



# Influence of Defects to $Zr_{65}Cu_{18}Ni_7Al_{10}$ Bulk Metallic Glass Properties Under Dynamic Compression

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The  $Zr_{65}Cu_{18}Ni_7Al_{10}$  bulk metallic glass with smaller diameter exhibits higher fracture strength under dynamic compression, which is ascribed to concentration of flow defect. The density of shear bands in the sample surface will increase with decreasing of the diameter, whereas, average distance and width of tear ridges in the fracture surface will increase with larger diameter. In addition, the volume of shear transformation zone can be estimated, which presents a ductile-to-brittle transition with the change of diameter. The physical graph of shear transformation zone can be obtained from the experimental analysis.

**Keywords:** metallic glass, mechanical testing, high-speed deformation, fractography, shear transformation zones

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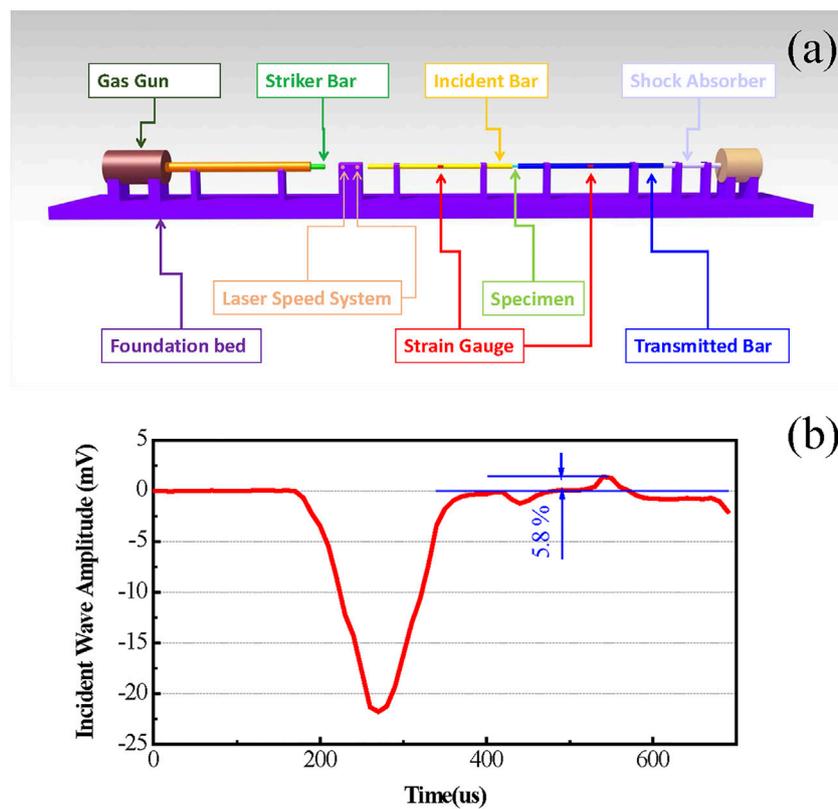
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## INTRODUCTION

In view of its amorphous nature of the long-range disordered feature, bulk metallic glasses (BMGs) exhibit high strength and hardness, large elastic limit, high fracture toughness and excellent corrosion resistance compared with their crystalline counterparts (Johnson, 1999; Schuh et al., 2007; Qiao and Pelletier, 2014; Xie et al., 2019; Tao et al., 2021). However, studies about mechanical properties concluded that the vast majority of BMGs exhibit poor plastic deformation at room temperature (Schuh and Nieh, 2003; Wang, 2012; Zhang et al., 2016a). Recently, a typical ductile-to-brittle transition was found in BMGs with high applied strain rates (Jiang et al., 2008; Li et al., 2016). Different deformation behaviors of BMGs upon quasi-static and dynamic compressions at different strain rates and temperatures have been investigated (Sunny et al., 2008; Liu and Liu, 2011; Huang et al., 2012), these researches showed either positive or negative correlation between fracture stress and strain rate. In addition, the size effect of BMGs have been extensively studied by experimental analysis and molecular dynamics simulation under quasi-static biaxial loading condition (Huang et al., 2007; Wu et al., 2015; Zhang et al., 2016b; Wang et al., 2016; Yang et al., 2016). With a faster cooling rate, smaller BMG containing more free volume or shear transformation zones (STZs) will be conducive to the formation of shear bands and further enhance its plasticity under quasi-static compressions (Huang et al., 2007; Zhang et al., 2016b). Up till now, size effect on the dynamic mechanical properties are rarely reported. Most investigations, if any, focus on the experimental phenomena. However, the physical mechanisms remain unclear. The fundamental mechanism of size effect on dynamic mechanical properties still requires further study.

