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# Numerical simulation of low reynolds number 2-d rough blade compressor cascade

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In this paper, a double-circular-arc cascade was used as the research object to study the effect of roughness on the compressor cascade under the conditions of different incoming attack angles. The Reynolds number of the cascade is  $1.5 \times 10^5$ . The profile of the suction surface is changed to simulate the change in surface roughness. A uniformly distributed 238 control points were selected on the suction surface profile of the smooth blade. The 238 points are increased randomly along with the normal phase of the wall, and a smooth curve connects these points to obtain a rough suction surface profile. In this study, numerical simulation methods were used to calculate the cascade flow field under 5 kinds of attack angles and 6 kinds of suction surface roughness (including smooth surface). In this paper,  $k - \omega$  SST and large eddy simulation were used for numerical simulation calculation with the model respectively. The calculation results show that the effect of roughness on the flow field of the cascade is different under different angles of attack. When  $70 < k_s^+ < 130$  the rough suction surface can increase the lift coefficient of the compressor blade.

#### KEYWORDS

roughness, boundary layer, compressor cascade, suction, large eddy simulation

#### 1 Introduction

Due to the different geographical locations and the social environment, the air will be mixed with solid or liquid particles such as dust, pollen, aerosols, etc (Bons 2010). The air in coastal areas will also contain salt spray particles (Mishra 2015). It is necessary to continuously consume a large amount of air to ensure the stable operation of the gas turbine. Inevitably, small-sized pollutant particles in the air will also enter the gas turbine. These solid or liquid particles will mix with the oil mist and adhere to the surface of the gas turbine components (Chen and Griffith, 2022b). When the gas turbine is running the compressor is the first part that comes into contact with the air, so the deposition of pollutant particles on the wall is also the most obvious. The deposition of solid or liquid particles in the air on the surface of the compressor components will eventually cause the wall to become rough. Especially when the surface of the compressor and gas turbine. Alessio Suman (Suman et al., 2017) et al. analyzed the distribution of particles on compressor blades through computational fluid dynamics (CFD).

The analysis has shown that particular fluid-dynamic phenomena and airfoil shape influence the pattern deposition. In the past, many scholars have conducted related studies on the influence of roughness on compressor performance (Chen et al., 2021; Chen and Griffith, 2022a). Mark P. Simens and Ayse G. Gungor used a flat plate with a strong inverse pressure gradient as the research object to study the effect of roughness on laminar separation bubbles. These numerical experiments indicate that laminar separation and turbulent transition are mainly affected by the type, height, and location of the roughness element (Simens and Gungor 2013). Jongwook Joo et al. (Joo et al., 2017). Regularly arranged rectangular blocks on the surface of the blade to simulate the rough surface, and used large eddy simulation to calculate the effect of roughness on the flow field of the NACA65 compressor cascade. Nicola Aldi et al. used NASA Stage 37 and NASA Stage 35 as research objects, simulated by imposing different spanwise distributions of surface roughness to analyze its effect on compressor performance (Nicola et al., 2013; Sun et al., 2018). Ravi J Chotalia and Dilipkumar Bhanudasji Alone (Chotalia and Alone 2017) studied the performance degradation of the single-stage transonic axial flow compressor due to the uniform roughness created by the aerosols. Yonghong Tang et al. found that doubling or halving the roughness value of all components of a centrifugal compressor in the hydraulically rough flow regime will result in a change in the peak efficiency of the compressor by about 1.72% (Tang et al., 2019). A test program on the GE J85-13 jet engine quantified the increased surface roughness and the distribution of salt deposits in an axial compressor (Syverud and Bakken 2006). Dongjae Kong (Kong, Jeong, and Song 2017) conducted experimental measurements on the smooth and rough compressor cascade passage flow field.

In recent years, compressor design technology has become more and more advanced, and compressor performance parameters have gradually improved. However, it can be noticed that the compressor blades are now more twisted than before. During the operation of the compressor, the gas flow around the blades is more complicated. It becomes complicated to analyze the influence of rough wall surface on the flow field around the compressor blade. Cascade passages are often used to analyze the two-dimensional flow process of compressor blades. The blades are straight along the span, so the transmission movement of the surrounding airflow along the blade height direction disappears. Therefore, this study used double-curved cascades as the research object and uses a combination of steady-state numerical simulation and Large-eddy simulation (LES) to study the effect of blade surface roughness on the performance of compressor cascades.

