **Eqs 2**, **3**: the

wall nucleation term, the two expansion terms, the two bubble interaction terms, the two convection terms, the condensation term, and the intergroup mass transfer term. The wall nucleation term,  $\langle \phi_{WN,I} \rangle$ , and the expansion terms,  $\langle \phi_{EXP} \rangle$ , are modeled and described in Brooks et al. (2014). The bubble interaction terms,  $\langle \phi_{J} \rangle$ , and convection term,  $\langle \phi_{Conv} \rangle$ , are described in Kumar and Brooks (2018). The condensation term for subcooled boiling flow,  $\langle \phi_{CO,I} \rangle$ , is modified by Zhu et al. (2019). The intergroup transfer term,  $\langle \phi_{\Delta m 12} \rangle$ , is modeled by Kumar and Brooks (2018) and is written as,

$$\langle \phi_{\Delta \dot{m}_{12}} \rangle = \sum_{j} \langle \phi_{j,12} \rangle + \max(0, \langle \phi_{V,12} \rangle) + \min(0, \langle \phi_{V,21} \rangle)$$
(4)

where,

$$\langle \phi_{V,12} \rangle = \langle \phi_{\Gamma,12} \rangle + \langle \phi_{DP,12} \rangle, \qquad \langle \phi_{V,21} \rangle = \langle \phi_{\Gamma,21} \rangle + \langle \phi_{DP,21} \rangle$$
(5)

where,

$$\begin{split} \langle \phi_{\Gamma,12} \rangle &= \chi_1 \left( \frac{D_c}{\langle D_{sm,1} \rangle} \right)^2 \left( \frac{\langle a_{i,1} \rangle}{\langle \alpha_{g,1} \rangle} \right) \\ &\times \left( \frac{\langle \Gamma_{g,1} \rangle - \left( \langle \eta_{CO,1} \rangle + \langle \eta_{WN} \rangle \right) \langle \rho_g \rangle}{\langle \rho_g \rangle} \right) \quad (6) \\ \langle \phi_{DP,12} \rangle &= -\chi_1 \left( \frac{D_c}{\langle D_{sm,1} \rangle} \right)^2 \left( \frac{\langle a_{i,1} \rangle}{\langle \alpha_{g,1} \rangle} \right) \left( \frac{\langle \alpha_{g,1} \rangle}{\langle \rho_g \rangle} \right) \langle \langle v_{g,1} \rangle \rangle \frac{d \langle \rho_g \rangle}{dz}, \end{split}$$

$$\langle \phi_{\Gamma,21} \rangle = \chi_2 \left( \frac{D_c}{\langle D_{sm,2} \rangle} \right)^2 \left( \frac{\langle a_{i,2} \rangle}{\langle \alpha_{g,2} \rangle} \right) \left( \frac{\langle \Gamma_{g,2} \rangle}{\langle \rho_g \rangle} \right), \tag{8}$$



FIGURE 3 | The Run 1 results of group-1 intergroup mass transfer coefficient (A), the source of group-1 interfacial area concentration (B), and the interfacial area concentration (C) calculated by correlation Sun et al. (2004).



FIGURE 4 | The Run 1 results of group-1 intergroup mass transfer coefficient (A), the source of group-1 interfacial area concentration (B), and the interfacial area concentration (C) calculated with analytical-NT correlation (Zhu et al., 2020).









$$\langle \phi_{DP,21} \rangle = -\chi_2 \left( \frac{D_c}{\langle D_{sm,2} \rangle} \right)^2 \left( \frac{\langle a_{i,2} \rangle}{\langle \alpha_{g,2} \rangle} \right) \left( \frac{\langle \alpha_{g,2} \rangle}{\langle \rho_g \rangle} \right) \langle \langle \nu_{g,2} \rangle \rangle \frac{d \langle \rho_g \rangle}{dz}$$

$$\tag{9}$$

where the terms  $\alpha$ ,  $\rho$ ,  $\eta$ ,  $\chi$ ,  $D_{sm}$ ,  $D_c$  are void fraction, density, volume change rate per unit mixture volume, intergroup mass transfer coefficient, Sauter mean diameter, and critical bubble diameter at the group boundary, respectively. The subscripts V,  $\Gamma$ , DP, 12, and 21 denote volume change, interphase mass transfer, pressure change, group-1 to group-2, and group-2 to group-1, respectively.