It is well known that there are no dislocations or grain boundaries in BMGs. The frozen-in “defects” of BMGs, which can be depicted as local fluctuations of density, is understood as the source of macroscopic plasticity and the origin of structural relaxation (Castellero et al., 2008). The “defects” can be described by physical models under different circumstances, such as STZs (Argon, 1979), free



**FIGURE 1 | (A)** Schematic diagram of the SHPB device for dynamic compression experience; **(B)** a suitable incident wave for loading BMGs samples to guarantee the reliability of data.

volume (Spaepen, 1977; Qiao et al., 2015), liquid-like core (Huo et al., 2013; Li et al., 2013), flow units (Wang et al., 2014, 2014), quasi-point defects (Perez, 1990), coupling model (Ngai, 1998), or weakly bonded regions (Ichitsubo et al., 2005). In the current research, the conception of STZ will be used. As the review depicted (Wang et al., 2015), STZ can be pictured as an “event” or a consequence of plastic flow (Argon, 1979), and also can be regarded as a flow defect which triggers plastic flow (Langer and Lemaitre, 2005). Herein, both of these two definitions are used.

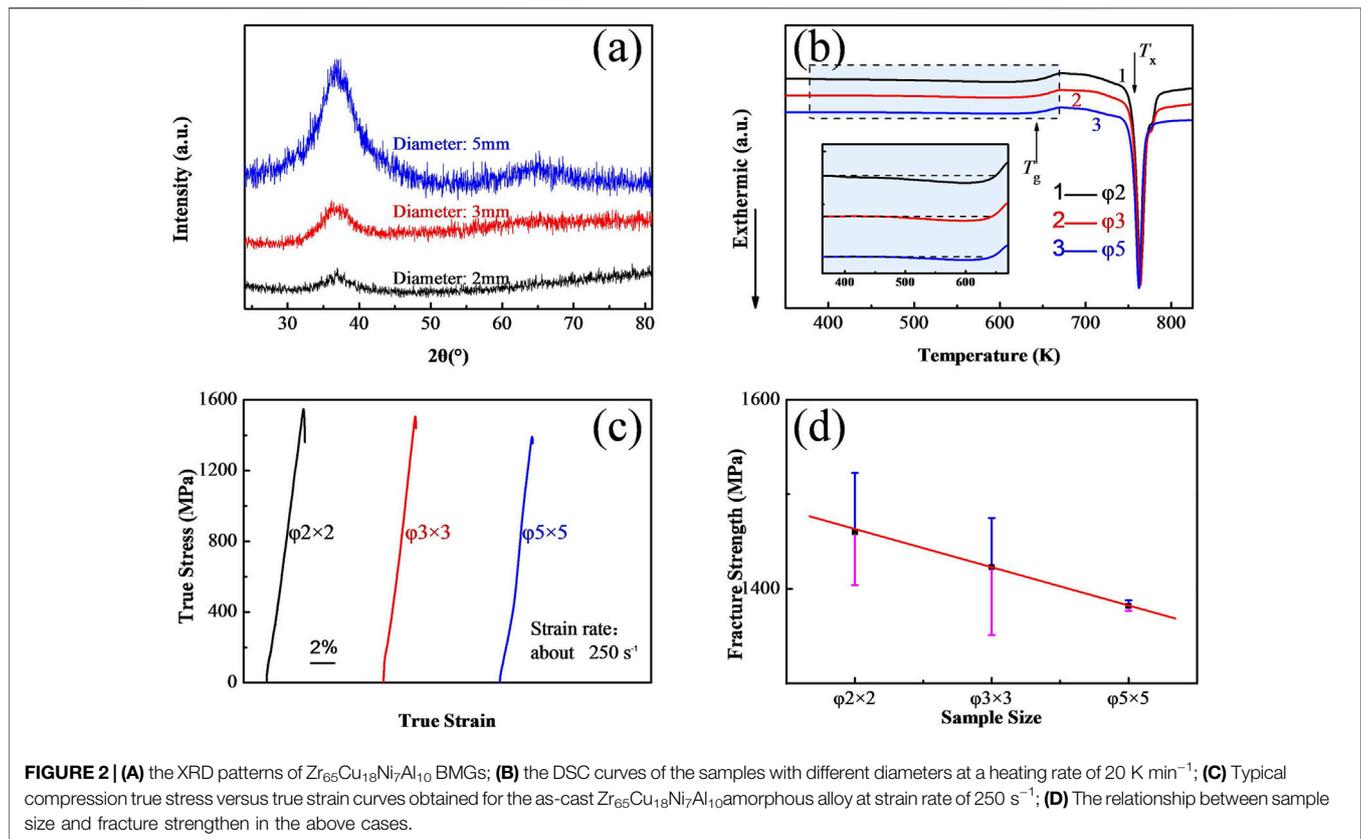
In this study, the size effect under high strain rate compression was investigated based on microstructure because mechanical properties are often determined by microstructure. By studying thermal and dynamic compressive properties of BMGs with the diameter of 2, 3, and 5 mm, the physical landscape between intrinsic microstructure (the number of STZs per unit volume) and dynamic compressive properties could be obtained.