## 2 Model development

#### 2.1 Geometric model

#### 2.1.1 Smooth cascade blade

A cascade was used to simulate the effect of roughness on the compressor. The blade profile is a compressor blade designed at



TABLE 1 Blade geometry.

228.6 mm	$\Sigma$	2.14
106.8 mm	AR	1.61
12.5 mm	Г	20.5°
914.4 μm	$\kappa_{I}$	53.0°
914.4 μm	$\kappa_2$	-12.0°
	Φ	65.0°
	106.8 mm 12.5 mm 914.4 μm	106.8 mm     AR       12.5 mm     Γ       914.4 μm     κ1       914.4 μm     κ2

the NASA Lewis Research Center. The blade section of the cascade is a double-circular-arc blade with 65 degrees of camber, a 20.5° stagger angle, a solidity of 2.14, and a 228.6 mm chord length. Both the radius of the leading and trailing edge is 0.9144 mm (Zierke and Deutsch 1990). The profile of the blade is shown in 1.1.11Figure 1. After the profile of the rough blade is generated, analyze the height change value of the control point on the suction surface to obtain the specific value of the roughness on the suction surface.

The equation for the pressure surface, suction surface, and camber line of the double-circular-arc blade used in the current study can be written as

$$x_p^2 + \left(y_p + 219.7\right)^2 = 246.8^2 \tag{1}$$

$$x_s^2 + (y_s + 149.5)^2 = 189.1^2$$
<sup>(2)</sup>

and

$$x_c^2 + (y_c + 179.4)^2 = 212.8^2$$
 (3)

The chord length of this cascade blade is 228.6 mm, and the radius of the leading edge and trailing edge are both 0.9144 mm.



The aspect ratio (AR) is 1.61. The geometric parameters of the blade profile are shown in Table 1.

### 3 Rough cascade blade

In past research, researchers often used two methods to study the effect of roughness on the compressor cascade. Mirko Morini et al. (Morini et al., 2011; Nicola et al., 2014; Shi et al., 2017) used the wall equation to simulate the roughness on the compressor blade, while Sun Hai-Ou et al. (Sun et al., 2021) changed the profile of the blade.

Previous research (Syverud and Bakken 2006) data has shown that the roughness on the suction surface of the compressor blade has a greater impact on performance. So only the suction side of the blade profile was been altered in this study. The rough blade profile is shown in Figure 2.

In this research, the method of replacing the blade profile was used to simulate the rough wall surface, to describe the rough wall surface more realistically. 238 control points were selected on the suction surface profile of the smooth blade, and each control point was moved at a random distance in the direction perpendicular to the wall. A spline curve is formed by the newly obtained 238 control points. The shape of this curve is regarded as the profile of the rough blade, as shown in Figure 2.

#### 3.1 Numerical mesh

Irregular contours on rough walls cause inconvenience to mesh generation. The quality and structure of the grid determine the accuracy and convergence of the numerical simulation, especially for large-eddy simulations. As shown in Figure 3, O-Block was used around the blade surface and the number of the mesh is 160 million. The *y*<sup>+</sup> value is less than 1, 1500 nodes was set at chord direction on the blade. The number of mesh nodes in the boundary layer of the blade surface is more than 40.  $\Delta x^+$  is one of the parameters for evaluating the quality of the grid, which is defined as follows.

$$\Delta x^{+} = \frac{u_{\tau} \cdot \Delta x}{v} = \frac{y^{+}}{y} \cdot \Delta x \tag{4}$$

The range  $\Delta x^+$  of the grid in this study is  $1 < \Delta x^+ < 10$ .