In **Eqs 6, 8**, the interphase mass transfer rate per unit volume for group-n vapor,  $\Gamma_{g,n}$ , is a key term which cannot be measured directly and is hard to estimate in subcooled boiling flow with the theoretical models. However, this term can be determined based on the one-dimensional two-group mass balance equations (Ishii and Hibiki, 2011),

$$\frac{\partial \rho_g \langle \alpha_{g,1} \rangle_B \langle \langle \nu_{g,1} \rangle \rangle}{\partial z} = \Gamma_{g,1} - \Delta \dot{m}_{12}$$
(10)

$$\frac{\partial \rho_g \langle \alpha_{g,2} \rangle \langle \langle \nu_{g,2} \rangle \rangle}{\partial z} = \Gamma_{g,2} + \Delta \dot{m}_{12}$$
(11)

The intergroup mass transfer term,  $\Delta \dot{m}_{12}$ , is difficult to model in subcooled boiling flow, so it is assumed to be small relative to  $\Gamma_{g,1}$  in **Eq. 9**. The relative magnitude of  $\Delta \dot{m}_{12}$  to  $\Gamma_{g,2}$  in **Eq. 10** is inconsequential to the calculation of  $\langle \phi_{\Delta \dot{m}12} \rangle$  due to the suppression using the min function in **Eq. 4**. Therefore,  $\Delta \dot{m}_{12}$  is assumed to be negligible in **Eqs 9**, **10**. Even though in some instances (e.g., when  $\chi_1$  is nonzero)  $\Delta \dot{m}_{12}$  will not be negligible relative to  $\Gamma_{g,1}$ , the analysis uses the same  $\Gamma_{g,1}$  to isolate the effect of  $\chi_1$  on  $\langle \phi_{\Delta \dot{m}12} \rangle$ . More rigorous analysis to the two-group IATE can be applied once  $\Gamma_{g,n}$  is modeled correctly.

In Eqs. 6, 8, the intergroup mass transfer coefficient for group-n vapor,  $\chi_n$ , is another important parameter in the multigroup formulation and needs to be modeled accurately. The group-1 intergroup mass transfer coefficient,  $\chi_1$ , and the group-2 intergroup mass transfer coefficient,  $\chi_2$ , can be related analytically (Kumar and Brooks, 2018) as  $\chi_2 = (\chi_1 \cdot n_{b,1}) / n_{b,2}$  where  $n_{b,n}$  is the number density of groupn bubbles. The group-1 intergroup mass transfer coefficient was first developed by Sun et al. (2004), then Zhu et al. (2020) proposed three new correlations: an analytical correlation Nukiyama-Tanasawa distribution, an analytical with correlation with Rosin-Rammler distribution, and an empirical correlation. The correlations for the group-1 intergroup mass transfer coefficient are summarized in **Table 1.** In the equations,  $D_c^*$  is the dimensionless critical bubble diameter which is calculated  $D_c^* = 1.5 D_c / D_{sm,1}$ . G (z,x) and G(z) are the upper incomplete Gamma function and the complete Gamma function. In addition, a constant value of 0,

1, or 0.01 has been assigned to  $\chi_1$  in some works (Zhu et al., 2020).

The two analytical correlations are based on the distributions proposed by Nukivama and Tanasawa (1939), and by Rosin and Rammler (1933). Both are single-peaked, right-skewed distributions; but the height and the location of the peak are different with the same characteristic parameter, i.e., the bubble nondimensional diameter. The factors in the power-law and exponential function of the nondimensional diameter are different in the two analytical correlations. Consequentially, in the range of  $D_{\text{sm},1}/D_c \in [0.2, 0.8]$ , the intergroup mass transfer coefficient calculated by the analytical correlation with Nukiyama-Tanasawa distribution is larger than the correlation with Rosin-Rammler distribution. The empirical correlation is developed using an annulus condensing dataset (Kumar et al., 2019), an annulus flashing dataset (Kumar et al., 2019), an annulus boiling dataset (Bottini et al., 2020), and a pipe flashing dataset (Ooi et al., 2020). The empirical correlation predicts larger intergroup mass transfer coefficient values in the range of  $D_{\text{sm},1}/D_c \in [0.2,0.5]$ ; thus it better captures the data-points trend of the four datasets compared with the analytical correlations.

