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# Mechanical, thermal, and structural investigations of chemically strengthened Na<sub>2</sub>O–CaO–Al<sub>2</sub>O<sub>3</sub>–SiO<sub>2</sub> glasses

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For a series of conventional soda-lime-silicate glasses with increasing Al<sub>2</sub>O<sub>3</sub> content, we investigated the thermal, mechanical, and structural properties before and after K<sup>+</sup>-for-Na<sup>+</sup> ion-exchange strengthening by exposure to molten KNO<sub>3</sub>. The Al-for-Si replacement resulted in increased glass network polymerization and lowered compactness. The glass transition temperature ( $T_g$ ), hardness ( $H$ ) and reduced elastic modulus ( $E_r$ ), of the pristine glasses enhanced monotonically for increasing Al<sub>2</sub>O<sub>3</sub> content.  $H$  and  $E_r$  increased linearly up to a glass composition with roughly equal stoichiometric amounts of Na<sub>2</sub>O and Al<sub>2</sub>O<sub>3</sub> where a nonlinear dependence on Al<sub>2</sub>O<sub>3</sub> was observed, whereas  $H$  and  $E_r$  of the chemically strengthened (CS) glasses revealed a strictly linear dependence.  $T_g$ , on the other hand, showed linear increase with Al-for-Si for pristine glasses while for the CS glasses a linear to nonlinear trend was observed. Solid-state <sup>27</sup>Al nuclear magnetic resonance (NMR) revealed the sole presence of AlO<sub>4</sub> groups in both the pristine and CS glasses. <sup>23</sup>Na NMR and wet-chemical analysis manifested that all Al-bearing glasses had a lower and near-constant K<sup>+</sup>-for-Na<sup>+</sup> ion exchange ratio than the soda-lime-silicate glass. Differential thermal analysis of CS glasses revealed a “blurred” glass transition temperature ( $T_g$ ) and an exothermic step below  $T_g$ ; the latter stems from the relaxation of residual compressive stresses. The nanoindentation-derived hardness at low loads and <5 mol% Al<sub>2</sub>O<sub>3</sub> showed evidence of stress relaxation for prolonged ion exchange treatment. The crack resistance is maximized for molar ratios  $n(M_{(2)}O)/n(Al_2O_3) \approx 1$  for the CS glasses, which is attributed to an increased elastic energy recovery that is linked to the glass compactness.

## KEYWORDS

ion exchange, soda-lime-silicate glass, aluminosilicate glass, strengthening, nanoindentation, glass transition temperature, MAS NMR, network connectivity

## 1 Introduction

The readily tuned properties of glasses render them immensely useful for numerous applications (Shelby, 2005; Axinte, 2011; Zanotto and Mauro, 2017). However, a high brittleness is one disadvantage with glass (Hand and Tadjiev, 2010; Wondraczek et al., 2011), which is typically mitigated by increasing the glass thickness in many applications. Besides improved recycling procedures (Carr and Kim, 2017), a more sustainable utilization of thinner but stronger glasses reduces the amount of raw materials as well as the transportation costs of both precursors and products. There are several options for improvements, such as strengthening the glass and increasing its crack resistance (Wondraczek et al., 2011; Januchta et al., 2017; Karlsson and Wondraczek, 2021). For instance, the glass strength can be improved either by protective coatings or by introducing *residual compressive stress* (RCS) at the surface (Musil et al., 2002; Karlsson et al., 2010; Wondraczek et al., 2011; Varshneya et al., 2022). RCS may be implemented either by a conventional thermal tempering or by chemical strengthening (Wondraczek et al., 2011; Gross, 2019; Karlsson and Wondraczek, 2021).

Chemical strengthening is often referred to as *ion-exchange strengthening* because it typically involves replacing small (alkali metal) cation species by larger ones introduced from a molten salt (Wondraczek et al., 2011; Gross, 2019; Karlsson and Wondraczek, 2021), such as  $\text{Li}^+$  or  $\text{Na}^+$  substitutions by  $\text{K}^+$ . The cation size-mismatch induces compressive stresses in the glass structure, which are subsequently relieved by a rapid relaxation process followed by a slower viscosity-dependent one (Varshneya, 2016; Macrelli et al., 2019). Therefore, the ion-exchange strengthening process is performed at a temperature around or below the *strain point* (a glass viscosity fixpoint,  $\log_{10} \eta = 14.5$  dPa s). Some chemical strengthening procedures involve multiple consecutive ion-exchange processes, which are beneficial for producing an engineered stress profile (ESP) with a controlled RCS profile in the surface (Green et al., 1999). *Chemically strengthened* (CS) glasses are nowadays widely used for flat-screen covers in handheld electronic devices (Varshneya and Bihuniak, 2018), pharmaceutical cartridges (Morandotti and Zuccato, 2018), and for impact-resistant front windows in high-speed trains, aircrafts, and high-security windows (Sheikh et al., 2020); their utilization is expected to expand to various architectural (Zaccaria et al., 2021), automotive (Jacoby, 2018), solar (Allsopp et al., 2020), and flexible photonics (Macrelli et al., 2020) purposes.

Current CS glasses conform to the aluminosilicate (AS) or aluminoborosilicate glass systems (Gross, 2019; Karlsson and Wondraczek, 2021). Such “high-performance” CS glasses often incorporate large amounts of Al, which increase both the melting temperature and the melt viscosity (Silverman, 1939; Cheng et al., 2013; Karlsson, 2021), and thereby making the glasses comparatively costly to produce. The melt viscosity is maximized around a unity molar ratio of the metal oxide to

alumina (Toplis et al., 1997; Webb et al., 2007),  $n(\text{M}_{(2)}\text{O})/n(\text{Al}_2\text{O}_3) = 1$  [henceforth referred to as “ $\text{M}_{(2)}\text{O}/\text{Al}_2\text{O}_3$ ”]. Conventional soda-lime-silicate (SLS) glasses are in general unsuitable for chemical strengthening because they suffer from higher stress relaxation than high performance AS glasses (Varshneya and Kreski, 2012; Erdem et al., 2017; Güzel et al., 2019; Macrelli et al., 2019; Sun and Dugnani, 2020). Nonetheless, given that SLS glasses dominate the industrial market (Hand and Tadjiev, 2010; Deng et al., 2020), their minor doping with Al is a promising route towards an optimal balance between high chemical strengthening performance and lower production costs, and thereby more sustainable glasses (Wallenberger and Bingham, 2010; LaCourse, 2018).

To reach these goals, the structure of the glass and its properties must be contrasted before and after the chemical strengthening process. To explore and improve the understanding of the complex chemical strengthening mechanisms (Varshneya, 2010a; Macrelli et al., 2019; Karlsson and Wondraczek, 2021), we report herein on effects from  $\text{K}^+$ -for- $\text{Na}^+$  cation exchange on physical properties and local structures of the industrially relevant compositions  $16\text{Na}_2\text{O}-12\text{CaO}-x\text{Al}_2\text{O}_3-(72-x)\text{SiO}_2$  for increasing Al contents ( $x$  mol%). For both the pristine and ion-exchanged glasses, we discuss the trends of the glass transition temperature ( $T_g$ ), the density ( $\rho$ ) and its related molar volume ( $V_m$ ) and atom packing density (compactness;  $C_g$ ), crack resistance ( $CR$ ) along with the nanoindentation-derived values of the hardness ( $H$ ) and reduced elastic modulus ( $E_r$ ). The degrees of  $\text{K}^+$ -for- $\text{Na}^+$  ion exchange was estimated by  $^{23}\text{Na}$  magic-angle spinning (MAS) nuclear magnetic resonance (NMR) experiments. The Al speciation was probed by  $^{27}\text{Al}$  MAS NMR, which revealed that all structures comprised essentially only  $[\text{AlO}_4]^-$  groups, regardless of the application of chemical strengthening.

## 2 Materials and methods

### 2.1 Glass preparation and composition analysis

Six glass samples were prepared with nominal compositions  $16\text{Na}_2\text{O}-12\text{CaO}-x\text{Al}_2\text{O}_3-(72-x)\text{SiO}_2$  with  $x = \{0, 4, 8, 12, 16, 20\}$  mol%; see Table 1. The glasses were melted in 200 g glass batches in Pt-Rh crucibles at 1500°C for  $x = \{0, 4, 8\}$  and 1600°C for  $x = \{12, 16, 20\}$  using  $\text{Na}_2\text{SO}_4$  as a fining agent. The melts were quenched on a steel plate before being annealed at their respective calculated glass transition temperature. Full details of the glass sample preparation are given in ref. (Karlsson, 2021).

Parts of the prepared glass were crushed on a steel plate using a steel hammer and then powdered using a mill from Retsch (model RM200) equipped with an agate mortar and pestle. The powdered glasses were sieved into two fractions,  $<45 \mu\text{m}$  and  $<90 \mu\text{m}$ , using steel sieves and a Retsch interval shaking table (model RV).

**TABLE 1** Analyzed and nominal glass compositions (in brackets), density ( $\rho$ ), molar volume ( $V_m$ ), compactness ( $C_g$ ), oxygen packing density ( $C_O$ ), glass transition temperature ( $T_g$ ), average network connectivity ( $\bar{N}_{BO}^{Si}$ ), and the K<sup>+</sup>-for-Na<sup>+</sup> exchange ratio<sup>d,e</sup> ( $IE$ ). Parameters based on nominal composition are given in parentheses. All data originate from the glass powders (see Section 2.2.1), except for density ( $\rho$ ) and thus indirectly also molar volume ( $V_m$ ), compactness ( $C_g$ ) that originate from the glass monoliths (see Section 2.2.2).

Label	Analyzed (nominal) composition (mol%)					$M_{(2)O}/Al_2O_3$	$\rho$ (g/cm <sup>3</sup> )	$V_m$ (cm <sup>3</sup> /mol)	$C_g$	$C_O$ (mol/cm <sup>3</sup> )	$T_g^a$ (°C)	$\bar{N}_{BO}^{Si}$ <sup>b</sup>	Na/K exchange	
	Na <sub>2</sub> O	K <sub>2</sub> O	CaO	Al <sub>2</sub> O <sub>3</sub>	SiO <sub>2</sub>								$IE_{ICP}^c$ (%)	$IE_{NMR}^d$ (%)
$\pm\sigma^e$	±2	±2	±2	±4	±5	—	±0.003	±0.3	±0.0003	±0.001	±5	--	±4	±0.2
Al0	14.4 (16.0)	—	12.0 (12.0)	0.8 (0.0)	72.8 (72.0)	20.6 (∞)	2.510	23.98	0.5105 (0.5221)	0.07272 (0.07208)	574	3.30 (3.22)	—	—
Al4	14.6 (16.0)	—	12.0 (12.0)	4.7 (4.0)	68.7 (68.0)	3.1 (4.0)	2.523	24.51	0.5098 (0.5211)	0.07266 (0.07211)	605	3.36 (3.29)	—	—
Al8	15.5 (16.0)	—	11.8 (12.0)	8.6 (8.0)	64.1 (64.0)	1.8 (2.0)	2.534	25.05	0.5092 (0.5200)	0.07237 (0.07213)	632	3.42 (3.38)	—	—
Al12	15.9 (16.0)	—	12.5 (12.0)	13.5 (12)	58.0 (60.0)	1.2 (1.3)	2.538	25.79	0.5058 (0.5175)	0.07169 (0.07195)	675	3.49 (3.47)	—	—
Al16	18.6 (16.0)	—	13.7 (12.0)	17.8 (16.0)	49.9 (56.0)	1.0 (1.0)	2.551	26.40	0.5044 (0.5168)	0.07028 (0.07201)	721	3.42 (3.57)	—	—
Al20	16.3 (16.0)	—	11.9 (12.0)	19.9 (20.0)	51.8 (52.0)	0.8 (0.8)	2.564	26.59	0.5067 (0.5164)	0.07201 (0.07210)	728	3.68 (3.69)	—	—
Al0-CS 5h	4.1	11.2	13.5	0.7	70.5	21.4	—	—	—	—	—	3.20	73.2	73.6
Al4-CS 5h	4.1	11.2	13.3	5.0	66.5	3.1	—	—	—	—	—	3.29	73.4	--
Al8-CS 5h	5.4	10.7	12.7	9.2	61.9	1.8	—	—	—	—	—	3.37	66.3	65.4
Al12-CS 5h	4.7	11.0	12.9	14.3	57.1	1.1	—	—	—	—	—	3.50	69.9	—
Al16-CS 5h	6.2	11.6	14.3	19.0	48.9	0.9	—	—	—	—	—	3.46	65.3	66.6
Al20-CS 5h	6.2	11.1	13.1	23.4	46.2	0.7	—	—	—	—	—	3.70	64.1	67.2

<sup>a</sup>Glass transition temperature ( $T_g$ ) was determined by the onset of the first endothermic step of the DTA, curve; *q.v.* Figure 5.