### 4 Numerical simulation

The influence of the angle of attack of the cascade on the flow field around the blade is also considered.  $k_s^+$  is a dimensionless parameter representing roughness, which is defined as follows:

$$k_s^+ = R_e \frac{k_s}{c} \sqrt{\frac{C_f}{2}}$$
(5)

$$C_f = \left[2.87 + 1.58 \lg \frac{c}{k_s}\right]^{-2.5}$$
(6)

The geometric mean roughness Ra satisfies the relationship of  $k_s = 6.2R_a$  between  $k_s$ . Where  $C_f$  is the wall shear stress coefficient, and c is the chord length of the blade. The Re in the formula is the Reynolds number using the chord as the characteristic length. Schlichting's (Schlichting and Gersten 1965) research shows that when  $k_s^+ < 5$ , the roughness has almost no effect on the flow loss, and the surface is Hydraulically smooth; when  $5 < k_s^+ < 70$ , the airflow loss is determined by Re and  $k_s^+$  together, which is the Transitional rough; when  $k_s^+ > 70$ , the loss of airflow is no longer affected by Re and is only related to  $k_s^+$ , which is fully rough. In this study, a total of 6 types of roughness blades (including smooth blades) were numerically simulated, and 5 different incoming attack angles (*i*) were included, as shown in Table 2.

Figure 4 shows the computational domain model of the cascade. The single-passage numerical simulation requires less computer and can increase the calculation speed of numerical simulation. The boundary conditions of the calculation model are as follows:

The Inlet Reynolds number is  $1.5 \times 10^5$ , and the total temperature of Inlet is 288.15 K. The incoming flow contained a freestream turbulence level of 10%, which is the same with the measurement. The cascade blade was set to smooth, adiabatic and non-slip wall condition, static pressure outlet condition was used on the outlet surface. The inlet Reynolds number is based on the blade chord.

In the steady-state calculation, the  $k - \omega$  SST turbulence model was selected as the turbulence model, and the  $\gamma - \theta$ transition model was used. LES was used to calculate the unsteady state of airflow. The principal idea behind LES is to reduce the computational cost by ignoring the smallest length scales, which are the most computationally expensive to resolve, via low-pass filtering of the Navier–Stokes equations. It is advantageous in solving the gas flow state near the wall.



TABLE 2 Incoming attack angle (i) and wall roughness ( $k_s^+$ ).

$k_s^+i$	0.00	11.80	65.05	105.94	141.21	196.47
$0^{\circ}$	S-0	R1-0	R2-0	R3-0	R4-0	R5-0
-5°	S-5	R1-5	R2-5	R3-5	R4-5	R5-5
-8°	S-8	R1-8	R2-8	R3-8	R4-8	R5-8
+5°	S+5	R1+5	R2+5	R3+5	R4+5	R5+5
$+8^{\circ}$	S+8	R1+8	R2+8	R3+8	R4+8	R5+8



#### 5 Simulation result

At different angles of attack, the  $C_p$  distribution on the airfoil surface is shown in Figure 5, the  $C_p$  is defined as follows:

$$C_p = \frac{P - P_{in}}{P_{in}^* - P_{in}} \tag{7}$$

It can be seen from Figure 5 that under different surface roughness, the distribution  $C_p$  varies with the angle of attack almost the same. Compared with the case of 0° angle of attack, at a negative angle of attack, the pressure coefficient of the blade surface decreases. But the trend shows the difference when the angle of attack is positive. When  $i = +5^\circ$  the  $C_p$  is much higher

than the case of  $i = 0^{\circ}$  before 0.7 chord, but lower after 0.7 chord. However, this position became 0.35 chord when  $i = +8^{\circ}$ . Notice the general fluctuation of the pressure coefficient curve on the suction side of the rough blade. When reading the wall pressure data, directly select the suction surface profile as the target object. The irregularity of the rough blade profile causes data fluctuations on the curve.

Under the same angle of attack, when the roughness of the suction surface is different, the distribution of the pressure coefficient is shown in Figure 6. The simulation results show that when the *i* is positive, the  $C_p$  hardly changes with the roughness of the suction surface.

When *i* is less than or equal to 0, the  $C_p$  changes with the roughness value. When  $k_s^+$  is greater than 105.94, the pressure coefficient of the rough suction surface is smaller than that of the smooth blade, and the pressure coefficient continues to decrease as the roughness increases. When  $k_s^+$  is less than 105.94, the pressure coefficient of the rough blade surface is greater than that of the smooth blade. With the increase  $k_s^+$ , the pressure coefficient of the rough blade gradually approaches the pressure coefficient of the smooth blade.

