



High-Sensitivity RF Choke-Enhanced Dipole Antenna-Coupled Nb₅N₆ THz Detector

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In this study, we demonstrate an Nb₅N₆ terahertz (THz) detector with radio frequency (RF) choke-enhanced dipole antenna structure for 0.3 THz detection. The maximum electric field intensity of 218 V/m is obtained by optimizing the parameters of the dipole antenna with RF choke. Compared to a dipole antenna without RF choke, the electric field intensity of that with RF choke is improved by 2.6 times. The RF choke-enhanced dipole antenna-coupled Nb₅N₆ THz detector is fabricated and characterized. The measured maximum responsivity of the detector is 1100 V/W at 0.308 THz, and the corresponding noise equivalent power (NEP) is 6.4×10^{-12} W/Hz^{1/2}. The measured response time of the Nb₅N₆ THz detector is as low as 8.46 µs. Furthermore, the Nb₅N₆ THz detector is applied to a homemade THz transmission imaging system for demonstrating its performance. The THz imaging results of a blade and access card show that the contrast of the blade image is sharp and the components hidden within the access card are clearly visible. This indicates that the Nb₅N₆ THz detector can be used in THz imaging, particularly in THz active imaging, which will have greater application prospects.

Keywords: Nb_5N_6 microbolometer, THz detector, RF choke, antenna, THz imaging

INTRODUCTION

In recent years, with the development of science and technology, the terahertz (THz) waves have attracted worldwide interest. THz technology is applicable in many fields such as physics, life science, materials science, astronomy, information technology, and national defense security. THz detectors are crucial components in terahertz technology, which have extensive application prospects in spectrum analysis, medical imaging, radio astronomy, security checks, biosensing, etc. [1-6]. At present, the research on uncooled THz detectors involve Golay cells [7], Schottky diodes [8], fieldeffect transistors (FETs) [9], and bolometers [10–12]. Among them, bolometers have been attracting more and more attention because of easy fabrication, integration with readout circuits, high sensitivity, low-noise equivalent power (NEP), wide response spectrum, and so on. To date, several materials have been used as the bolometers, such as VO_x [13-16], YBCO [17-19], titanium (Ti) [20], niobium (Nb) [10, 21], platinum (Pt) [22], poly-Si [23], poly-Si-Ge alloys [24], etc. Although polycrystalline VO_x exhibits a temperature coefficient of resistance [TCR = (1/R) $(dR/dT) \sim -2\%/K$ at room temperature, it undergoes a semiconductor-metal transition at ~ 67°C, which reduces the dynamic range of the bolometer when working in the semiconductor phase [25]. Amorphous YBCO has been demonstrated as an uncooled bolometer, but its stoichiometric composition is difficult to control, and the resistance and TCR are affected by the oxygen

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FIGURE 1 (A) Schematic diagram of the planar dipole antenna on the Si substrate. (B) The location and size of the Nb₅N₆ film in the planar dipole antenna. (C) The electric field intensity distribution of the half-wave dipole antenna with $L_1 = 164 \mu m$ at 0.295 THz. (D) The electric field intensity varies with L_1 in the half-wave dipole antenna in the range of 0.25–0.35 THz. (E) The electric field intensity distribution of the full-wave dipole antenna with $L_1 = 164 \mu m$ at 0.295 THz. (D) The electric field intensity varies with L_1 in the half-wave dipole antenna in the range of 0.25–0.35 THz. (E) The electric field intensity distribution of the full-wave dipole antenna with $L_1 = 404 \mu m$ at 0.301 THz. (F) The electric field intensity varies with L_1 in the full-wave dipole antenna in the range of 0.25–0.35 THz.

concentration in films [26]. The low TCR value of metallic materials (two orders of the magnitude lower than that of semiconductors) results in low responsivity of detectors. In polycrystalline Si and Si-Ge alloys, the resistivity can be maintained sufficiently low by controlling the doping level, although it limits the TCR to approximately 2.5–3%/K. However, the processing of polycrystalline materials requires temperatures as high as 650°C for achieving the desired crystallinity [23, 24]. Therefore, it is desirable to combine the advantages of both semiconductors (high TCR for high responsivity) and metals (low resistivity) for the development of fast and sensitive THz detectors.