## EXPERIMENTAL PROCEDURES

Because the alloy show superior plastic processing capacity, i.e. the supercooled liquid region is  $\sim 120$  K (Wang et al., 2015; Zhang et al., 2016a), the master alloys of  $Zr_{65}Cu_{18}Ni_7Al_{10}$  (atom %) were re-melted at least five times to guaranty its chemical homogeneity. Rod-like BMGs with different diameters were

then fabricated by employing the copper mould casting method. The amorphous nature of as-cast specimens was examined by the X-ray diffraction (XRD, Philips PW3830). The thermal properties and relaxation behavior of  $Zr_{65}Cu_{18}Ni_7Al_{10}$  BMG were determined by differential scanning calorimetry (DSC, NETZSCH DSC 200F3). Cylindrical samples with different sizes of  $\phi 2 \times 2$  mm,  $\phi 3 \times 3$  mm,  $\phi 5 \times 5$  mm were prepared and both ends were polished carefully to ensure dustless and parallelism for the dynamic mechanical tests (Zhang et al., 2016a; Zhang et al., 2016b). The height error of each sample is less than 0.05 mm. The dynamic compressive experiments were performed on a split Hopkinson pressure bar (SHPB) apparatus at an average strain rates of  $250 \text{ s}^{-1}$ . The sample surface and fracture morphology of the BMGs samples were observed by the scanning electron microscopy (SEM, TESCAN MIRA 3 XMU).

Dynamical compression experiments were performed by SHPB technique at high strain rates, as illustrated in Figure 1A. Specimens with an aspect ratio of one were sandwiched between the incident bar and the transmitted bar. Taking the higher strength of BMGs into account, two protection plates (12.7 mm in diameter) made of tungsten core and titanium-alloy ring were attached to the ends of bars between the samples. The size of tungsten core (6 mm in diameter) was designed in order to match the wave impedance of the steel bar.



Copper impulse shaper was placed between the striker bar and the incident bar to shape the incident wave, aiming at maintaining a constant-strain rate loading during the dynamical compression experiment. The specimen ends were lubricated to minimize friction. We examined the effectiveness of impulse shaping and protection plates by compression test. **Figure 1B** shows a suitable incident wave for loading BMGs samples. In addition, the low amplitude of the reflected wave implies that introduction of protection plates has negligible effect on the experimental results. Thus, relatively reliable experimental results could be obtained, which formed the basis of current analysis. Eight test pieces were prepared for each test and at least four valid data were obtained.