So for the blades of the cascade, the lift coefficient is the area value enclosed by the pressure coefficient curve in Figure 5 and Figure 6. The abscissa of Figure 7 is the lift coefficient of the blade, and the ordinate is the angle of attack and the equivalent sand roughness on the suction surface of the cascade blade, the lift coefficient is defined as follows:

$$c_L = \frac{L}{\frac{1}{2}\rho v_{\infty}^2 l} \tag{8}$$

Where *L* is the lift,  $\rho$  is the density,  $\nu_{\infty}$  is the far-field velocity, and *l* is the chord length. For the cascade blade  $L = L_p - L_s$ , and when *l* is the dimensionless chord length, the range is from 0 to 1. Where the density is considered to be constant.









The pressure coefficient is along the y-axis direction, as shown in Figure 1. The effect of the roughness of the suction surface on the lift coefficient of the blade is not completely negative. When  $70 < k_s^+ < 130$  the rough suction surface can increase the lift coefficient of the compressor blade. However, this phenomenon only occurs when  $i < 0^\circ$ . When the angle of attack of the cascade is positive, the roughness on the suction surface of the blade no longer affects the lift coefficient. Figure 7 shows that a suitable negative angle of attack has a positive effect on the lift coefficient of the compressor cascade blade.

The premature separation of the airflow on the lowroughness wall destroys the original gas flow state near the wall and also causes the thickness of the boundary layer near the trailing edge of the blade to be significantly reduced.

As the roughness of the blade surface increases, the separation bubbles disappear. The boundary layer separation gradually thickens, and the position of the airflow separation is advanced. These phenomena can be observed in Figure 9.

The isentropic Mach number on the suction surface of the blade is shown in Figure 8. Starting from 0.7 chord length, the isentropic Mach value begins to change due to the increase in roughness. In addition, when  $k_s^+ = 11.80$  and  $k_s^+ = 65.05$ , separation bubbles appear at 0.2 chord length. When the wall surface is smooth, the airflow maintains a laminar state before 0.7 chord length, after which the airflow gradually changes to a turbulent state after passing through the transition zone. However, the wall becomes rough, the laminar flow is disturbed, and becomes turbulent flow at 0.2 chord length. It can be found in Figure 9 that the separation bubble disappears at 0.4 chord length, and the airflow near the wall becomes laminar again. After 0.4 chord length, the airflow into the wall has been maintained in a laminar state until the 0.8 chord length position. Figure 10 shows the velocity distribution near the suction surface of the cascade. The abscissa is the position of the chord length, and the ordinate is the distance from the suction surface.

In previous studies, the inlet velocity of the cascade was often used as the far-field velocity when calculating the airflow displacement thickness on the suction surface. Figure 10 shows the velocity distribution in the cascade passage. It can be easily found from the figure that the velocity of the airflow is different at different chord length positions. So our team used different far-field velocities at different position when calculating the displacement thickness.

Figure 11 shows the performance of the displacement thickness on the suction surface under different roughness. Since the incoming Mach number is very small, the density is assumed to be constant when calculating the displacement thickness. Overall, the displacement thickness increases with roughness. However, when the roughness is small ( $k_s^+$  = 11.80), the displacement thickness of the blade back surface begins to increase at x/c = 0.2, which is greater than the displacement thickness of the smooth blade cascade. But when x/c = 0.4, the displacement thickness value of the rough blade is equal to the case of the smooth wall, and is smaller than that of the smooth wall after that.

In this paper,  $x_d$  represents the value of x when it is equal to the displacement thickness of the smooth blade. The position of  $x_d$  lags with increasing roughness. When the roughness is 11.80, the critical position is at 0.46 chord, and when the roughness is 65.05, the critical position is at 0.6 chord.