<sup>b</sup>Average number of bridging oxygen atoms at the SiO<sub>4</sub> groups, calculated either from each analysed or nominal glass compositions (values based on nominal within parentheses) according to the procedure of (Edén, 2011) and assuming the absence of Al-NBO bonds and that all Al sites are four coordinated.  $\bar{N}_{BO}^{Si}$  for the K<sup>+</sup>-exchanged glass was calculated from the analysed compositions.

<sup>c</sup>Extent of K<sup>+</sup>-for-Na<sup>+</sup> exchange ( $IE_{ICP}$ ), as calculated by  $IE_{ICP} = n(K_2O)/[n(Na_2O) + n(K_2O)]$  of the ICP-analysed glass compositions.

<sup>d</sup>Extent of K<sup>+</sup>-for-Na<sup>+</sup> exchange ( $IE_{NMR}$ ), where the remaining of the Na<sup>+</sup> reservoir is given by (1- $IE$ ).  $IE_{NMR}$  was calculated from the difference of the integrated <sup>23</sup>Na NMR signal intensities before ( $I_{pristine}$ ) and after ( $I_{CS}$ ) cation exchange, according to  $IE_{NMR} = (I_{pristine} - I_{CS})/I_{pristine}$ .

<sup>e</sup>Data uncertainties specified as  $\pm 1\sigma$ .

The glass compositions were determined by inductively coupled plasma optical emission spectroscopy (ICP OES) with an Aviro 200 instrument (Perkin Elmer). For each glass sample, 0.5 g of powdered material was dissolved in 15 ml 48% HF, 1 ml H<sub>2</sub>SO<sub>4</sub> (95%–97%) and 0.5 ml HNO<sub>3</sub> (65%) altogether diluted to 50 ml using deionized water. The solution was evaporated, and the residual oxides were dissolved in 10 and 30 ml HCl and deionized H<sub>2</sub>O respectively under mild heating on a hot plate for 60 min (occasionally deionized H<sub>2</sub>O was added to replace evaporated H<sub>2</sub>O). The solution was next transferred to a 250 ml beaker and 150 ml of deionized H<sub>2</sub>O was added, whereupon the solution was evaporated to a volume of  $\approx$ 100 ml. It was subsequently transferred to a volumetric flask and diluted to a total volume of 250 ml. The cation contents of the solutions were measured by ICP OES, whereas the corresponding O contents were determined from the charge

requirements of the analyzed cations. Periodic Table Mix one for ICP measurements (SIGMA ALDRICH) was used as reference standards [33 elements in 10% HNO<sub>3</sub>(aq)].

## 2.2 Chemical strengthening

### 2.2.1 Glass powders

About 35 g of reagent grade KNO<sub>3</sub> salt (ACROS Organics; 99% purity) and 5 g of the <45  $\mu$ m sieved glass powder, *q.v.*, Section 2.1, were heated to 450°C in separate fused quartz crucibles in a Nabertherm muffle furnace (model N2OH). The glass powder was then put into the molten KNO<sub>3</sub> and was stirred. One more stirring was performed after 2.5 h. After 5 h, the crucibles were taken out from the furnace and left to cool to ambient temperature on a refractory brick. Deionized water was

added under magnetic stirring to dissolve the  $\text{KNO}_3$ , whereafter, the glass powder was separated from the aqueous solution using a Munktell 3-grade filter paper and a Büchner funnel.

### 2.2.2 Glass monoliths

For each prepared glass sample, two glass monoliths of approximate size  $2 \times 2 \times 2$  cm were cut out by a diamond saw. The glasses were polished in the Struers LaboForce100 equipment using the pertinent sample holder fixated by the UniForce equipment. SiC #320 sheets (MD Gekko) were used for rough polishing before fine polishing by the sheets MD Plan, MD Dac and MD Nap (in that order) with their corresponding polishing solutions. The samples were cleaned using Struers Lavamin equipment between the polishing steps.

The polished glass pieces were then subjected to ion-exchange strengthening in a Hybe muffle furnace (model KUT180) at  $450^\circ\text{C}$ . The glass pieces were heated in a wire net basket separately to a stainless-steel container with the salt. The salt bath consisted of 99% purity  $\text{KNO}_3$  from ACROS Organics. The salt bath was first melted at  $450^\circ\text{C}$  before the basket with the glass pieces was put inside for 5 and 16 h. The basket was then taken out and the furnace was turned off with the door left a little open so that the cooling rate was ca.  $3\text{--}5^\circ\text{C}/\text{min}$ . The salt bath temperature was controlled to be within  $\pm 5^\circ\text{C}$  to the target using a thermocouple of type K using a P655 logger from Dustmann Electronic GmbH. The treatment temperature,  $450^\circ\text{C}$ , was chosen specifically to give some stress relaxation (Svenson et al., 2016). The glass series from these treatments are called Al-CS 5 h and Al-CS 16 h.

## 2.3 Glass characterization

### 2.3.1 Solid state nuclear magnetic resonance experiments and data analysis

All  $^{23}\text{Na}$  (spin 3/2) and  $^{27}\text{Al}$  (spin 5/2) solid-state NMR experiments were performed with a Bruker Avance-III spectrometer at a magnetic field of 14.1 T, which provided the Larmor frequency of  $-158.7$  MHz for  $^{23}\text{Na}$  and  $-156.4$  MHz for  $^{27}\text{Al}$ . A fine powder ( $<10$   $\mu\text{m}$  particles) of each glass specimen was filled in a 3.2 mm zirconia rotor that underwent magic-angle spinning (MAS) at 24.00 kHz. To obtain NMR spectra that quantitative reflect the site populations (Edén, 2020), short radio-frequency (rf) pulses were employed:  $0.40$   $\mu\text{s}$  ( $12^\circ$  flip angle; 83 kHz nutation frequency) for  $^{23}\text{Na}$ , and  $0.32$   $\mu\text{s}$  ( $10^\circ$  flip angle; 87 kHz nutation frequency) for  $^{27}\text{Al}$ . The relaxation delays were 5 s ( $^{23}\text{Na}$ ) and 2 s ( $^{27}\text{Al}$ ), with 2024–10240 and 4048–8192 signal transients accumulated per  $^{23}\text{Na}$  and  $^{27}\text{Al}$  NMR acquisition, respectively, with the precise number depending on the Na and Al content of the glass.  $^{23}\text{Na}$  and  $^{27}\text{Al}$  shifts are quoted relative to 0.1 M  $\text{NaCl}(\text{aq})$  and 1.0 M  $\text{Al}(\text{NO}_3)_3(\text{aq})$ , respectively.

Iterative numerical fitting of the MAS NMR spectra were used for determining the average values of the isotropic chemical shift,  $\bar{\delta}_{\text{iso}}^E$ , and quadrupolar product,  $\bar{C}_{Q\eta}^E$ , of each detected  $E = \{^{27}\text{Al}, ^{23}\text{Na}\}$  species in the respective  $^{27}\text{Al}$  and  $^{23}\text{Na}$  NMR spectra.  $\bar{C}_{Q\eta}^E$  is the average quadrupolar product over all  $N$   $^{27}\text{Al}$  (or  $^{23}\text{Na}$ ) sites in the structure,

$$\bar{C}_{Q\eta}^E = \sqrt{\frac{1}{N} \sum_{j=1}^N (C_Q^j)^2 (1 + \eta_j^2/3)} \quad (1)$$

where  $C_Q^j = e^2 q_j Q/h$  and  $\eta_j$  are the quadrupolar coupling constant and asymmetry parameter of the electric field gradient (EFG) tensor of  $^{27}\text{Al}^j/^{23}\text{Na}^j$ , respectively (Edén, 2020). Each NMR spectrum was fitted by assuming one  $^{27}\text{Al}$  or  $^{23}\text{Na}$  nuclear site with NMR-parameter distributions (see Section 3.1). The best-fit parameters did not vary significantly if the central-transition (CT) peak was fitted alone or if the spinning sidebands of the satellite transitions were also included. The fitting was performed by in-lab developed software that minimizes the root-mean-square deviation between numerically simulated NMR spectra and the experimental counterpart while varying the targeted  $\{\bar{\delta}_{\text{iso}}^E, \bar{C}_{Q\eta}^E, W_{\text{iso}}^E\}$  parameter triplet; e.g., see (Iftekhar et al., 2012) for further information. Here,  $W_{\text{iso}}^E$  is the width of the (as-assumed) Gaussian distribution of isotropic chemical shifts, whereas the Czjzek model accounted for the quadrupolar-product distribution (Czjzek et al., 1981).

The numerically exact simulations accounted for both first- and second-order quadrupolar interactions. The carousel-averaging COMPUTE algorithm (Levitt and Edén, 1998) was employed to efficiently exploit the time periodicity and for analytical averaging over the “ $\gamma$ ” Euler angle, which together with a  $^{\text{ROSE}}\text{LEBhemi18817}$  set of  $\{\alpha, \beta\}$  orientations (Stevansson and Edén, 2006) ensured a fully converged powder average. The smallest time-integration step in the solution of the Schrödinger equation was 81 ns. The initial density and detection operators were  $\rho(t=0) = I_x$  and  $I^+$ , respectively.

### 2.3.2 Scattered light polariscope experiments

The surface RCS of the CS glass monoliths,  $\sigma_s$ , were determined using SCALP-05, a scattered light polariscope equipment from GlasStress Ltd. It is a well-described method to determine RCS in glass, see (Hödemann et al., 2016; Sun et al., 2019; Hödemann et al., 2020). The measurement depth was 0.75 mm, the stress optic coefficient ( $\Delta C$ ) was assumed to be constant,  $\Delta C = 2.72$   $\text{TPa}^{-1}$ , and the refractive index ( $RI$ ) was assumed to be 1.52 for  $x = \{0, 4, 8, 12\}$  and 1.53 for  $x = \{16, 20\}$  where  $x$  is the  $\text{Al}_2\text{O}_3$  content; *q.v.*, Section 2.1, and Table 1. The glass monoliths were cleaned with conventional window cleaning solution and dried clean with a wipe. A thin layer of Cargille immersion oil (liquid code 5040,  $RI = 1.5200$  at 632.8 nm) was sprayed onto the sample. Three repeated measurements were performed for each sample and the average value was used. The stress was calculated from the gradient of the fitted optical retardation curve which starts from two to three pixels into

the sample, i.e., ca. 10  $\mu\text{m}$ , hence the errors of the determined stresses are estimated to be  $\pm 30\%$ . In fact,  $\sigma_s$  is rarely overestimated with SCALP-05 (only for ultra-high  $\sigma_s$ ) but for simplicity the error is displayed as both under- and overestimation.

### 2.3.3 Simultaneous thermal analysis experiments

The <45  $\mu\text{m}$  glass powder (*q.v.*, Section 2.1) and the CS glass powder (*q.v.*, Section 2.2.1) were run in a simultaneous thermal analysis (STA), NETSZCH STA 409 PC/PG. Both the differential thermal analysis (DTA) and thermal gravimetric analysis (TGA) signals were recorded.  $\text{Al}_2\text{O}_3$  crucibles were used for the measurements that were run under  $\text{N}_2$ -gas flow. The heating rate was 10 K/min up to 1200°C. An empty crucible was used for baseline correction. The glass transition temperature ( $T_g$ ) was determined by using the Proteus software provided by NETSZCH using standard procedure where each  $T_g$  value was determined from the onset of the first endothermic step of the DTA curve (Zheng et al., 2019). The TGA analyses of the glass powders, available in the dataset; *q.v.*, Karlsson et al. (2022), revealed no significant mass losses during the measurements.