The Nb₅N₆ films made in our laboratory exhibit a high TCR of 0.7%/K, which is higher than that of metals [27–29]. In addition,

the resistivity of $Nb_5N_6~(<5\times10^{-3}~\Omega~cm)$ is considerably lower than that of the other semiconductors, resulting in reduction of the impedance in comparison with a-Si or VO_x devices of the same dimensions, which is advantageous to impedance matching when integrated with planar antennas. Moreover, the Nb_5N_6 films include other advantages, such as simple fabrication, low cost, large area production, and easy control. Hence, Nb_5N_6 thin films are promising candidate material for THz bolometer detector at room temperature.

Generally, a THz bolometer detector consists of a THz antenna and a thin film microbridge. As the key component of the detector, the antenna plays an important role in coupling THz radiations to the film microbridge; hence, the performance of the detector can be improved by appropriately designing and optimizing the antenna structure. Radio frequency (RF) choke structures are often used in the antenna design for improving the antenna performance, such as bandwidth enhancement, multiple band operation, gain enhancement, and radiation pattern shaping [30–33]. Their functions are usually to provide direct current (DC) bias and prevent RF signals from leaking to the bias electrodes. Considering the feature of the RF choke, we designed an RF choke-enhanced dipole antenna to improve the performance of the Nb₅N₆ THz detector. Compared to that without RF choke, the electric field intensity of the RF choke-enhanced dipole antenna to improved. In addition, we applied the Nb₅N₆ THz detector to a THz transmission imaging system and studied the THz imaging of a blade and access card.

Design and Fabrication

We choose 0.3 THz as the center frequency of the dipole antenna for design and optimization, because it is located in the atmospheric transmission window of THz signals, and its loss is small in the atmosphere. To maximize the response of dipole antenna under the 0.3 THz radiation, the antenna length must match the effective wavelength of the surrounding media. The length *l* of the resonant antenna [34] is given by $l = \lambda / \sqrt{\epsilon_{eff}}$, where λ is the free-space wavelength and ε_{eff} is the effective dielectric constant; within the effective media approximation, ε_{eff} is the average of the dielectric constants of Si and air, which are 1 and 11.9, respectively. The calculated effective wavelength of the dipole antenna on high resistivity Si is 394 µm for 0.3 THz detection. The geometry of the dipole antenna is depicted in Figure 1A. The antenna was placed on a thermally oxidized Si substrate with a period of $800 \,\mu\text{m} \times 800 \,\mu\text{m}$, located in the center of the coordinate system. The thickness of the Si substrate, the SiO₂ on the Si surface, and Au layer were 500, 200, and 200 nm, respectively. A 200 nm thick Au layer was provided on the back of the Si substrate as a Fabry-Perot (F-P) cavity to further enhance the coupling efficiency of the antenna to THz signals. In the simulation model, the X and Y directions were both set as periodic boundary conditions, whereas the $\pm Z$ directions were both set as open (add space) boundary conditions. The plane wave excitation was setup to simulate the normal incidence of a uniform plane wave on the antenna and its incidence direction was defined as along the Z-axis. As the wave polarized along the X-axis was of our interest, we defined the plane wave excitation amplitude of 1 V/m on the X-axis, while zero on the Y-axis and Z-axis. An electric field monitor probe was placed in the center of the antenna to detect the changes in electric field intensity of the antenna structure (red point in Figure 1A). Therefore, the electric field enhancement performance of the antenna could be obtained by comparing the probe field strength with the plane wave excitation amplitude. The length and width of the dipole antenna were L_1 and 10 µm. To integrate the dipole antenna with the Nb₅N₆ microbolometer THz detector, the Nb₅N₆ film was located in the center place of the antenna indicated by the brown area in Figure 1B. The size of the Nb₅N₆ film microbridge was 3 μ m \times 12 μ m, and the length of the transition area between the film microbridge and dipole antenna was 10 µm. The size of the electrodes (pads) used for providing bias current to the THz