## RESULTS AND DISCUSSION

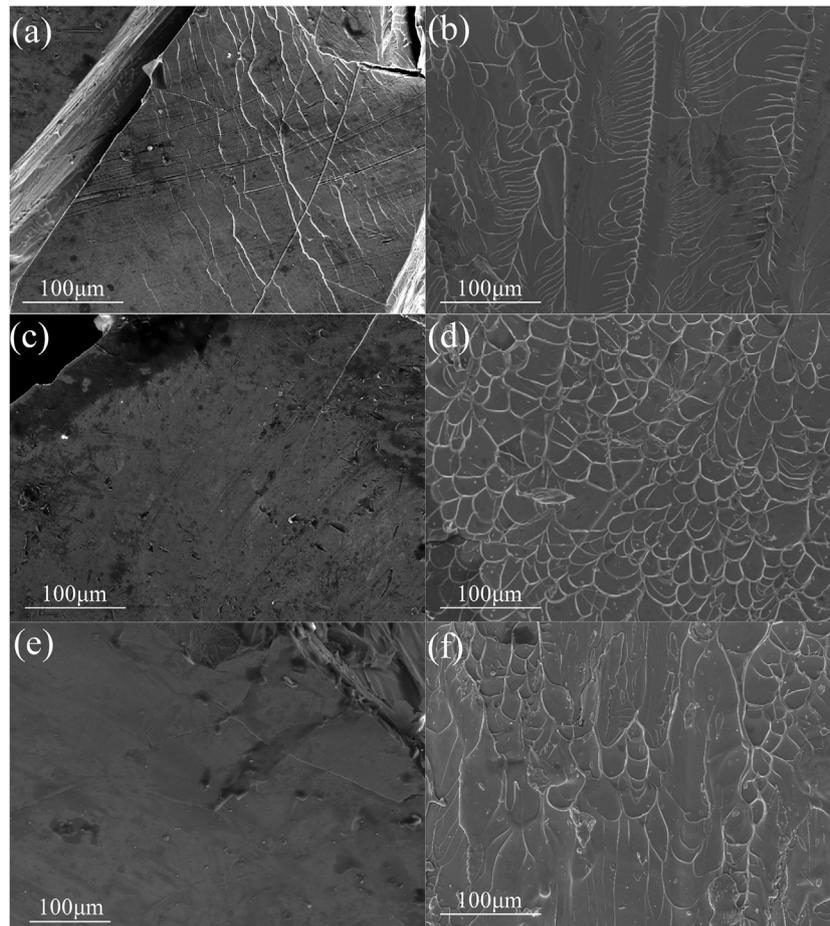
**Figure 2A** shows the XRD patterns of  $Zr_{65}Cu_{18}Ni_7Al_{10}$  BMGs with different diameters. The typical diffraction humps were exhibited, which suggest amorphous structure of the alloys. The bigger the diameter, the stronger the intensity owing to the length of deflection.

Typical DSC curves were conducted at heating rate of  $20\text{ K min}^{-1}$ , as shown in **Figure 2B**. Each DSC curve exhibits a single endothermic reaction, which corresponds to the glass transition process, which followed by a significant exothermic event corresponding to a crystallization behavior in the

**TABLE 1 |** Thermal properties of samples measured from their DSC curves and dynamic compression performances obtained by the SHPB tests.

Parameters	Diameter (mm)		
	2	3	5
$T_g$ (K)		643	
$T_x$ (K)		757	
Structural relaxation exothermic heat, $\Delta H$ (J/g)	0.201	0.100	0.067
Relative STZ concentration $V_{STZ}/\Delta V^*$	1/6.9	1/13.9	1/20.7
Fracture strength (MPa)	1547.8	1506.1	1391.3
Average distances of tear ridges ( $\mu\text{m}$ )	7.048	17.776	27.040
Average width of tear ridges ( $\mu\text{m}$ )	1.561	2.541	3.816

supercooled liquid region. As displayed in **Table 1**, the glass transition temperature  $T_g = 643\text{ K}$  and onset temperature of crystallization  $T_x = 757\text{ K}$  were acquired from the DSC line chart, respectively. These experimental results agree well with the previously published data (Zhang et al., 2016a; Zhang et al., 2017). It is noted that smaller sample with the same composition contains larger structural relaxation exothermic heat, because of the faster cooling rate (**Table 1**). As discussed in the previous literature (Zhang et al., 2016b), structural relaxation exothermic heat can be considered as characterize “flow defect” in BMGs, called STZs (Langer and Lemaître, 2005), which play an essential role on the mechanical properties (Huang et al., 2007; Zhang et al., 2016b). As shown in **Table 1**, structural relaxation exothermic heat (involved energy per unit mass) with the



**FIGURE 3** | SEM images of the samples surfaces and fracture surfaces with diameters of 2 mm [(A,B)], 3 mm [(C,D)], and 5 mm [(E,F)], respectively. In (A), the shear bands are dense, and the red arrows only point out two types of shear bands. The red arrows in (C) and (E) point out all shear bands.

diameters of 2, 3, and 5 mm are  $0.201$ ,  $0.100$ , and  $0.067 \text{ J g}^{-1}$ , respectively.