### 2.3.4 Density, molar volume and compactness

Densities were measured using Archimedes method in deionized water ( $\rho_w = 0.998 \text{ g/cm}^3$ ). The water temperature was kept within  $21.5^\circ\text{C} \pm 0.5^\circ\text{C}$  for the density measurements. Density values of the glass samples ( $\rho$ ) were calculated using

$$\rho = \frac{m_d}{m_d - m_w} \cdot \rho_w \quad (2)$$

where  $m_d$  is the weight of the dry glass sample,  $m_w$  the weight of the sample immersed in deionized water and  $\rho_w$  being the density for water. Sample weights varied between 6 and 10 g and the densities were estimated to be reproducible to  $\pm 0.003 \text{ g/cm}^3$ .

The molar volume of the glass sample was calculated using

$$V_m = \frac{\sum_i x_i M_i}{\rho} \quad (3)$$

where  $x_i$  is the molar fraction,  $M_i$  the molar mass of element  $i$  and  $\rho$  the density of the glass. The compactness (atom packing density) was calculated according to

$$C_g = \frac{\sum_i x_i V_i}{V_m} \quad (4)$$

where  $V_i$  is the volume contribution from oxide  $i$  ( $i = M_xO_y$ ) calculated from the ion radii,  $r_M$  and  $r_O$  (Rouxel, 2007), by  $V_i = \frac{4}{3}\pi N_A (x r_M^3 + y r_O^3)$ , and  $N_A$  is Avogadro's constant. The ionic radii {1.18, 1.12, 0.39, 0.26, 1.35} in  $\text{\AA}$  were taken from (Shannon, 1976) with the coordination numbers (CN) {8, 8, 4, 4, 2} for the ions { $\text{Na}^+$ ,  $\text{Ca}^{2+}$ ,  $\text{Al}^{3+}$ ,  $\text{Si}^{4+}$ ,  $\text{O}^{2-}$ } (Cormier and Neuville, 2004; Le Losq et al., 2014).

The oxygen packing density ( $C_O$ ), which is the molar amount of oxygen per unit volume of the glass (Ray, 1974; Grammes et al., 2020), was calculated according to

$$C_O = \frac{N_A}{V_m} \sum_i x_i n_i^O \quad (5)$$

where  $n_i^O = \{1, 1, 3, 2\}$  is the stoichiometric coefficient of the amount of O in the respective oxide { $\text{Na}_2\text{O}$ ,  $\text{CaO}$ ,  $\text{Al}_2\text{O}_3$ ,  $\text{SiO}_2$ } component.

### 2.3.5 Nanoindentation experiments

Hardness ( $H$ ) and the reduced elastic modulus ( $E_r$ ) were determined by nanoindentation on polished glass samples before and after ion-exchange treatment using the Oliver-Pharr method (Oliver and Pharr, 1992).  $H$  is defined by the maximum applied load ( $F_m$ ) divided by the projected contact area ( $A_p$ );  $H = F_m/A_p$ .  $A_p$  is calculated by a polynomial fitting,  $A_p(h_c) = C_0 h_c^2 + C_1 h_c^1 + C_2 h_c^{1/2} + C_3 h_c^{1/4} + \dots + C_8 h_c^{1/128}$ , where  $C_x$  is the indenter specific factors ( $C_0 = 24.56$  for a perfect Berkovich tip) and  $h_c$  is the real contact depth.  $h_c$  takes into account the sink-in effect which is calculated by  $h_c = h_m - \frac{\epsilon F_m}{S}$ , where  $h_m$  is the maximum penetration depth,  $\epsilon$  is a tip factor ( $\epsilon = 0.72$  for a Berkovich tip) and  $S$  is the stiffness as determined from the slope upon unloading,  $S = \frac{\partial P}{\partial h}$ . The reduced elastic modulus ( $E_r$ ) is determined through  $E_r = \frac{\sqrt{\pi}}{2\beta} \frac{S}{\sqrt{A_p}}$ , where  $\beta$  is a geometrical tip factor ( $\beta = 1.034$  for a Berkovich indenter).

Nanoindentation was performed using a NHT<sup>2</sup> instrument from Anton Paar with indentation loads 1, 5, 10, 15, 25, 50 and 75 mN. The number of indents was 40 for 1 mN load and 20 for the other load levels. The indentations had the following settings: acquisition rate 10 Hz, loading and unloading rate two times the max load per min (i.e., 2 mN/min for the max load of 1 mN etc.), holding time at max load 10 s, approach speed 4  $\mu\text{m/s}$  and the stiffness threshold 500  $\mu\text{N}/\mu\text{m}$ . The Poisson's ratio ( $\nu$ ), that is affecting the determination of  $E_{IT}$  by  $E_{IT} = (1 - \nu^2) / (\frac{1}{E_r} - \frac{(1 - \nu^2)}{E_i})$ , where  $E_i$  and  $\nu_i$  are the elastic modulus and Poisson's ratio of the indenter tip.  $E_{IT}$  is only reported on in the dataset; see Karlsson et al. (2022), was hypothesized to be constant,  $\nu = 0.23$ , although  $\nu$  can be slightly affected by the increasing  $\text{Al}_2\text{O}_3$  content (Luo et al., 2016b; Pönitzsch et al., 2016; Ren et al., 2019) and affected by the chemical strengthening (Ryou et al., 2020). Collected nanoindentation data are available in the dataset (*q.v.*, Karlsson et al., 2022) and for 75 mN are selected data given in Supplementary Table S1.

### 2.3.6 Microindentation experiments

The crack resistance ( $CR$ ) was measured on polished glass monoliths before and after ion-exchange treatment using a Micro-Combi Tester (MCT) from CSM Instruments. The Vickers microindentations were used to produce radial crack statistics in the glass samples. The  $CR$  method is described in (Sundberg et al., 2019), which follows the original procedure of (Kato et al., 2010). 15 indents were made for each load using a



Vickers indenter. The least square method was used for fitting to the Weibull cumulative function,

$$PCI = 1 - e^{-(x_p/x_c)^m} \quad (6)$$

where  $PCI$  is the probability of crack initiation,  $x_p$  is the load in N,  $x_c$  is the characteristic value in N and  $m$  the Weibull modulus. The  $CR$  is then defined as the load when the  $PCI$  is 50%. However, the Weibull cumulative function,  $x_c$ , can likewise be used as a relative measure of the  $CR$ , cf. previous paper (Karlsson, 2022). The Vickers indentations were made with an acquisition rate of 10 Hz using a linear loading with an (un)loading rate of two times the max load per min (i.e., 2N/min for the max load 1N etc.). The holding time at load was 15 s. The approach speed was 8  $\mu\text{m}/\text{min}$  and the retract speed 16.6  $\mu\text{m}/\text{min}$ . The contact force was 30 mN and the contact stiffness threshold 25,000  $\mu\text{N}/\mu\text{m}$ . The indentations were made in the load range up to 30N and in the displacement range up to 1000  $\mu\text{m}$ . All indentations, both nano and micro were performed at  $23^\circ\text{C} \pm 2^\circ\text{C}$  and relative humidity of  $40\% \pm 10\%$ . We estimated the error of the  $CR$  and  $x_c$  to be  $\pm 20\%$  except for the samples with  $m > 10$  where the error is lower (cf., Supplementary Table S2).

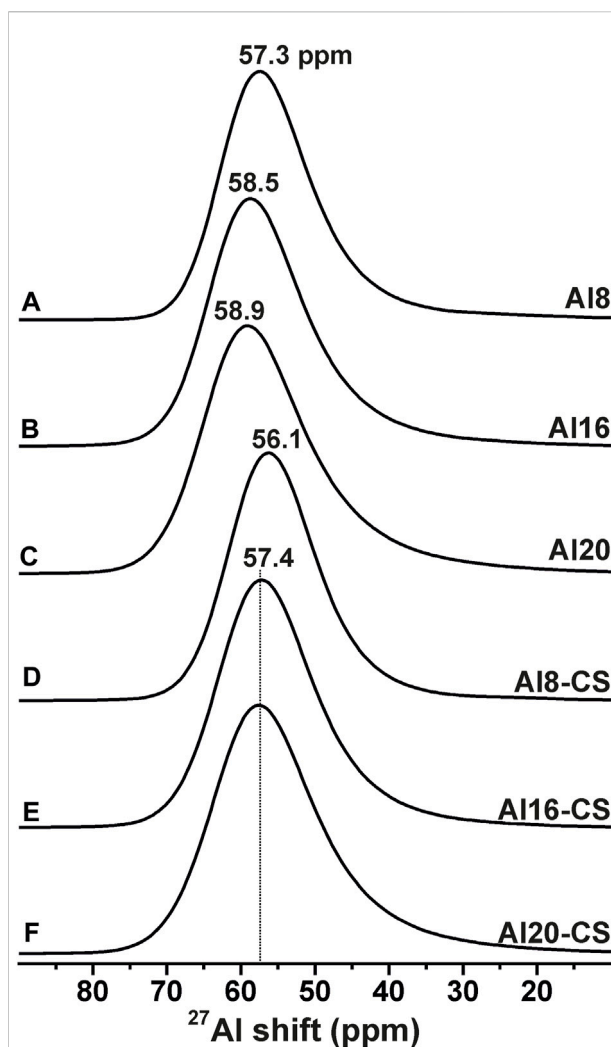
## 3 Results

### 3.1 Solid state nuclear magnetic resonance

#### 3.1.1 $^{27}\text{Al}$ magic-angle spinning nuclear magnetic resonance

Owing to the low/moderate cation field strength of the  $\text{Na}^+$  and  $\text{Ca}^{2+}$  ions of the  $\text{Na}_2\text{O}-\text{CaO}-\text{Al}_2\text{O}_3-\text{SiO}_2$  glasses, a traditional network model is expected to hold well for their structures, meaning that the AS glass network only involves  $[\text{AlO}_4/2]^-$  groups, surrounded by  $\text{SiO}_4/2$  tetrahedra that accommodates all non-bridging oxygen (NBO) ions (Engelhardt and Michel, 1987; Edén, 2020). The sole presence of tetrahedrally coordinated Al species is verified by the  $^{27}\text{Al}$  MAS NMR spectra presented in Figures 1A–C, which were recorded from the Al8, Al16, and Al20 glasses. The observed  $^{27}\text{Al}$  shifts at the peak maxima,  $\delta_{\text{max}}$ , remain in the 57–59 ppm range, and are consistent with  $\text{AlO}_4$  groups (Engelhardt and Michel, 1987; Edén, 2020). The slight “tailing” of the  $^{27}\text{Al}$  NMR peak shapes towards lower  $^{27}\text{Al}$  shifts stems from distributions of  $^{27}\text{Al}$  quadrupolar products,  $\bar{C}_{Q\eta}^{\text{Al}}$  (Edén, 2020).

For increasing  $\text{Al}_2\text{O}_3$  content, the NMR-peak maxima in Figure 1 displace concurrently towards higher shifts. From the best-fit  $^{27}\text{Al}$  NMR parameters listed in Table 2 follows that the observed peak displacements originate from a concomitantly increased (average) isotropic  $^{27}\text{Al}$  chemical shift ( $\bar{\delta}_{\text{iso}}^{\text{Al}}$ ) when  $\text{SiO}_2$  is replaced by  $\text{Al}_2\text{O}_3$ . This is most transparent from Figure 2A that plots both  $\delta_{\text{max}}$  and  $\bar{\delta}_{\text{iso}}^{\text{Al}}$  against the  $\text{Al}_2\text{O}_3$  content, revealing



**FIGURE 1**  
 $^{27}\text{Al}$  NMR spectra recorded at 14.1 T and 24.00 kHz MAS from the  $\text{Na}_2\text{O}-\text{CaO}-(\text{Al}_2\text{O}_3)-\text{SiO}_2$  glasses before (A–C) and after (D–F) exchange with  $\text{K}^+$  ions. The peak maximum (in ppm) is indicated at the top of each NMR spectrum.

a linear relationship between  $\bar{\delta}_{\text{iso}}^{\text{Al}}$  and  $n(\text{Al}_2\text{O}_3)$ . Along with the increased isotropic  $^{27}\text{Al}$  chemical shift, the quadrupolar products also manifest a concurrent minor increase, which accounts for roughly half of the net  $\sim 2$  ppm increase in the full-width-at-half-maximum height (fwhm) observed between the  $^{27}\text{Al}$  NMR responses from the Al8 and Al20 glasses; the remaining  $\sim 1$  ppm peak-broadening stems from a very minor increased chemical-shift dispersion ( $W_{\text{iso}}^{\text{Al}}$ ).