detector was set to $150 \,\mu\text{m} \times 300 \,\mu\text{m}$. According to microwave antenna theory, the electric field intensity changes with frequency are simulated when the length of the dipole antenna is approximated half wavelength, as shown in Figure 1D. In the 0.25-0.35 THz frequency range, the electric field intensity initially increases and then decreases with the increase of L_1 . The electric field intensity attains maximum value of 75 V/m at 0.2945 THz when L_1 is 164 µm. Figure 1C shows the electric field intensity distribution of the half-wave dipole antenna near the surface of Si substrate at 0.2945 THz when $L_1 = 164 \,\mu\text{m}$. The electric field is mainly distributed in the gap area of the half-wave antenna where the Nb₅N₆ film is located. This shows that the Nb₅N₆ film can receive maximum THz energy from the antenna coupling. Additionally, we also simulate the changes in the electric field intensity of the full-wave antenna in the frequency range of 0.25-0.35 THz, as shown in Figure 1F. It can be observed that as L_1 increases, the electric field intensity shows a trend of first increasing and then decreasing. When $L_1 = 404 \,\mu\text{m}$, the best electric field intensity of the full-wave antenna is 85 V/m at 0.301 THz. The electric field intensity of the full-wave antenna is 10 V/m greater than that of the half-wave antenna. The electric field intensity distribution of the full-wave antenna near surface of Si substrate at 0.301 THz is shown in Figure 1E. Compared to the half-wave antenna, the full-wave antenna has a higher local field enhancement. It indicates that the full-wave antenna has higher coupling efficiency to THz radiation signals. However, there is also strong electric field distribution on the dipole antenna and pads.

To further enhance the electric field intensity at the center of the dipole antenna, an RF choke structure was employed to prevent the coupled THz waves from leaking to the bias pads, as depicted in Figure 2A. By optimizing the location and the size of the RF choke, a maximum electric field intensity of 218 V/m at the center of the dipole antenna with RF choke is obtained at 0.301 THz, as shown in Figure 2B. The optimized parameters of the dipole antenna with RF choke are as follows: $L_2 = 66 \,\mu\text{m}$, $W_1 =$ 30 μ m, and $W_2 = 230 \mu$ m. Compared to the dipole antenna without RF choke, the maximum electric field intensity is improved by 2.6 times. The electric field intensity distribution of the antenna with RF choke near the surface of the Si substrate at 0.301 THz is displayed in Figure 2C. The antenna gap region has greater local field enhancement due to the suppression effect of RF choke structure. There is almost no electric field distribution on the bias pads, and the electric field is mainly concentrated between RF choke structures, which makes the electric field intensity at the center of the dipole antenna greater. When the RF choke-enhanced dipole antenna is integrated with the Nb₅N₆ microbolometer THz detector, the Nb₅N₆ film can absorb more THz energy, thus improving the sensitivity of the device. Figure 2D shows the far-field radiation patterns of the dipole antenna with RF choke in the E plane (Phi = 0°) and H plane (Phi = 0°) at 0.301 THz, with a maximum directivity of 8.34 dBi. This indicates that the dipole antenna with RF choke has a good property of receiving the normal incidence THz waves according to the reciprocity principle.

The fabrication of the RF choke-enhanced dipole antennacoupled Nb_5N_6 microbolometer THz detector was implemented







using micromachining technique. The process flow is illustrated in **Figure 3**. The Nb_5N_6 film with a thickness of 120 nm was first deposited on the thermally oxidized Si substrate by using RF magnetron sputtering, as shown in **Figure 3A**. The TCR of the Nb₅N₆ thin film is -0.7% K⁻¹ at 300 K. After the required pattern was formed through micro-processing technology, 5 nm thick Ti