## **RESULTS AND DISCUSSION**

The subcooled boiling database by Bottini et al. (2020) is used to validate the two-group IATE model with four different group-1 intergroup mass transfer coefficient correlations: correlation Sun et al. (2004), analytical-NT correlation, analytical-RR correlation, and empirical correlation (Zhu et al., 2020). The experiment was conducted in an internally heated annulus channel with inner and outer diameters of 19.1 and 38.1 mm, respectively. The test section has a heated length of 2.85 m and an unheated length of 1.63 m. Four-sensor conductivity probes at three ports are used to measure interfacial area concentration and other two-phase parameters for group-1 and group-2 bubbles in the heated section. The conditions by Bottini et al. (2020) considered for the benchmark of two-group IATE model are listed in **Table 2**.

Figures 3–6 present the calculation results of Run 1 in the database of Bottini et al. (2020) with two-group IATE model using the correlation Sun et al. (2004), the analytical-NT correlation, the analytical-RR correlation, and the empirical correlation (Zhu et al., 2020). The subfigure A) in each figure is the calculated group-1 intergroup mass transfer coefficients. The  $\chi_1$  calculated by the correlation Sun et al. (2004) is observed to exceed the physical limit, 2, stated by Ishii and Hibiki (2011), in the lower region of the heated section where void fraction is small; thus, it is manually set to be two in the calculation. The cause of the nonphysical value is the power-law relation between  $\chi_1$  and  $\alpha_1$  in the correlation Sun et al. (2004), shown in the equation in **Table 1**. The  $\chi_1$ calculated by the correlationSun et al. (2004) is also observed to decrease when the void fraction increases along the heated length. On the contrary,  $\chi_1$  calculated by analytical-NT correlation, analytical-RR correlation, and empirical correlation (Zhu et al., 2020) increases along the heated length in the lower test section. This trend is more physically appropriate: the group-1 intergroup mass transfer coefficient should increase when more group-1 bubbles are present as  $\chi_1$  evaluates the amount of group-1 bubbles transferring to group-2 bubbles. Compared to the  $\chi_1$ calculated by the correlation Sun et al. (2004),  $\chi_1$  calculated by the other three correlations also has a smaller range, the maximum of which is ~0.18 for the three cases.

In the figures, the subfigure B) shows the source terms of group-1 interfacial area concentration: the intergroup mass transfer, wall nucleation, condensation, expansion, convection, and bubble interaction. The sum of the terms is shown with a solid blue line in the subfigures, and the intergroup mass transfer term is shown as a dashed black line. The subfigure C) is the comparison of simulated interfacial area concentration with experimental data for group-1 bubbles presented as red dots and red line. In the source term subfigures with the correlation Sun et al. (2004), the sum of the terms is negative or even infinite, dominated by a strong intergroup mass transfer term caused by the large  $\chi_1$ mentioned above. The intergroup term is so large on account of the unphysically large  $\chi_1$  term that any  $a_{i,1}$  generated by other means is immediately transferred to group-2, resulting in unphysical growth in group-2  $a_i$  while keeping  $a_{i,1}$  nearzero. This highlights the necessity of a new group-1 intergroup mass transfer coefficient correlation. With the analytical-NT correlation, and analytical-RR new correlation, the sums of terms calculated by 2-G IATE are positive values in the whole heated section. The calculated  $a_{i,1}$ is positive and increases along the heated length. The results show an improvement of the new analytical and empirical group-1 intergroup mass transfer coefficient correlations compared to the correlation Sun et al. (2004). In addition, compared with the  $\chi_1$  scale by analytical correlations, the constant value of 0 or 1 should not be assumed to  $\chi_1$ , which give under- and over-prediction of intergroup transfer term. The constant value of 0.01 should not be assigned to the whole heated section. It overpredicts in the low void fraction section and affects the section later. Thus, the correlations which can dynamically predict  $\chi_1$  are needed.