The negatively charged  $[\text{AlO}_4]^-$  moieties are expected to be charge-compensated primarily by the  $\text{Na}^+$  cations, whereas the remaining  $\text{Na}^+$  reservoir, along with that of  $\text{Ca}^{2+}$ , mainly act as glass network modifiers that introduce—and charge balance—NBO anions at the  $\text{SiO}_4$  groups (Mysen et al., 1981; Lee and Sung, 2008; Gambuzzi et al., 2014a). Consequently,

TABLE 2  $^{27}\text{Al}$  and  $^{23}\text{Na}$  NMR parameters before and after chemical strengthening.<sup>a</sup>

Glass sample	$^{27}\text{Al}$ NMR parameters					$^{23}\text{Na}$ NMR parameters				
	$\delta_{\text{max}}^{\text{Al}}$ (ppm)	fwhm (ppm)	$\delta_{\text{iso}}^{\text{Al}}$ (ppm)	$\bar{C}_{\text{O}^{\eta}}^{\text{Al}}$ (MHz)	$W_{\text{iso}}^{\text{Al}}$ (ppm)	$\delta_{\text{max}}^{\text{Na}}$ (ppm)	fwhm (ppm)	$\delta_{\text{iso}}^{\text{Na}}$ (ppm)	$\bar{C}_{\text{O}^{\eta}}^{\text{Na}}$ (MHz)	$W_{\text{iso}}^{\text{Na}}$ (ppm)
$\pm\sigma^b$	$\pm 0.10$	$\pm 0.2$	$\pm 0.6$	$\pm 0.3$	$\pm 0.3$	$\pm 0.2$	$\pm 0.3$	$\pm 0.6$	$\pm 0.5$	$\pm 0.4$
Al0	—	—	—	—	—	-9.7(-4.9)	24.4(24.2)	-1.9(3.7)	3.33(3.69)	18.6(16.8)
Al8	57.3(56.1)	13.9(13.7)	61.9(60.3)	5.46(4.98)	9.0(9.7)	-12.6(-6.8)	22.2(21.2)	-5.3(0.6)	2.85(3.24)	18.5(16.4)
Al16	58.6(57.0)	14.8(14.9)	63.7(61.9)	5.82(5.46)	9.6(10.4)	-12.6(-7.1)	21.0(19.3)	-5.8(-0.3)	2.70(3.09)	18.0(15.0)
Al20	58.9(57.4)	16.0(16.2)	64.6(62.9)	6.18(5.94)	10.2(10.8)	-12.4(-6.9)	20.9(18.9)	-5.8(-0.4)	2.64(3.03)	18.1(14.8)

<sup>a</sup>NMR shifts at the peak maximum,  $\delta_{\text{max}}^{\text{Al}}$  or  $\delta_{\text{max}}^{\text{Na}}$ , of the respective  $^{27}\text{Al}$  and  $^{23}\text{Na}$  peaks of the MAS NMR spectra along with their corresponding full width at half maxima (fwhm) and the triplet of best-fit parameters  $\{\delta_{\text{iso}}^E, \bar{C}_{\text{O}^{\eta}}^E, W_{\text{iso}}^E\}$  for  $E = \{\text{Al}, \text{Na}\}$ ; see Section 2.3.1 for notation. The values within parentheses are those obtained *after* chemical strengthening.

<sup>b</sup>Data uncertainties specified as  $\pm 1\sigma$ .

the progressive  $\text{Al}_2\text{O}_3$ -for- $\text{SiO}_2$  substitutions (at constant  $\text{Na}_2\text{O}$  and  $\text{CaO}$  contents) are accompanied by a concurrently increased polymerization of the AS glass network. This feature is reflected by an increase of the average number of bridging oxygen (BO) atoms per  $\text{SiO}_4$  group ( $\bar{N}_{\text{BO}}^{\text{Si}}$ ) (Edén, 2011), from  $\bar{N}_{\text{BO}}^{\text{Si}} = 3.3$  of the parent  $16\text{Na}_2\text{O}-12\text{CaO}-72\text{SiO}_2$  glass to  $\bar{N}_{\text{BO}}^{\text{Si}} = 3.7$  of the Al-richest sample of stoichiometry  $16\text{Na}_2\text{O}-12\text{CaO}-20\text{Al}_2\text{O}_3-52\text{SiO}_2$  (Table 1). The glass-network polymerization for increasing Al content is also reflected in enhanced physical glass properties, such as an increased glass transition temperature, hardness, and reduced elastic modulus (see Section 3.3 and Section 3.5). Such trends for increasing Al content are well documented (Cheng et al., 2013; Takahashi et al., 2015; Ragoen et al., 2017b).

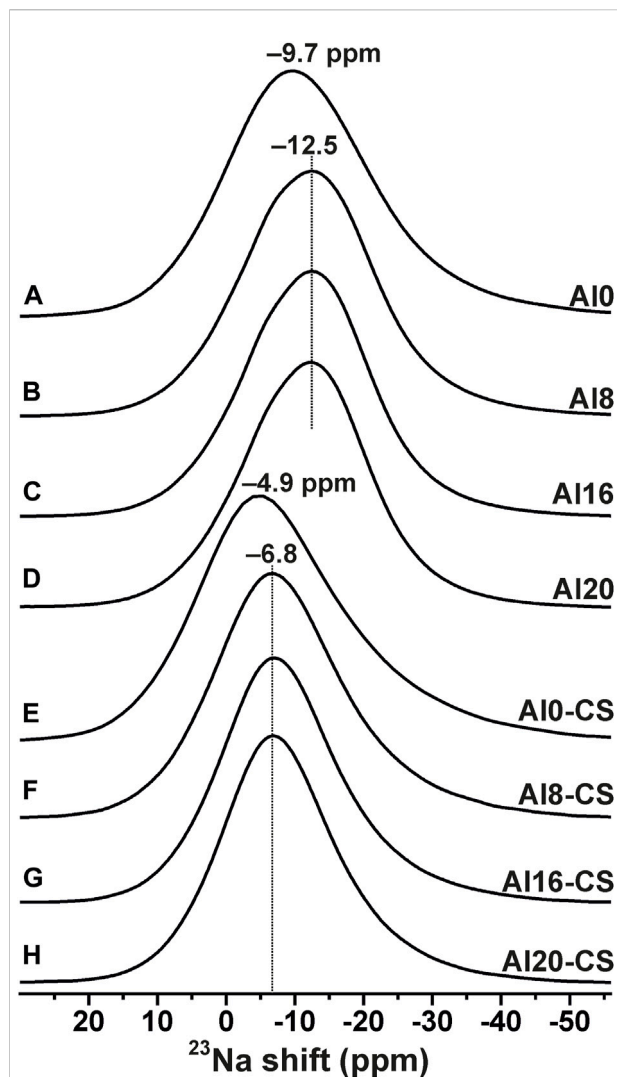
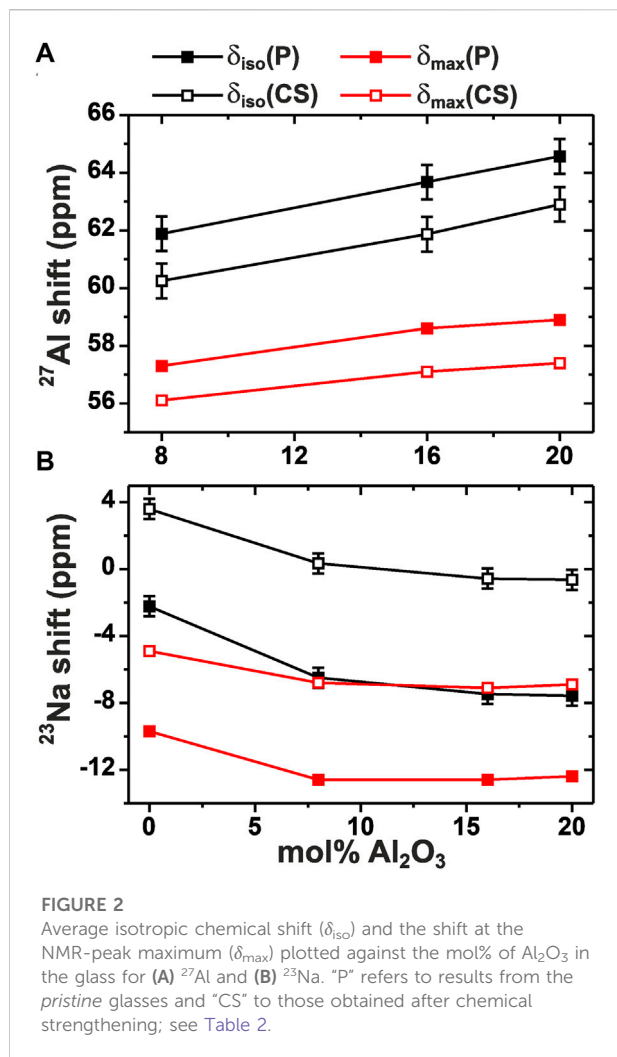
Upon the  $\text{K}^+$ -for- $\text{Na}^+$  exchange, the  $^{27}\text{Al}$  MAS NMR spectra (Figures 1D–F) reveal only minor alterations, which underlie the following NMR-parameter trends (Table 2): (i) Although the linear correlation between the isotropic chemical shift and  $n(\text{Al}_2\text{O}_3)$  remains intact after chemical strengthening, the ion-exchanged glass exhibits a 1–2 ppm lower  $\delta_{\text{iso}}^{\text{Al}}$  value at a fixed Al content (Figure 2A), which is attributed to the partial  $\text{K}^+$ -for- $\text{Na}^+$  replacement, where  $\text{K}^+$  also charge compensate the  $[\text{AlO}_4]^-$  groups. (ii) Despite that the net fwhm value of the  $^{27}\text{Al}$  MAS NMR peak remains essentially unaltered after chemical strengthening (Table 2), the quadrupolar product is slightly reduced by 0.25–0.5 MHz. Apparently, the accompanying NMR-peak narrowing is counter-balanced by a  $\sim 0.7$  ppm broadening from a slightly larger chemical-shift dispersion, thereby giving an overall unaltered  $^{27}\text{Al}$  NMR peak-width (Table 2). Although the emphasized shift dispersion is consistent with the larger chemical/structural disorder created by mixed  $\text{K}^+/\text{Na}^+$  cations in the second coordination sphere of the  $^{27}\text{Al}^{[4]}$  nuclei, the decrease in  $\bar{C}_{\text{O}^{\eta}}^{\text{Al}}$  is surprising and difficult to rationalize, as it suggests an overall enhanced local ordering around the  $^{27}\text{Al}$  nuclei of the  $\text{AlO}_4$  groups.

### 3.1.2 $^{23}\text{Na}$ magic-angle spinning nuclear magnetic resonance

We next consider the  $^{23}\text{Na}$  MAS NMR results, where Figure 3A depicts the NMR spectrum from the SLS glass. As anticipated from the spin-3/2 nature of  $^{23}\text{Na}$ , its resonance is broad and slightly tailed towards lower shifts (*vide supra*). On the introduction of Al in the glass network, the  $^{23}\text{Na}$  NMR peak maximum shifts from  $-9.7$  ppm (Figure 3A) to around  $-12.5$  ppm in the spectra from the Al-bearing glasses Al8, Al16, and Al20; see Figures 3B–D and Table 2. This trend stems from an isotropic chemical shift decrease from  $\delta_{\text{iso}}^{\text{Na}} = -1.9$  ppm in the SLS glass towards the asymptotic value  $\delta_{\text{iso}}^{\text{Na}} = -5.8$  ppm for the two Al-richest glasses. This shift effect is attributed to the *partial* relocation of the  $\text{Na}^+$  cations from being associated with the Si–NBO moieties of the silicate glass to instead be positioned around the  $[\text{AlO}_4]^-$  tetrahedra. Note, however, that the slight  $^{23}\text{Na}$  shift displacements of  $\delta_{\text{max}}$  and  $\delta_{\text{iso}}^{\text{Na}}$  may also stem from an altered distribution of Na coordination numbers.