The typical dynamic compression true stress versus true strain curves are obtained at strain rate of  $250 \text{ s}^{-1}$  with the diameters of 2, 3, and 5 mm, respectively, as shown in **Figure 2C**. The fracture strength of the samples with diameter of 2 mm are 1,404, 1,522, and 1,455 MPa, respectively. That diameter of 3 mm are 1,444, 1,351, and 1,475 MPa, respectively. That diameter of 5 mm are 1,381, 1,378, 1,338, and 1,381 MPa, respectively. Although the data have some variability, the general trend has been determined. And the variability is due to the experimental technology to a certain extent. As we all know, the difficulty and accuracy of dynamic mechanical test are lower than that of quasi-static mechanical test. The average fracture strength of 2, 3, and 5 mm are 1460.3, 1423.3, and 1382.0 MPa, respectively. The correlation between sample size and fracture strength is given in **Figure 2D**, also the error bars are indicated in the figure. There exists an obvious trend that fracture strength decreases with the increasing of sample size. Based on previous

investigation (Huang et al., 2007; Zhang et al., 2016b), this may be caused by the number of STZ.

At the sample surfaces, as can be seen in **Figures 3A,C,E**, the density of sample's shear bands with the diameter of 2 mm is obviously larger than that of 3 and 5 mm. As observed in fractography (**Figures 3B,D,F**), both average distance and average width of tear ridges decrease with increasing of the sample diameter.

Structural relaxation exothermic heat  $\Delta H$  (i.e., the exothermic event during relaxation) prior to  $T_g$  could be regarded as an indicator of the excess STZs,  $V_{STZ}$ , described by an empirical equation as follow (Slipenyuk and Eckert, 2004; Huang et al., 2007; Evenson and Busch, 2011):

$$\Delta H = \beta' V_{STZ} \quad (1)$$

where  $\beta'$  is a constant.  $\Delta H$  have been obtained by integrating the insert diagram of **Figure 2B**.

Jiang et al. (2011) proposed a method to predict the characteristic volume of an STZ  $v_{STZ}$  in BMGs ambient temperature  $T$  under the stain rate of  $\dot{\gamma}$ :

$$\Omega = \frac{kT}{4R\zeta\mu_T\gamma_T^2(1-2\mu_T\gamma_T/(\pi\mu_0\gamma_0))^{3/2}} \left( \ln \frac{\omega_0}{C\dot{\gamma}} - \frac{\Delta V^*}{V_{STZ}} \right) \quad (2)$$

where  $\Delta V^*$  represents the required excess STZs to transition and no less than an atom volume;  $R$  is the fold ratio  $\approx 1/4$ , and  $\zeta$  is a correction factor = 3;  $k$  is Boltzmann's constant and  $\mu_0\gamma_0$  represents the threshold shear resistance at 0 K;  $\gamma_0 = 0.036$  is the critical yield strain at 0 K,  $\mu_0$  is the shear modulus at 0 K,  $\mu_T$  denotes the temperature related shear modulus, indicated by  $\mu_T = \mu_0 - d\mu_0/dT \times T$  GPa/K. The yield elastic limit has the form of  $\gamma_T = 0.036 - 0.016(T/T_g)^{2/3}$ . Here  $\ln(\omega_0/C\dot{\gamma}) \approx 30$  could be considered as a reasonable work at the typical strain rate of  $10^{-4} \sim 10^{-2} \text{ s}^{-1}$ , where  $\omega_0$  is the frequency of shear phonon of nm wavelength ( $\sim 10^{13}$  Hz);  $\dot{\gamma}$  is the strain rate, and  $C$  is a dimensionless constant of order unity (Johnson and Samwer, 2005). The detail definitions can be found in the related literatures (Johnson and Samwer, 2005; Jiang et al., 2011).

On the other hand, Jiang et al. (Jiang et al., 2008) presented that there exists a transition from STZ to tension transformation zones at the strain rate of  $10^6 \text{ s}^{-1}$ . That means when the strain rate is  $10^6 \text{ s}^{-1}$ , STZ volume  $\Omega$  tends to zero.