The shape factor ( $H_{12}$ ) of the suction surface of the blade is shown in the Figure 12, the shape factor of the smooth blade is almost kept around 2, and slightly increased near the leading and trailing edges. When the roughness of the  $k_s^+$  = 11.80, the shape





factor value increases from the leading edge of the blade and suddenly increases after the chord length of 0.11, where the airflow starts to transition. From the shape factor, it can be found that the appearance of roughness induces an early transition of the boundary layer on the back surface of the leaf. But when the roughness is 11.80, the airflow on the back surface of the blade returns to the laminar state again at the chord length of 0.75. After the boundary layer is regenerated, the shape factor of rough leaves is also slightly reduced compared to smooth blade.

Figure 13 shows the outlet airflow angle of the cascade. The outlet angle of the cascade varies with the roughness of the suction surface and is affected by the angle of attack of the incoming flow. When the angle of attack of the incoming flow is negative, the airflow angle of the outlet increases with the increase of roughness. When the angle of attack of the incoming flow is positive, even if the suction surface of the cascade becomes rough, the airflow angle at the outlet will hardly change. Once the angle of attack of the incoming flow is greater than 0, the increase of the reverse pressure gradient when the flow passes through the cascade passage will cause the airflow to





transition in advance. The rough wall surface on the suction surface will be completely covered by the boundary layer with increased thickness. Rough suction surface no longer has an interference effect on the mainstream area.

The velocity distribution along the x direction was measured at position 1. Position 1 is at 150% chord length from the leading edge of the blade. The effect of roughness on the velocity distribution is linear. As the roughness increases, the minimum velocity decreases linearly, as shown in Figure 14. The position of the minimum velocity is biased towards the suction surface of the blade.





Take the speed extreme point of the smooth blade as the zero point, and the extreme point position is positive toward the pressure surface and negative toward the suction surface. When  $k_s^+ \leq 105.94$ , the rough suction surface has a positive influence on the speed extreme point, and when  $k_s^+ > 105.94$ , the rough wall has a negative influence on the speed extreme point position.



Figure 15 is the contour of the velocity extreme point position and the velocity minimal value. When i is positive, the roughness does not affect the velocity distribution on the B-to-B plane. Meanwhile, as the roughness increases, the velocity of the blade exit is less sensitive to the angle of attack of the incoming flow. As the roughness of the suction surface increases, the velocity gradient at the exit of the cascade increases. The mixing loss of the secondary flow of the blade wake and the main flow increases due to the increase of the velocity gradient. The increase in roughness will also

reduce the sensitivity of the cascade flow field to the angle of attack of the incoming flow.

The calculation result of LES is shown in Figure 16 Figure 17to Figure 18. Vorticity in the cascade passage is shown in the figure. In continuum mechanics, vorticity is a pseudovector field that describes the local spinning motion of a continuum near some point (the tendency of something to rotate), as would be seen by an observer located at that point and traveling along with the flow. It is an important quantity in the dynamical theory of fluids and provides a convenient framework for understanding a variety of complex flow phenomena, such as the formation and motion of vortex rings (Guyon et al., 2002; Higham et al., 2015).

For a smooth cascade passage, the vorticity on the suction surface starts to change at 50% chord length. At the beginning of 0.36 chord length of the suction surface, the high-value area of the vorticity begins to become concentrated and gathers into a cluster. This phenomenon indicates that the airflow changes from laminar flow to turbulent flow from this position. The vortex first moves tightly against the suction surface force towards the trailing edge of the blade, and its size is gradually increasing. At 0.8 chord length, the vortex starts to break away from the suction surface. Unlike the suction surface, the airflow on the pressure surface does not turn into turbulence until it passes the trailing edge, and a larger vortex V2 is formed. The vortex V1 on the suction surface collides with the vortex V2 on the pressure surface downstream of the blade.

The vortex V2 is impacted by the high-speed, small-scale vortex V1 at the trailing edge of the blade. V2 is crushed and divided into a large-scale vortex and multiple small-scale vortices. The large-scale vortices continue to move downstream along the main flow direction, and the small-







scale vortices are gradually consumed during the movement and almost disappear at T6. The shape of V1 changes from the previous round shape to a long strip shape at T4, and it returns to a round shape at T5. At t4, V1 is deformed under the influence of the vortex generated by the trailing edge of the pressure surface. However, in a short period, V1 regained its round shape under the influence of the mainstream, and the scale of the vortex was not affected.