Relative to the pristine glasses, drastical  $^{23}\text{Na}$  shift-changes are observed in the NMR spectra obtained from the series of  $\text{K}^+$ -for- $\text{Na}^+$  exchanged glasses displayed in Figures 3E–H, where the shift at the peak maximum is increased by 5–6 ppm (Table 2). As for  $^{27}\text{Al}$ , this deshielding trend mirrors that of the isotropic  $^{23}\text{Na}$  chemical shift (Figure 2B), which is increased by  $\sim 6$  ppm upon chemical strengthening. The same  $^{23}\text{Na}$  deshielding trend has been reported for Na–Ca–Si–O glasses upon chemical strengthening by  $\text{K}^+$  incorporation (Ragoen et al., 2017a), and accords also with elevated-pressure induced effects observed in  $^{23}\text{Na}$  MAS NMR from AS glasses (Lee, 2010).

The dependence of the isotropic  $^{23}\text{Na}$  chemical shift on structural parameters is complicated but is known to increase for a decrease in either the Na coordination number ( $\text{Na}^{[p]}$ ) or the average Na–O distance (Xue and Stebbins, 1993; George and Stebbins, 1995; Charpentier et al., 2004; Gambuzzi et al., 2014b;

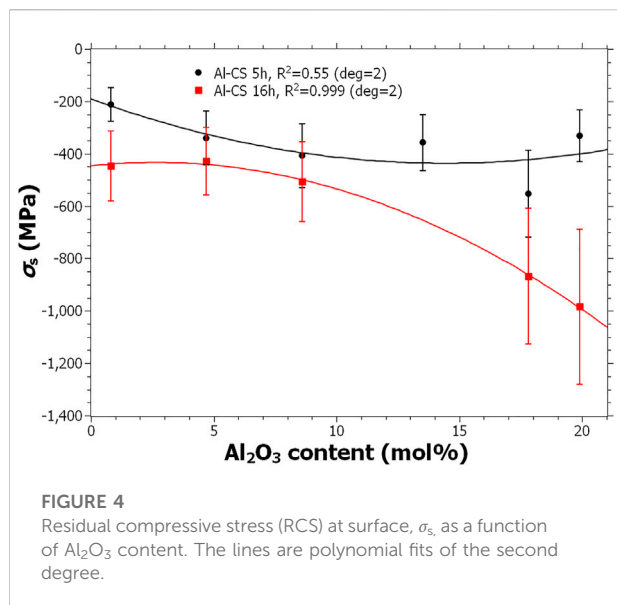


Yu et al., 2020), where we note that these structural parameters are usually correlated. Hence, the observed  $^{23}Na$  shift increase upon the  $K^+$ -for- $Na^+$  exchange may stem from one or both of the following: (i) a shift in the distribution of  $Na^{[p]}$  average coordination towards lower  $p$ , where the larger  $K^+$  cation outcompetes  $Na^+$  for the highest coordination sites in the glass structure, or (ii) a net reduction in the average Na–O distances (with unaltered coordination numbers), as concluded by (Ragoen et al., 2018) or tacitly assumed (Ragoen et al., 2017a) from cation-exchanged Na/Ca-based silicate glasses. Moreover, unlike the scenario of Al, the partial replacement of  $Na^+$  cations by  $K^+$  leads to a minor increase in the  $^{23}Na$  quadrupolar product by  $\sim 0.4$  MHz, the elevation of which remains constant, regardless of the presence/absence of Al in the glass (or the precise Al content); see Table 2.

From the integrated  $^{23}Na$  NMR signal intensities of the pristine and ion-exchanged glasses in Figure 3 and the assumption of an equimolar  $K^+$ -for- $Na^+$  exchange, we

estimated that a majority of the  $Na^+$  cations (74%) was exchanged in the Al-free SLS glass (AI0), whereas a somewhat reduced exchange efficiency ( $\sim 65\%$ ) resulted for the Al-bearing counterparts. Incidentally, the exchange-degree trend of an initial large drop at no/low  $Al_2O_3$  contents, followed by near-constant values for increasing  $x \geq 8$  mol%, which match that of  $\delta_{iso}^{Na}$  in Figure 2B. The NMR-derived cation-exchange fractions agree very well with those estimated from the analyzed glass compositions (Table 1), but with a superior accuracy. The slightly lower cation exchange of the Al-rich glasses may be ascribed to their more polymerized—and thereby more rigid—AS glass networks, where the charge-balance requirement of the





tetrahedrally coordinated Al to form  $[\text{AlO}_4]^- \cdots \text{Na}^+$  pairs naturally reduce the  $\text{Na}^+$  mobility and thereby the cation exchange.

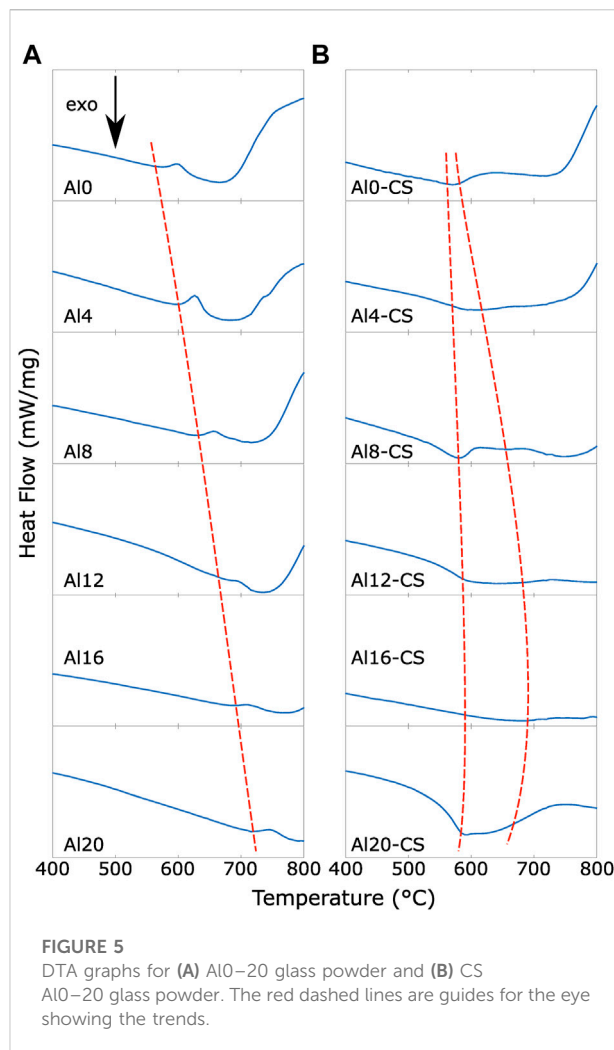
### 3.2 Residual compressive stress (RCS)

Figure 4 shows the RCS as a function of the  $\text{Al}_2\text{O}_3$  content and the recorded  $\sigma_s$  values are given in Supplementary Table S1. The 5 and 16 h treatment do not follow the same trend; the 5 h follows a positive second degree polynomial function while the 16 h follows a negative second degree polynomial function. The RCS results will be further discussed in relation to other properties in the forthcoming sections.

### 3.3 Glass transition temperature ( $T_g$ ) and sub- $T_g$ relaxation

Figure 5 displays the DTA data obtained from both the pristine and CS glass powders. All pristine samples exhibit a well-defined glass transition, with the  $T_g$  values increasing concurrently with the  $\text{Al}_2\text{O}_3$  content, in good agreement with literature findings (Cheng et al., 2013; Takahashi et al., 2015; Ragoen et al., 2017b). The determined  $T_g$  for the investigated glasses are given in Table 1. The second endothermic peak, which we are not investigating in this paper, we anticipate being the onset of the crystallization range of the glasses.

The  $T_g$ s of the CS glasses are diffuse and not reported due to their large uncertainties. In contrast to the pristine glasses, the first step of the DTA curve is characterized by an exothermic step (hereinafter referred to as a sub- $T_g$  step).



We anticipate the observed sub- $T_g$  step to stem from relaxation processes of the chemical-strengthening-induced RCS, analogous to the sub- $T_g$  step observed from mechanically induced excess enthalpy studied by (Martin et al., 2005). The sub- $T_g$  step coincides well to the point where stresses start to relax, i.e., slightly above the strain point. The Al8-CS glass manifests 2  $T_g$  peaks, which is attributed to the presence of two separate segments/phases of each glass particle they presumably differ primarily in their glass modifier composition, constituting ion-exchange-induced  $\text{K}^+$ -bearing outer segment and a core/bulk of the pristine Al8 glass (similar to a liquid-liquid phase separation (Zheng et al., 2020)). Such inhomogeneities rationalize the ill-defined  $T_g$  values of the CS glasses. Moreover, for the Al12-CS, Al16-CS and Al20-CS glasses, the endothermic steps are particularly diffuse, which is attributed to slower relaxation stemming from an increased viscosity due to the high Al content of these glasses (Karlsson, 2021). Al16-CS also exhibit a diffuse sub- $T_g$  step. Apparently, the Al16-CS

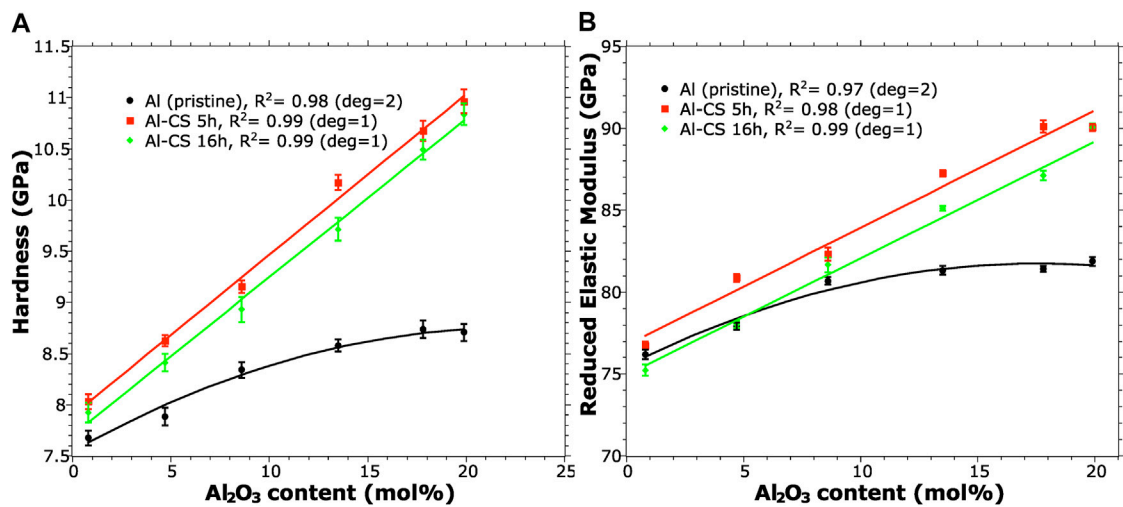


FIGURE 6

(A) Hardness ( $H$ ) and (B) Reduced Elastic Modulus ( $E_r$ ) for 75 mN nanoindentation as a function of  $\text{Al}_2\text{O}_3$  content in mol%. The solid lines are polynomial fittings with the  $R^2$  and polynomial degree given in the legend.