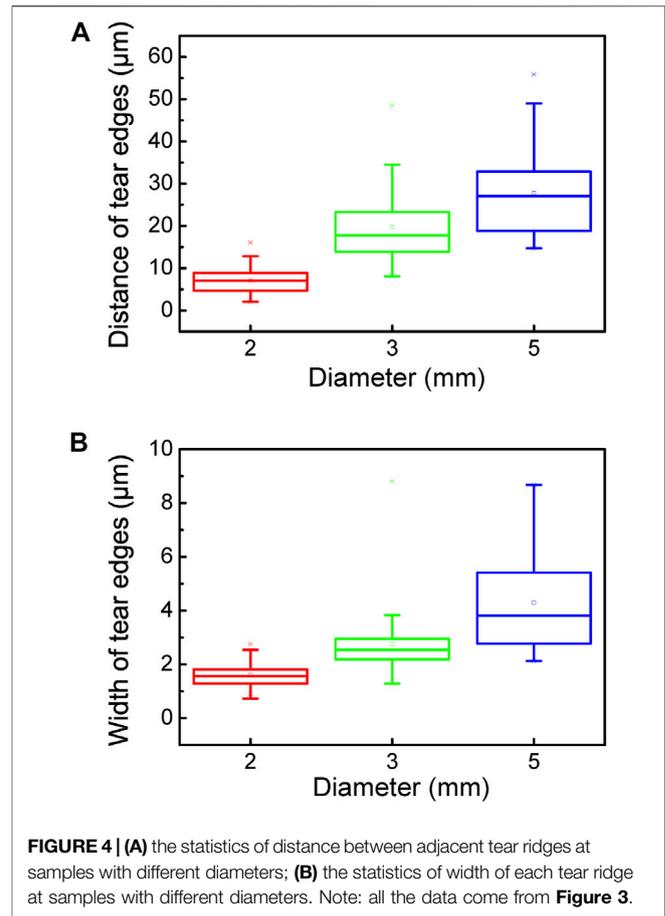
$$\ln \frac{\omega_0}{C\dot{\gamma}} - \frac{\Delta V^*}{V_{STZ}} = 0 \quad (3)$$

By assuming  $\ln(\omega_0/C\dot{\gamma}) \approx 30$  at the strain rate of  $10^{-4} \sim 10^{-2} \text{ s}^{-1}$ , one can obtain that  $\ln(\omega_0/C\dot{\gamma}) \approx 6.9$  at the strain rate of  $10^6 \text{ s}^{-1}$  for as-cast Vit1. The value is adopted in the current research. Therefore,  $\Delta V^*/V_{STZ} \approx 6.9$  for the BMG with the diameter of 2 mm, the relative STZ concentration  $V_{STZ}/\Delta V^* \approx 1/6.9$ , which is normalized by the critical volume  $\Delta V^*$  for the present component.

At  $250 \text{ s}^{-1}$ ,  $\ln(\omega_0/C\dot{\gamma})$  is approximated to be 15.2. In order to calculate the volume of an STZ at the temperature of 300 K, the parameter  $\mu_T$  requires to be determined. In this work,  $\mu_R = 30.3 \text{ GPa}$  ( $\mu_R = \mu_T|T = \text{Room temperature}$ ), which is close to that of  $\text{Zr}_{65}\text{Cu}_{15}\text{Ni}_{10}\text{Al}_{10}$  BMG (Lu et al., 2009), and  $\mu_0 = 31.5 \text{ GPa}$  with  $d\mu_0/dT = 4 \times 10^{-3} \text{ GPa/K}$ .

According to Eq. 2, the volume of STZ of the sample with diameter of 2 mm is  $1.32 \text{ nm}^3$ . This number is located between 0.8 and  $3.2 \text{ nm}^3$ , which is corresponding to the strain rate of 1800 and  $2 \times 10^{-4} \text{ s}^{-1}$  for as-cast Vit105 (Li et al., 2016). The average atomic radius can be statistically estimated by  $R = (\sum_i^n A_i r_i^3)^{1/3}$ , in which  $A_i$  is the atomic fraction and  $r_i$  is the atomic radius of each element (Pan et al., 2008), the average atomic radius of the present alloy is about 0.148 nm. Based on the dense-packing hard-sphere model of BMGs, the STZ includes about 97 atoms, which is less than the literature of Pan et al. (2008). This may be ascribed to the dramatically different strain rates.

In such circumstances, the BMGs with a larger diameter can be regarded as an annealing sample, because they both contain less STZ per unit volume. The  $V_{STZ}/\Delta V^*$  of the samples with diameter of 3 and 5 mm are estimated to be 1/13.9, 1/20.7 according to  $\Delta H$ . From Eq. 2, the volumes of STZ for the BMGs with diameter of 3 and 5 mm under  $250 \text{ s}^{-1}$  can be identified as 0.21 and  $-0.88 \text{ nm}^3$ , respectively.