Figure 7 shows the group-1 interfacial area concentration benchmark of 2-G IATE with four group-1 intergroup mass transfer coefficient correlations, i.e., the correlation Sun et al. (2004), the analytical-NT correlation, the analytical-RR correlation, and the empirical correlation (Zhu et al., 2020). The results calculated by the correlation Sun et al. (2004) are not physically correct: the group-1 values of  $a_i$  are zero as explained above, while the group-2 part is the contributor to the  $a_i$  values here. However, when using the analytical-NT correlation, the analytical-RR correlation, and the empirical correlation, the 2-G IATE is foundationally capable to handle the transition to group-2 bubbles in subcooled boiling flow. A lot more works are still required to improve the two-group IATE in subcooled boiling flow, such as an accurate two-group interphase mass transfer model for  $\Gamma_{q,n}$ .

At high void fractions, the two-group (2-G) IATE model is developed to account for the different mechanisms for heat and mass transfer of bubbles with different shapes and sizes. Spherical and distorted bubbles are categorized as group-1 bubbles while cap and slug bubbles are categorized as group-2 bubbles. Intergroup transfer of the IAC between groups tracks and partitions the gas phase into the bubble groups, and the accurate modeling of the intergroup mass transfer coefficient is important for the two-group IATE.

In this paper, the 2-G IATE is benchmarked with subcooled boiling experimental data from Bottini et al. (2020). Different intergroup mass transfer coefficient correlations are employed in the 2-G IATE model, which include the correlation Sun et al. (2004), the analytical correlation with Nukiyama-Tanasawa distribution (Zhu et al., 2020), the analytical correlation with Rosin-Rammler distribution (Zhu et al., 2020), and the empirical correlation by Zhu et al. (2020). The group-1 interfacial area concentration results show that the three modified correlations improve the physically incorrect prediction by the correlation Sun et al. (2004). With the modified correlations, the 2-G IATE is

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foundationally capable of predicting the group-1 interfacial area concentration in subcooled boiling flow. More rigorous analysis to the two-group IATE can be applied once two-group  $\Gamma g$  is modeled correctly. The IATE model will be coupled with system codes in the next phase to present the numerical improvement across flow regimes.

## DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/Supplementary Material, further inquiries can be directed to the corresponding author.

### AUTHOR CONTRIBUTIONS

LoZ: conceptualization, investigation, methodology, validation, formal analysis, visualization, writing—original draft, project administration; JB: writing—review and editing; CB: methodology, writing—review and editing, supervision; LuZ: writing—review and editing.

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# NOMENCLATURE

 $a_i$  interfacial area concentration  $[m^{-1}]$ 

- D diameter [m]
- G gap width [m]

 $\Delta \dot{m}_{12}$  intergroup mass transfer rate per unit mixture volume [kg/m³/s]

- $n_b$  bubble number density  $[m^{-3}]$
- $\mathbf{Z}$  axial position [m]

# Greek

- $\alpha$  void fraction [-]
- $\varGamma$  vapor generation rate to the k-phase per unit volume  $\rm [kg/m^3/s]$
- $\Delta$  difference
- $\eta$  volume source rate  $[s^{^{-1}}]$
- ho density [kg/m<sup>3</sup>]
- $\phi$  Interfacial area concentration source or sink term [1/m/s]
- $\chi$  intergroup mass transfer coefficient

# Superscripts

- \* nondimensional
- ' boundary value between cap bubble and slug bubble

## **Subscripts**

- group-1
   from group-1 to group-2
   group-2
   from group2 to group-1
   c critical value at the boundary of group-1 and group-2 bubbles
   CO condensation
   Conv convection
   DP pressure change
   Exp expansion
   g vapor phase
- **i** interaction
- **sm** Sauter mean
- V volume change
- $\boldsymbol{W}\boldsymbol{N}$  wall nucleation

# **Brackets**

- $\langle \rangle$  area-averaged
- $\langle \rangle_B$  bubble layer averaged
- $\left<\left<\right>\right>$  void-weighted area-averaged