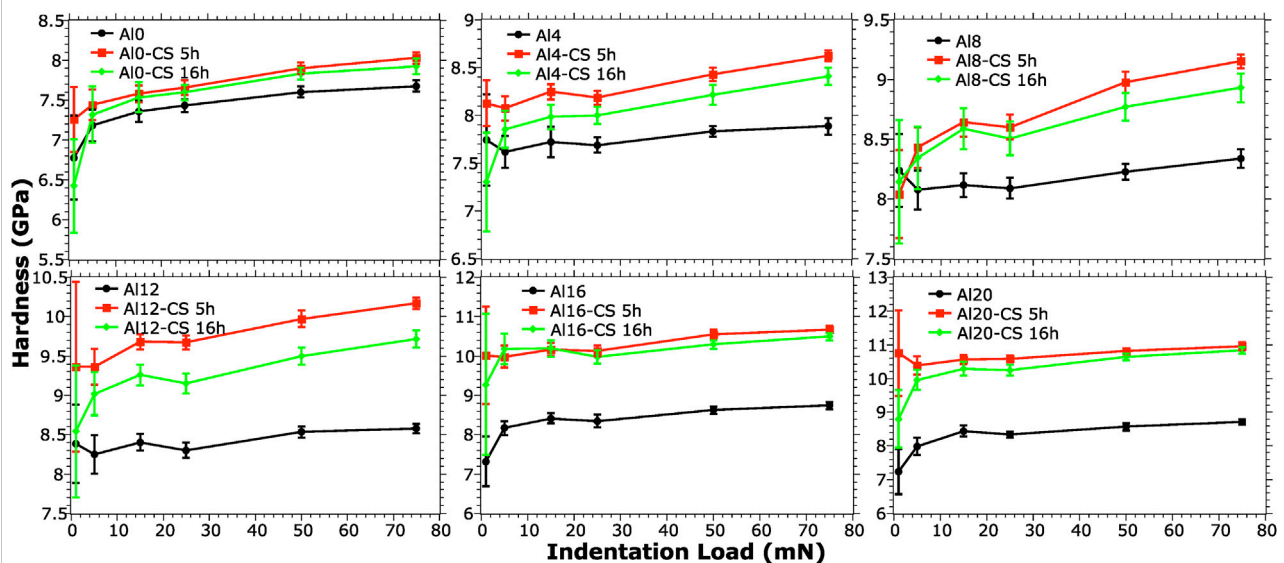


FIGURE 7

Hardness ( $H$ ) as a function of indentation load for each sample, pristine and CS-glasses for 5 and 16h respectively. The solid lines are guides for the eye.

glass exhibits the highest temperatures for the sub- $T_g$  step and  $T_g$  peak, as is shown by the dashed guiding lines in Figure 5.

The trends of Figure 5 correspond well to the recorded  $\sigma_c$  data shown in Figure 4 and given in Supplementary Table S1. Furthermore, the results of Figure 5 suggest that the

relaxation of RCS might be investigated by differential scanning calorimetry (DSC) measurement on ion-exchanged glass powder or discs in analogy with the method of (Yue et al., 2002) for fictive-temperature estimates. However, more work is required to verify this concept.

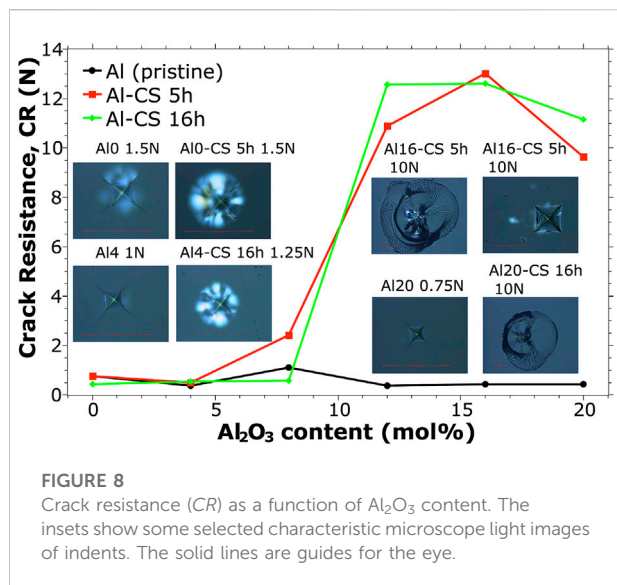


FIGURE 8

Crack resistance (CR) as a function of  $\text{Al}_2\text{O}_3$  content. The insets show some selected characteristic microscope light images of indents. The solid lines are guides for the eye.

### 3.4 Density, molar volume and compactness

The results of Table 1 reveal a concomitant increase in the glass density and molar volume when silica is replaced by alumina, in good agreement with previous findings (Cormier and Neuvill, 2004; Ragoen et al., 2017b). These results may be understood from the free-volume expansion that accompanies the increased average Na–O bond length upon the Na–NBO–Si→Na–BO–Al bond replacements occurring when  $\text{Na}^+$  gradually shifts from balancing the NBO anions of silicate groups to BO atoms of  $[\text{AlO}_4]^-$  moieties (Cormier and Neuvill, 2004; Webb et al., 2007; Mysen, 2021). It is thus reasonable that the free volume increase with increasing Al. In general, free volume should increase mobility, and is observed for  $\text{Na}^+$  self-diffusion coefficients (Frischat, 1975). However, for increasing Al content of the studied glasses, the exchange ratio (*IE*) is rather reduced; see Table 1. We emphasize that *IE* does not necessarily reveal the same trend as the effective diffusion coefficients as a function of  $\text{Al}_2\text{O}_3$  but do give information about the effectiveness of the ion exchange,  $C_{\text{K}^+}/(C_{\text{K}^+} + C_{\text{Na}^+})$ , where  $C_{\text{M}^+}$  is the concentration of the alkali cation  $\text{M}^+$  in the surface. In the ternary system ( $\text{Na}_2\text{O}-\text{Al}_2\text{O}_3-\text{SiO}_2$ ), additions of  $\text{Al}_2\text{O}_3$  decreases the effective diffusion coefficients, though *n.b.*, that  $\text{Na}_2\text{O}-\text{SiO}_2$  give a less effective ion exchange in the surface (Ragoen et al., 2017b).

Along with an increase in the free volume, the  $C_g$  is decreased for increasing alumina content. The decrease of the  $C_g$  has also been observed in the ternary  $\text{Na}_2\text{O}-\text{Al}_2\text{O}_3-\text{SiO}_2$  system (Bechgaard et al., 2016).  $C_g$  is linearly decreased for  $\text{Al}_2\text{O}_3 < 10$  mol% and then shifting to a nonlinear trend for  $\text{Al}_2\text{O}_3 > 10$  mol% which to some extent is caused by the deviation from the nominal

composition, see Table 1. The  $C_g$  can affect other properties (Giri et al., 2004; Rouxel, 2007) and thus the structural moieties influencing the  $C_g$  may be a key for understanding the network dilation coefficient besides glass network connectivity (Tandia et al., 2012; Wang et al., 2017). For instance, the change of packing between Na and K (in alkali silicates) is the highest change for all the alkali ions as shown from the study of (Giri et al., 2004) and is a reason to why Na and K is most suitable for chemical strengthening compared to other alkali-pairs. The  $C_O$  remains approximately invariant when  $\text{SiO}_2$  is replaced by  $\text{Al}_2\text{O}_3$ .

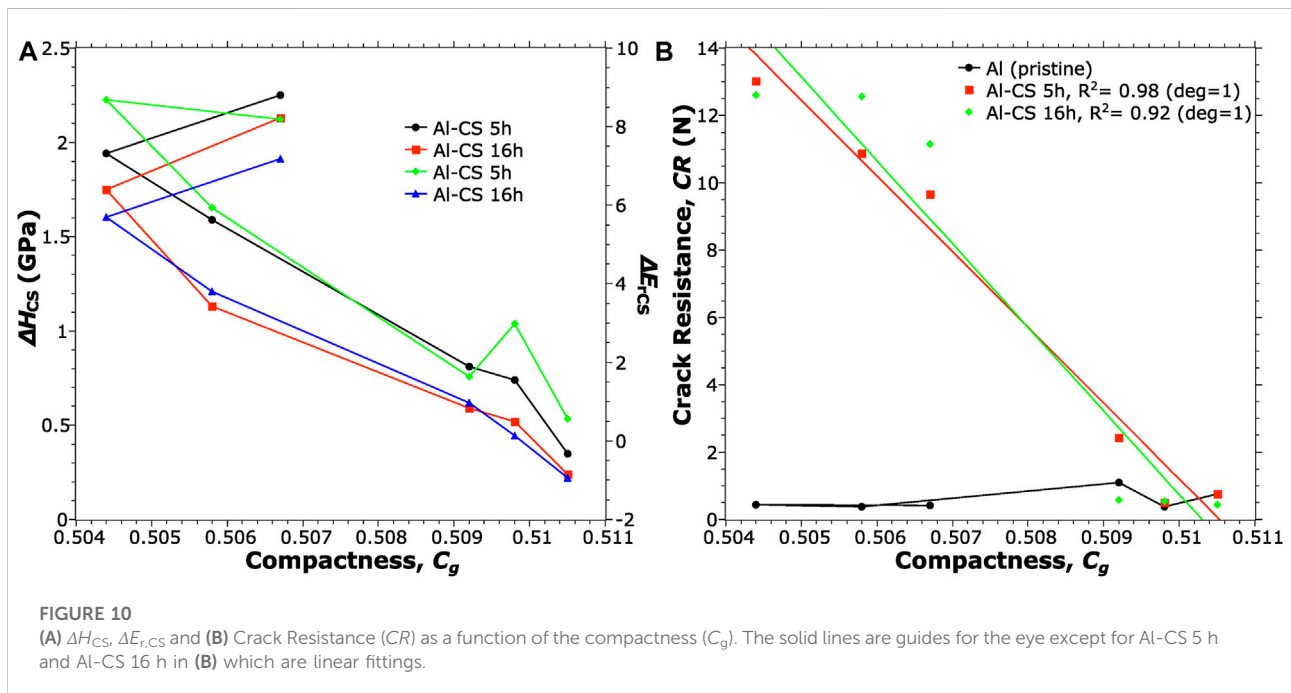
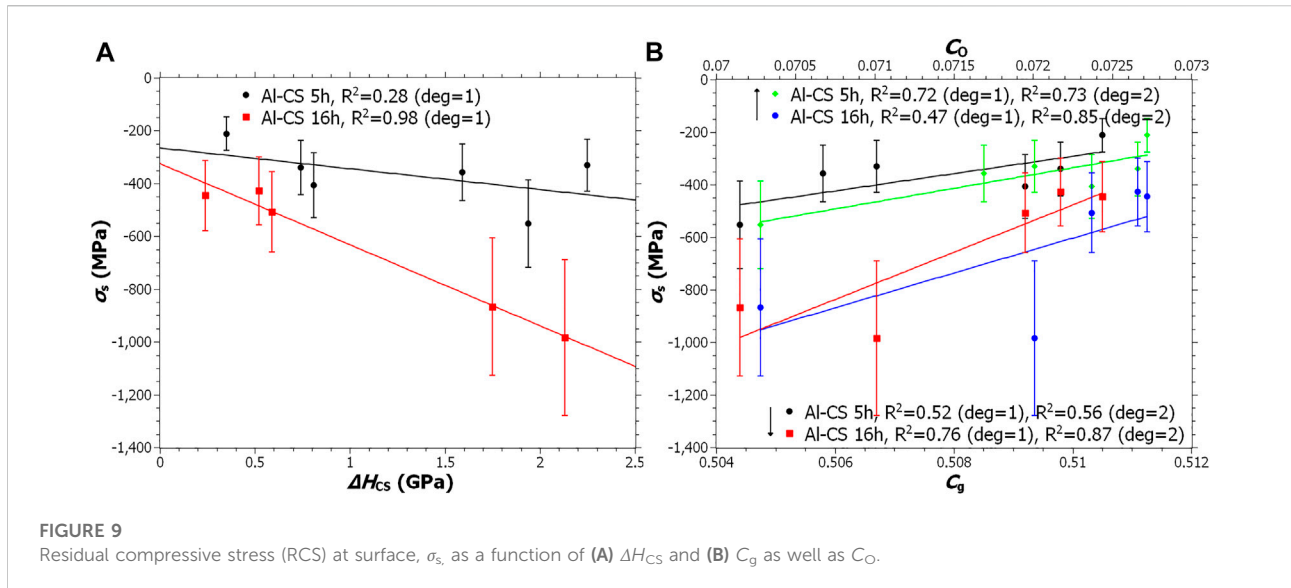
### 3.5 Hardness and reduced elastic modulus

Figure 6 shows  $H$  and  $E_r$  as a function of  $\text{Al}_2\text{O}_3$  content for 75 mN load. For the pristine glasses,  $H$  and  $E_r$  increase with the  $\text{Al}_2\text{O}_3$  content, at first linearly, but for  $M_{(2)}\text{O}/\text{Al}_2\text{O}_3 \approx 1$  nonlinearly. The results follow similar trends as reported previously (Yoshida et al., 2004). For the CS series, the trends are linear for both  $H$  and  $E_r$ . For  $\text{Al}_2\text{O}_3 < 5$  mol% in the CS 16 h series have lower and approximately the same  $E_r$  as for the untreated series.