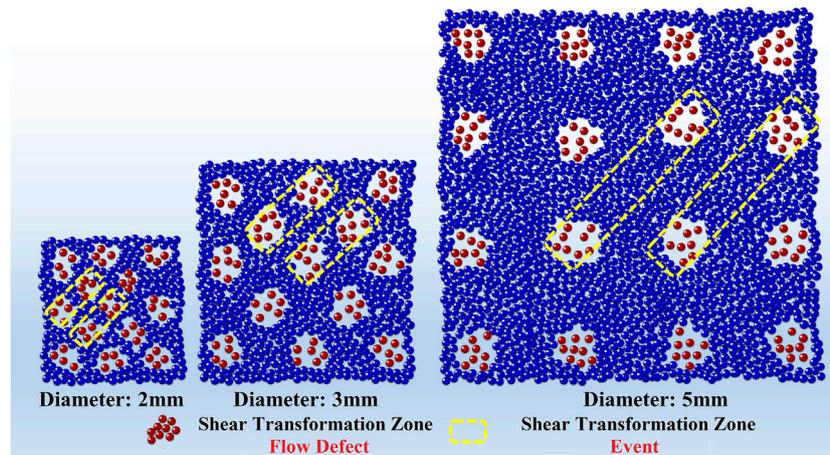


**FIGURE 4 | (A)** the statistics of distance between adjacent tear ridges at samples with different diameters; **(B)** the statistics of width of each tear ridge at samples with different diameters. Note: all the data come from **Figure 3**.

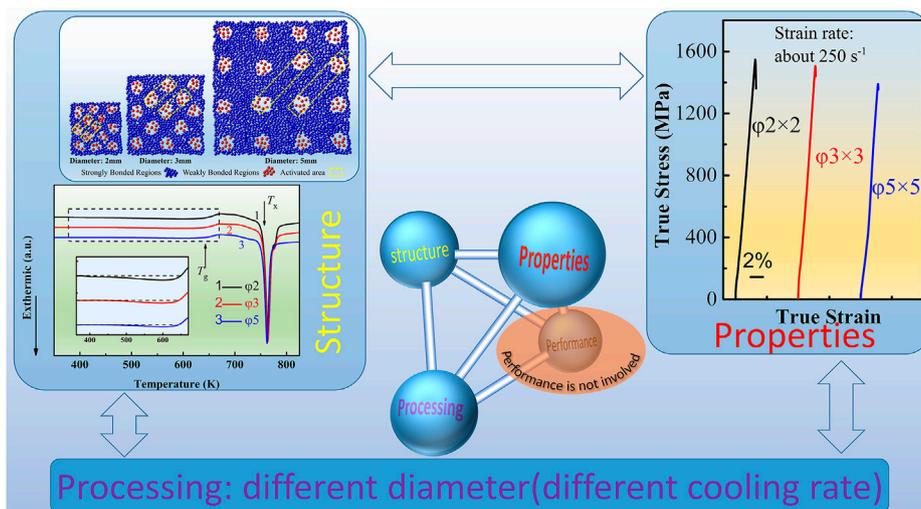
For sample with 5 mm in diameter, there appears a negative volume of STZ under  $250 \text{ s}^{-1}$ , which indicates that Eq. 2 cannot predict the volume of STZ and the obtained value is not reliable. This can be explained by a transition from ductile to brittle when sample diameter changes from 2 to 5 mm at the strain rate of  $250 \text{ s}^{-1}$ . Similar observations have also been achieved in the research of Li et al. (2016). Distinctly different fracture morphology confirm this point, as shown in Figure 3.

For further study, under the strain rate of  $10^{-4} \sim 10^{-2} \text{ s}^{-1}$  at 300 K, the volumes of STZ are calculated to be 3.68, 2.57 and  $1.47 \text{ nm}^3$  for samples with the diameter of 2, 3 and 5 mm, respectively. These values are close to  $3.89 \text{ nm}^3$  for  $\text{Zr}_{55}\text{Cu}_{25}\text{Ni}_{10}\text{Al}_{10}$  (Pan et al., 2008) and  $2.93 \text{ nm}^3$  for  $\text{Zr}_{52.5}\text{Cu}_{17.9}\text{Ni}_{14.6}\text{Al}_{10}\text{Ti}_5$  (Vit105) (Jiang et al., 2011). In general, under quasi-static conditions, the volume of STZ decreases with the increase of sample diameter rather than transitions.