When  $k_s^+ = 11.80$ , the airflow transitions at 0.36 chord length of the suction surface. The mountain-like profile on the rough suction surface interferes with the airflow near the wall. The airflow that has just been transformed into turbulence is disturbed by the profile, and the size of the vortex is reduced. The vortex disappears after the

turbulent flow of 5% chord length, and the airflow becomes laminar again. The pressure gradient along the flow direction causes the airflow to turn turbulence again at 0.6 chord. Different from the smooth surface, the vortex has not been away from the wall before the tail edge of the blade. In the flow direction of the airflow, the scale of the eddy is gradually increased. Furthermore, when  $k_s^+ = 11.80$  compared to the smooth wall surface, the scale of the vortex is reduced.

Figure 16 is the vorticity amount of the passage when  $k_s^+$  = 196.47. Referring to Schichting's research conclusion, the suction surface is fully rough now. When the suction surface is completely rough, the scroll is significantly increased, and the movement of the airflow has become

more complex. The transfer of airflow occurs at 0.33 chord of the suction surface, which is the same as  $k_s^+ = 11.80$ . However, the phenomenon of the airflow is retrogradely laminar flow disappears when the wall is Fully Rough. Due to the increase in roughness, the disturbance of the wall faces the airflow is more pronounced. The large-scale vortex near the wall is broken by the influence of the mountain-like rough profile. The vortex is broken into the main vortex and multiple small-scale vortices and the small-scale vortex is sandwiched between the wall and the main vortex. The main vortex is squeezed away and gradually moves away from the wall.

## 6 Conclusion

In this study, we simulated the effect of a 2D-rough surface on a double-circular-arc blade cascade. The profile of the suction surface was replaced to simulate a rough surface. From the above discussion, one may conclude that the sensitivity of the pressure coefficient distribution on the blade surface to the angle of attack does not correlate with the roughness at the suction of the blade. When the incoming flow angle of attack is positive, the velocity distribution of the trailing edge of the cascade along the b-b direction and the airflow angle at the outlet will not change with the roughness. When  $70 < k_s^+ < 130$  the rough suction surface can increase the lift coefficient of the compressor blade. However, this phenomenon only occurs when  $i < 0^\circ$ .

Unfortunately, only the effect of two-dimensional roughness on the airflow in the cascade passage was simulated in this study. The impact of the three-dimensional roughness of the blade surface on the flow of the cascade is more complicated. Meanwhile, the numerical calculation of the three-dimensional rough wall also puts forward a greater demand on the computer, which is currently difficult to supply the requirements. So one important future direction of studying the effect of the rough blade on the compressor cascade is establishing the relationship between the two-dimensional flow and the three-dimensional flow in the rough cascade. And more experiments should be conducted for the compressor cascade.

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### Data availability statement

The original contributions presented in the study are included in the article/supplementary material, further inquiries can be directed to the corresponding author.

## Author contributions

L-SW is mainly responsible for rough blade calculation and data collation in this research Z-YW is responsible for the guidance and assistance of numerical simulation calculations in this research Y-HW is responsible for providing guidance and assistance in the writing of the article in this research MW is responsible for providing advice and guidance on the research content in this research H-OS is responsible for revising and optimizing the article during the research.

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# Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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# Nomenclature

k<sub>s</sub><sup>+</sup> Dimensionless roughness
Re Reynolds number
Ra Surface Roughness
AR Aspect ratio
C<sub>f</sub> Wall shear stress coefficient

 $k_s$  Equivalent sand diameter

c Chord

i Incoming attack angle

 $\boldsymbol{c}_p$  Pressure coefficient

 $c_L$  Lift coefficient

Ma Mach number

V Velocity

 $\delta_1$  Displacement thickness  $H_{12}$  Shape factor  $\beta$  Exit angle

# Subscripts and superscripts

p Pressure side
s Suction side
LE Leading-edge
TE Trailing-edge
in Inlet
out Outlet
is Isentropic
\* Total