Figure 7 plots the hardness ( $H$ ) against the indentation load for each glass before and after chemical strengthening. Generally, regardless of the precise load and treatment time, the chemical strengthening give an increase in  $H$ . However, the 16 h CS glasses typically give lower hardness values than the 5 h CS glasses in all cases except for  $M_{(2)}\text{O}/\text{Al}_2\text{O}_3 \approx 1$  (Table 1) where the 16 h give higher hardness than 5 h for the 5 and 15 mN loads. *N.b.*, that at 1 mN loads and  $\text{Al}_2\text{O}_3$  contents  $< 10$  mol%,  $H$  is lower for the 16 h CS glasses than for the pristine glasses (see further discussion in Section 4.1). In Supplementary Figure S1, the  $\Delta H_{\text{CS}}$  as a function of load and compositions are shown, where  $\Delta H_{\text{CS}}$  is given by  $H_{\text{CS}} - H_{\text{pristine}}$  (Karlsson, 2022). In general, the trends are linearly increasing with the  $\text{Al}_2\text{O}_3$  content, the linear fitting  $R^2$  for the 5–75 mN loads of the 5 and 16h CS glasses gives 0.87–0.97 and 0.85–0.94 respectively. For 1 mN load, the trend for  $\Delta H_{\text{CS}}$  is entirely different; as a function of  $\text{Al}_2\text{O}_3$  content, it initially shows a decrease and a clear minimum for Al8 before increasing linearly. The results will be further discussed in Section 4.1.

### 3.6 Crack resistance

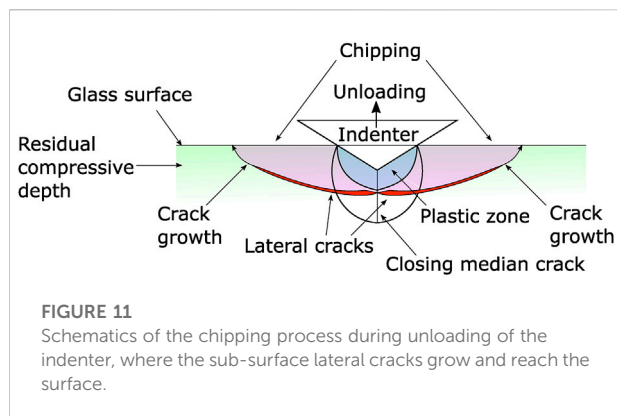
Figure 8 shows the dependence of  $CR$  on the  $\text{Al}_2\text{O}_3$  content. The pristine series show a marginal decrease with increasing  $\text{Al}_2\text{O}_3$  concentration, as previously observed for  $\text{Na}_2\text{O}-\text{Al}_2\text{O}_3-\text{SiO}_2$  glasses (Yoshida et al., 2004). For the 5 h CS glasses, there is a clear increase from  $\text{Al}_2\text{O}_3 > 8$  mol% to a maximum for  $M_{(2)}\text{O}/\text{Al}_2\text{O}_3 \approx 1$ ; see Table 1. The 16h CS glasses also show a maximum for  $M_{(2)}\text{O}/\text{Al}_2\text{O}_3 \approx 1$ . Both 5 and 16 h CS



glasses show a decrease for Al<sub>20</sub>. The results for Al<sub>0</sub> are lower than reported for a conventional SLS float glass (Sundberg et al., 2019), which could be explained by higher surface quality of float glass (Güzel et al., 2020) or sharpness of tip (Gross, 2012). CR for the Al-bearing glasses are similar to what has been reported by (Morozumi et al., 2015). The CR including the Weibull fitting results are given in Supplementary Table S2. The results and its explanation will be discussed in Section 4.2.

## 4 Discussion

In general, AS glasses frequently exhibit a turning point for properties  $M_{(2)}O/Al_2O_3 \approx 1$ . Especially transport properties, such as melt viscosity, exhibit clear maxima in both the ternary and the quaternary  $Na_2O-(CaO)-Al_2O_3-SiO_2$  systems (Toplis et al., 1997; Webb et al., 2007; Karlsson, 2021). However, other properties conventionally also give a trend-shift at  $M_{(2)}O/Al_2O_3 \approx 1$  (Smedskjaer et al., 2013; Cormier, 2021). This originates from the



preference for the alkali ions to act as charge compensators of the negatively charged  $\text{AlO}_4$  groups; *q.v.*, Section 3.1.1. This has now also been shown to account for effects for CS glasses as shown in Section 3.2, Section 3.3, Section 3.5, and Section 3.6. Furthermore, trends in physical properties as a function of the  $\bar{N}_{\text{BO}}^{\text{Si}}$  generally shows same behavior as for against  $\text{Al}_2\text{O}_3$  content, i.e., deviations from linearity around the  $M_{(2)}\text{O}/\text{Al}_2\text{O}_3 \approx 1$  for both pristine and CS glasses. In Section 4.1, respectively Section 4.2 will primarily the causes of the  $\Delta H_{\text{CS}}$  and the CR results be discussed.

#### 4.1 Hardness, reduced elastic modulus and residual compressive stress

For a progressing  $\text{K}^+$ -for- $\text{Na}^+$  ion exchange, the glass densifies while compressive stresses accumulate at the surface (Luo et al., 2016a); *q.v.* Figure 4. Along with hot compression,  $H$  is often increased by that kind of densification (Svenson et al., 2017). Therefore, it is reasonable that the hardness increases upon chemical strengthening treatment which thus is linked to the RCS (Jang, 2009), as has previously been indicated by (Svenson et al., 2016) and supported by the correlation between  $\sigma_s$  (*cf.*, Supplementary Table S1) and  $\Delta H_{\text{CS}}$  in Figure 9A. The  $\sigma_s$  data of 5h is likely less accurate than 16h due to the lower depth-of-layer and it is reflected in the fit.  $\sigma_s$  as a function of  $\Delta E_{r,\text{CS}}$  give a slightly better fit,  $R^2$  0.37 and 0.98 for 5 and 16h respectively. *N.b.*, that RCS is not of the same nature as for thermally strengthened glass (Kese et al., 2004). Free volume and thus also  $C_g$  (*q.v.*, Section 3.4) is linked to  $\nu$  (Rouxel, 2007) which is a governing property for the ability to densify; lower  $C_g$  gives higher densification-ability; see Figure 4 in (Rouxel et al., 2008), and related text for detailed reasoning. The support for this is manifested by Figure 9B, where  $\sigma_s$  as a function of  $C_g$  and even  $C_O$  give a fair correlation. The lack of perfect correlation is wrecked by a combination of the  $\sigma_s$  error and the complexity of the compressive stress buildup and relaxation, which is a function of other properties as well (Dugnani, 2020), thus more studies on this are needed.

Figure 10A plots  $\Delta H_{\text{CS}}$  and  $\Delta E_{r,\text{CS}}$  against  $C_g$ . The trend is quite linear until it turns after  $M_{(2)}\text{O}/\text{Al}_2\text{O}_3 \approx 1$  (Table 1). In contrast to a recent report on the  $\text{Na}_2\text{O}-\text{CaO}-\text{TiO}_2-\text{SiO}_2$  system (Karlsson, 2022), the present  $\text{Na}_2\text{O}-\text{CaO}-\text{Al}_2\text{O}_3-\text{SiO}_2$  glasses reveal nearly linear correlations between  $V_m$  and the ratio  $E_r/H$ , *cf.* Supplementary Figure S2.  $E_r/H$  is essentially a measure of the elastic volume recovery which will be discussed more in Section 4.2 (Rouxel, 2015). Similar correlation as for Figure 9A can be found for  $\sigma_s$  as a function of  $E_r/H$  as well.

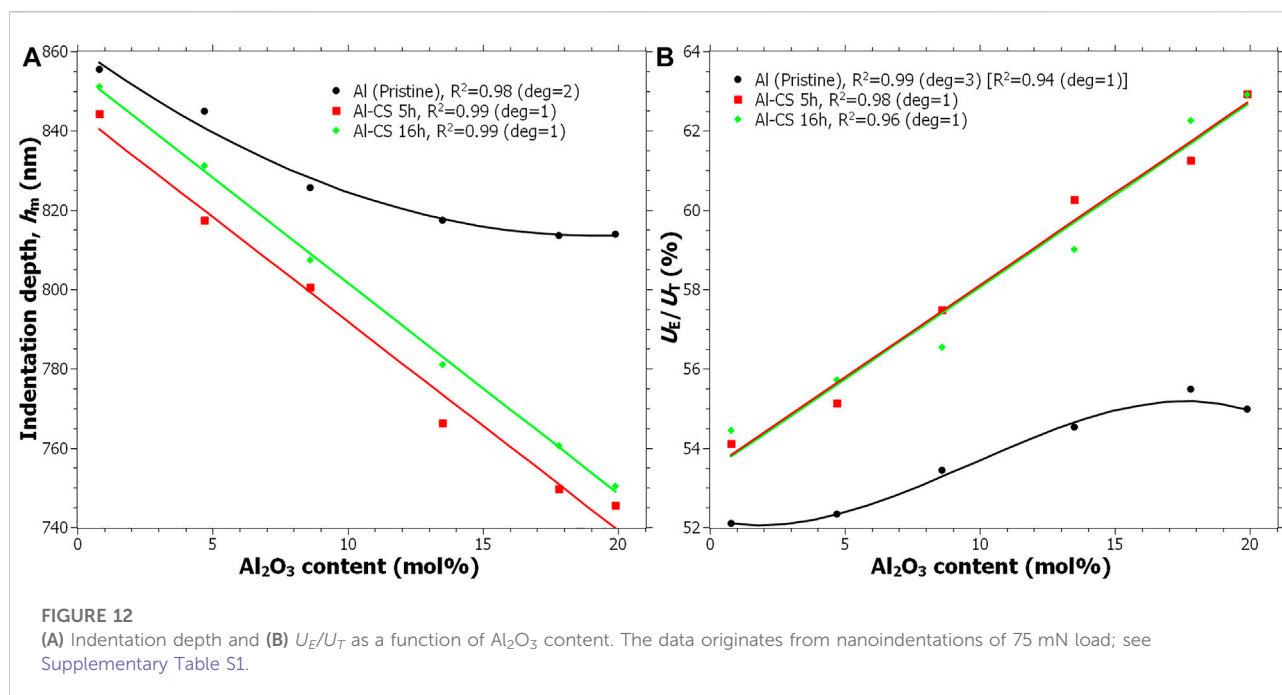
The mixed alkali effect (MAE) is well-known to primarily affect cation-transport-related properties, but has also been shown to influence mechanical properties (Kjeldsen et al., 2014; Aakermann et al., 2015; Grund Bäck et al., 2019); for  $n(\text{K}_2\text{O})/[n(\text{Na}_2\text{O}) + n(\text{K}_2\text{O})] \approx 1$ , both  $H$  and  $E$  is lower than for ratios  $n(\text{K}_2\text{O})/[n(\text{Na}_2\text{O}) + n(\text{K}_2\text{O})]$  close to 0 (Kjeldsen et al., 2014; Weigel et al., 2016). This applies also to isostatically compressed glasses (Aakermann et al., 2015). For CS glasses, the MAE is not manifested in the same way as for as-melted glass (Varshneya, 2010a).