At the liquid nitrogen temperature (77 K), the volumes of STZ are 0.95, 0.66 and  $0.38 \text{ nm}^3$  for samples with the diameter of 2, 3 and 5 mm, respectively. The volumes of STZ decrease continuously to be 0.21, 0.15 and  $0.08 \text{ nm}^3$  at 17 K for samples with the diameter of 2, 3 and 5 mm. These predictions are similar to 0.73 and  $0.17 \text{ nm}^3$  (as-cast Vit105) at the temperature of 77 and 17 K, respectively (Li et al., 2016). Only at 0 K will the volume of STZ reaches  $0 \text{ nm}^3$ . However, absolute zero can hardly be reached



**FIGURE 5** | Diagrammatic sketch of microscopic deformation mechanism on  $Zr_{65}Cu_{18}Ni_7Al_{10}$  BMGs with different diameters under dynamic compression.



**FIGURE 6** | The relationship among processing, structure and properties is reflected in the dynamic compression of  $Zr_{65}Cu_{18}Ni_7Al_{10}$  BMGs with different diameters.

based on the laws of thermodynamics. Thus, decreasing of temperature causes smaller volume of STZ and the limit is  $0 \text{ nm}^3$  at  $0 \text{ K}$ . Besides, it will not lead a transition, i.e., change of STZ volume from positive to negative.

Based on the above analysis, it can be concluded that there are more STZs in smaller samples, which can deduce that: 1) the volume of STZ in smaller samples are larger than the bigger one under the same condition; 2) a transition in the volume of STZ from positive to negative will take place under strain rate of  $250 \text{ s}^{-1}$  at  $300 \text{ K}$ ; 3) the volume of STZ will increase as the temperature increases for a given sample, where any transitions were improbable in theory.

However, it is still not clear how the STZs are distributed in a sample. The distribution of STZs is investigated based on the fractography. We have studied the average distance and average width of tear ridges. As shown in **Figure 4**, the average distances of

tear ridges are  $7.048$ ,  $17.776$  and  $27.040 \mu\text{m}$  in the fracture surface with the diameter of 2, 3, and 5 mm. Also, the corresponding average width of tear ridges are  $1.561$ ,  $2.541$ ,  $3.816 \mu\text{m}$ , respectively. All these data are obtained by SEM images (**Figure 3**), using pixel size and Pythagorean Theorem. Because of the equal area, the greater the total length of the shear band, the greater the shear band density. According to statistics, the total length of the shear bands in the SEM pictures of the surfaces of 2, 3 and 5 mm samples shown in **Figure 3** is  $5,029$ ,  $869$  and  $540 \mu\text{m}$ , respectively. In order to make it understood easier, the red arrows will be insert in **Figure 3**. In **Figure 3A**, the shear bands are dense, and the red arrows only point out two types of shear bands. The red arrows in **Figures 3C,E** point out all shear bands.

For the specimen with stronger STZ concentration in unit volume, the average distances of tear ridges will inevitably be

small on the fracture surface after impact load, which can be proved by experimental observation. If the dispersion degree of weak keys in the specimen is consistent, the average width of tear edges should behave similarly in specimens of different sizes. However, the experimental results show that the smaller the specimen, the smaller the average width of the tear edge. Through the above analysis, this paper speculates that the smaller specimen not only has the greater STZ concentration, but also the more dispersed. **Figure 5** shows the physical image of STZ distribution and deformation mechanism.

## CONCLUSION

In summary, the sample with smaller size exhibits better fracture strength under dynamic compression, which is ascribed to the higher concentration of STZ. The volume of STZ has been predicted. A transition from ductile to brittle at samples with different diameters is observed at the strain rate of  $250\text{ s}^{-1}$ . In addition, the volume of STZ will increase at elevated temperatures. Based on fractography, the physical distribution image of STZ is displayed. The current analysis provides several microstructure information that could be drawn from STZ concentration, which is caused by different cooling rate, i.e. processing. This can be explained by **Figure 6**.

This research have built two relationships: 1) the smaller the sample, the greater the strength, 2) the smaller the sample, the more dispersed the flow defect.

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## DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/**Supplementary Material**, further inquiries can be directed to the corresponding author.

## AUTHOR CONTRIBUTIONS

CZ and BH: conceptualization and project administration. CZ and DZ: testing, data acquisition, resources, and writing—original draft. BH: writing—review and editing. All authors contributed to the article and approved the submitted version.

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## SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fmats.2021.820158/full#supplementary-material>

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