For our CS glasses the origin of  $\Delta H_{\text{CS}}$  cannot be caused by the MAE, instead it must be caused by RCS and its accompanied densification. However, at 1 mN loads (*q.v.*, Figure 6 and *cf.* Supplementary Figure S1) notable trend-breakers can be observed where  $\Delta H_{\text{CS}}$  are negative for Al0, Al4 and Al8. We hypothesize this, and the fact that 16h CS glasses generally show lower  $\Delta H_{\text{CS}}$ , to be caused by viscosity driven stress relaxation. Thus,  $H$  then resembles the  $\text{K}_2\text{O}$ -rich end of the MAE of  $H$  (i.e., giving a lower value than for pristine glass (Kjeldsen et al., 2014; Aakermann et al., 2015)). Thus, for relaxed CS glasses we believe MAE plays a role. We hope to confirm this reasoning in a future study.

#### 4.2 Crack resistance

Upon application of an indenter on a glass surface, the energy is dissipated by elastic and plastic deformation (Yoshida, 2019; Rouxel et al., 2021; Varshneya et al., 2022). These processes differ in that when the load is removed, the elastic deformation is completely reversed, whereas the plastic deformation leaves a permanent imprint. The energy dissipation by plastic deformation can be categorized into two processes: densification, and shear flow (Yoshida, 2019; Rouxel et al., 2021; Varshneya et al., 2022). At sufficiently high indenter force the ability of the glass to deform elastically and plastically is not sufficient, instead the energy is dissipated as cracks. These cracks can either form during loading or during unloading (Yoshida, 2019; Rouxel et al., 2021; Varshneya et al., 2022). The crack type depends on the glass composition, but brittle materials exhibit most frequently radial and half-penny cracking (Yoshida, 2019; Rouxel et al., 2021), i.e., cracks originating from the corners of the Vickers indenter.





In CS glasses, (Morris et al., 2004), showed that the radial crack initiation is inhibited by RCS, but suggested that lateral cracking, i.e., sub-surface cracks parallel to the surface; see Figure 11, is enhanced by the RCS. However, Morris *et al.* stated that lateral cracks are not as critical for the glass-strength degradation as the radial counterparts, at least not until they reach the surface. Lateral and Median cracking are sub-surface cracking (Cook and Pharr, 1990), and are not recorded in the CR test. This suggests that RCS, along with the accompanied densification gradient, enhances the CR by obstructing sub-surface cracking to reach the surface. Chipping, defined as sub-surface cracking reaching the surface and is manifested by giving a crater; see Figure 11, then occurs at sufficiently high load, *q.v.*, Figure 8 insets for 10N of Al16-CS and Al20-CS. Chipping is then taken to imply  $PCI = 100\%$  if it affects all sides of the indent. Chipping was clearly observed for  $\text{Al}_2\text{O}_3$  contents  $>10$  mol%. This reasoning is confirmed by a recent study by (Kang et al., 2020) where they observed median crack initiation upon indentation, likely beneath the RCS zone. However, their  $\text{Al}_2\text{O}_3$  content is not mentioned in the paper. Xiaoyu *et al.* concluded that radial cracks are more easily initiated in CS AS glass with about 5 wt%  $\text{Al}_2\text{O}_3$  (Xiaoyu et al., 2018), which is in the range of 3–4 mol%; see Table 1, and thus agrees with our results.

As already stated, chipping was for  $\text{Al}_2\text{O}_3 > 10$  mol% frequently observed to occur during unloading. In thin glass central tensile stresses from the balancing RCS could potentially induce such behavior, however, our monoliths are quite thick (ca. 2–2.5 cm), so we expect the central tension to be marginal.

Instead, the clear difference in CR is governed by the  $\text{Al}_2\text{O}_3$  content.  $\text{Al}_2\text{O}_3 > 10$  mol% reduces the “fast” network relaxation by having a favorable structure, i.e., low alkali-induced NBO structure, a high reduced elastic modulus and lower  $C_g$  (see Figure 10B). We already implied in Section 4.1, that lower  $C_g$  is linked to a higher ability for densification (Rouxel et al., 2008; Varshneya et al., 2015) and thus a higher ability for RCS buildup which also increases the CR; *q.v.*,  $\sigma_s$  data in Supplementary Table S1. In Figures 10A,B linear relation between CR and  $C_g$  is shown,  $R^2$  is 0.98 and 0.92 for 5 and 16h respectively. In addition, viscous relaxation of the RCS is also inhibited by the higher viscosity that the  $\text{Al}_2\text{O}_3$  content gives (Karlsson, 2021).

Varshneya previously described ion-exchange strengthening to be analogous to indentation (Varshneya, 2010b; a), having both elastic and plastic processes that governs the buildup of RCS. It was predicted that the compositions that exhibit less indentation depth and thus those having higher elastic modulus should perform better as CS glass. For our CS glasses the indentation depth at max load,  $h_m$ , as a function of the  $\text{Al}_2\text{O}_3$  content is decreased; *q.v.*, Figure 12, which supports the prediction of Varshneya. The indentation depth,  $h_m$ , follows linear but opposite trends as  $H$  as a function of  $\text{Al}_2\text{O}_3$  content. In a previous paper, a rough correlation between the ratio  $E_r/H$  and CR was found (Karlsson, 2022). It is more frequently used in its inverted form,  $H/E_r$ , where materials having  $H/E_r > 0.1$  exhibit excellent tribological properties (Musil et al., 2002; Calahoo et al., 2016; Varshneya et al., 2022).  $H/E_r$  can be directly linked to the elastic recovery energy ( $U_E/U_T$ ) in the indentation process by a proportionality factor,  $\kappa$ , which can be approximated to 5.17 in

the range 0.08–0.12 for  $H/E_r$  (Yetna N'Jock et al., 2016).  $U_E/U_T$  upon indentation is given by

$$\frac{U_E}{U_T} = \kappa \frac{H}{E_r} \quad (7)$$

and the calculated values for  $U_E/U_T$  are given in [Supplementary Table S1](#) together with  $H^3/E_r^2$ , which instead is a measure of the resistance to plastic deformation (Calahoo et al., 2016). *N.b.*, that we assume the same  $\kappa$  also for those samples having  $H/E_r > 0.12$ . Both  $U_E/U_T$  and  $H^3/E_r^2$  show similar trends as  $H$  and  $h_m$  as a function of  $\text{Al}_2\text{O}_3$  content, cf. [Supplementary Figure S3](#). Thus,  $U_E/U_T$  and  $H^3/E_r^2$  also show similar trends as a function of  $CR$ ; see [Supplementary Figure S3](#). Thus,  $H$  and  $E_r$  are keys for understanding  $CR$  through the previously described relations. Indentation induced cracking is governed by mismatching stresses upon loading or unloading and a better balance between elastic and plastic dissipation should give higher  $CR$ . For the studied glasses it occurs when  $U_E/U_T > 57\%$ . An analysis of the indentation stress fields, and its effect of RCS in CS glasses would be useful to elucidate more on this.

The crack initiation mode, radial or median, can be dependent on the composition apart from the RCS. For  $\text{Al}_2\text{O}_3 < 10$  mol% there is sub-surface cracking but also much easier radial cracking. However, the  $CR$  test only considers the radial crack statistics and is thus unfavorable for the CS glasses with  $\text{Al}_2\text{O}_3 < 10$  mol%, as radial cracks are still easily formed (like isotatically compressed glasses (Zhang et al., 1995; Svenson et al., 2014)). The radial crack lengths for the CS glasses are clearly shorter, *q.v.* [Figure 8](#), consistent with observations in borosilicate CS glasses (Talimian and Sglavo, 2017). However, this measure was not considered in the present study.

## 5 Summary and conclusion

Our study on Al-doped soda-lime-silicate glasses upon  $\text{K}^+$ -for- $\text{Na}^+$  exchange manifested a complex property-dependence on their relative Si and Al contents.  $T_g$ ,  $H$  and  $E_r$  of the pristine glasses were enhanced monotonically for increasing  $\text{Al}_2\text{O}_3$  content.  $H$  and  $E_r$  showed for  $M_{(2)}\text{O}/\text{Al}_2\text{O}_3 \approx 1$  nonlinear trends, whereas the CS counterparts were found to be strictly linearly increasing with the Al-for-Si replacement.  $T_g$  and  $\sigma_s$  on the other hand, was linear for the pristine glasses and linear to nonlinear for the CS glasses as a function of  $\text{Al}_2\text{O}_3$  content. Several physical properties, along with the silicate network connectivity, of both series of pristine and chemically strengthened glasses depended roughly linearly on their alumina content, except for clear deviations observed for glasses around the  $M_{(2)}\text{O}/\text{Al}_2\text{O}_3 \approx 1$  composition.

$^{27}\text{Al}$  MAS NMR experiments revealed a dominance of  $[\text{AlO}_4]^-$  groups throughout all pristine and ion-exchanged glass structures, with the exchange manifested only by a minor  $^{27}\text{Al}$  deshielding upon the replacement of  $\text{Na}^+$  by the less

electronegative  $\text{K}^+$  cations. Significant alterations were observed in the  $^{23}\text{Na}$  shifts observed by  $^{23}\text{Na}$  MAS NMR, which presumably reflects an overall lowering of the (average) Na coordination numbers because the larger  $\text{K}^+$  cations outcompete those of  $\text{Na}^+$  for the highest-coordination sites in the glass structure. The extent of cation exchange was also estimated by  $^{23}\text{Na}$  MAS NMR, revealing a markedly higher degree of exchange ( $\approx 74\%$ ) of the Al-free SLS glass relative to the Al-bearing glasses, all of which revealed an essentially identical exchange level ( $\sim 65\%$ ) of the  $\text{Na}^+$  cations; this observation is attributed to the more polymerized networks of the Al-bearing glasses and the preference for the monovalent  $\text{Na}^+$  cation to balance the  $[\text{AlO}_4]^-$  groups, which naturally impedes its mobility and thereby exchange by  $\text{K}^+$ .

Differential thermal analysis data revealed that the ion-exchanged glasses exhibit an exothermic step just below a blurred  $T_g$ . We propose that these characteristics stem from a combination of stress relaxation of the residual compressive stresses induced by the  $\text{K}^+$ -for- $\text{Na}^+$  ion exchange.

The chemical strengthening consistently increased the hardness, except for glasses incorporating  $< 10$  mol%  $\text{Al}_2\text{O}_3$  for the specific conditions of an 16 h exchange period and an applied load of 1 mN, which resulted in a lower hardness than that of the pristine glass. We attribute these observations to a combination of stress relaxation and the mixed alkali effect. Stress relaxation is also manifested by the fact that the 16 h treatment time consistently yielded lower  $H$  and  $E_r$  values relative to glasses treated for 5 h. These effects were less obvious for the Al-rich glasses with  $n(M_{(2)}\text{O})/n(\text{Al}_2\text{O}_3) \approx 1$  and on the  $\sigma_s$  results regardless of composition.

The crack resistance of the pristine glasses becomes slightly lower with increasing  $\text{Al}_2\text{O}_3$ . The chemical strengthening has a pronounced enhancement on the crack resistance for glasses with  $n(M_{(2)}\text{O})/n(\text{Al}_2\text{O}_3) \approx 1$ . The chemical strengthening increases the elastic energy of recovery and the resistance to plastic deformation. The observed linear correlation between the crack resistance and the compactness suggests that  $C_g$  is a key parameter for the buildup of compressive stresses in the vitreous structure. Thus, our results suggest that the crack resistance ( $CR$ ) of a chemically strengthened glass is a property governed by both the composition and structure, as well as residual compressive stress, where the  $CR$  is favored by a polymerization increase and compactness decrease; the latter is also linked to the residual compressive stress increase.

## Data availability statement

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation. The

datasets generated for this study can be found in the Swedish Service Repository <https://snd.gu.se/en>. DOI: <https://doi.org/10.5878/kp36-vd61>.

## Author contributions

SK: conceptualization, methodology, validity, formal analysis, investigation, resources, data curation, writing - original draft, writing - review and editing, visualization, project administration, funding acquisition. RM: investigation, data curation, writing - review and editing. SA: investigation, data curation, writing - review and editing. MP: methodology, formal analysis, investigation, data curation. JA: validity, resources, writing - review and editing. BS: formal analysis. ME: methodology, validity, formal analysis, investigation, resources, writing - original draft, writing - review and editing, visualization.

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## Conflict of interest

Authors MP and JA were employed by GlasStress Ltd.

The remaining authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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## Supplementary material

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fmats.2022.953759/full#supplementary-material>

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