membrane and a 200 nm thick Au membrane were deposited on the formed pattern (Figure 3B) and the Ti membrane function as an adhesive layer in this process. Further, the Au antenna pattern was formed by using the lift-off technology, as shown in Figure 3C. Subsequently, the Nb₅N₆ thin film was patterned into Nb₅N₆ film microbridge using lithography and reactive ion etching (RIE) technique (Figure 3D). The resistance of the Nb5N6 microbolometer THz detector depends on the dimensions of the Nb5N6 film microbridge. Finally, a photoresist protection pattern was generated on the formed pattern using the same lithography technique, with only two open areas on both sides of the Nb5N6 film microbridge (Figure 3E). After the photoresist pattern was formed, the SiO₂ layer on the surface of the Si substrate was dry-etched through RIE to expose Si pattern only in the two open areas. The RIE was performed in a gas mixture of CF_4 and O_2 at the total pressure of 4 Pa and RF power of 150 W for 3 min. The optimized mixture ratio of CF4 and O2 was 30:10. The opening formed on Si was further etched by using RIE to create an air-cavity (Figure 3F). The etching conditions are as follows: the flow of SF₆ is 40 sccm, chamber pressure is 8 Pa, and the RF power is 70 W. The air-cavity was formed under the Nb₅N₆ film microbridge with a SiO₂ supporting layer through RIE etching, because the etching rate of Si was much higher than that of SiO₂. In the etching process of the Si air-cavity, we adopted the technology of stepwise etching, etching for 3 min each step, and etching for three steps in total. The advantages of stepwise etching are that, on the one hand, it is conducive to directly observe whether the air-cavity under the Nb₅N₆ film microbridge is formed; on the other hand, when the air-cavity is etched, the photoresistance on the surface of Si substrate will be heated as little as possible to prevent it from solidifying, which is conducive to the removal of photoresistance in the later stage.

The optical image of the fabricated RF choke-enhanced dipole antenna-coupled Nb_5N_6 THz detector and the SEM micrograph of an Nb_5N_6 air-bridge are displayed in **Figure 4**. The fabricated Nb_5N_6 THz detector includes four main parts: a planar dipole antenna, Au bias pads, an Nb_5N_6 film microbridge, and an etched air-cavity. In **Figure 4B**, the Nb_5N_6 film microbridge is thoroughly freestanding from the Si substrate with a 200 nm SiO₂ support layer. The SiO₂ layer mainly plays a role in preventing the Nb₅N₆ film microbridge from fracture or collapse. The air-cavity isolates the heat conduction between the Nb₅N₆ film microbridge and the Si substrate, which further improves the sensitivity of the THz detector. The DC resistance of the fabricated Nb₅N₆ THz detector is 0.76 kΩ.

CHARACTERIZATION OF THE NB₅N₆ THZ DETECTOR

Typical performance parameters of THz detectors are responsivity (\Re) and NEP. Responsivity is related to the efficiency of THz detector and is defined as the RF or optical response for a given input power. It can be expressed as follows [35, 36]:

$$\mathfrak{R} = \frac{\Delta U}{P_{in}},\tag{1}$$

$$P_{in} = \frac{A_{eff}}{S_{beam}} \cdot P_{beam},\tag{2}$$

where $\triangle U$ is the output voltage of the Nb₅N₆ THz detector, P_{in} is the input power received on the effective area of the Nb₅N₆ THz detector, and S_{beam} and P_{beam} are the spot size and total power of the THz beam at the focus, respectively. The measured total power of the THz beam at the focus is 0.5 mW.

Measuring the power received by the detector is essential to calculate responsivity and NEP, but this power is a function of effective area A_{eff} and impinging of received power density. The effective area A_{eff} is considered as the pixel area that effectively captures the impinging power and it is a function of receiving antenna directivity and signal frequency. Thus, if D is the directivity of receiving antenna, the effective area of the THz detector is given by [37, 38]

$$A_{eff} = \frac{\lambda^2}{4\pi} \cdot D, \tag{3}$$



where λ is the signal wavelength. For a simulated value of D = 8.6 dBi and free-space wavelength $\lambda = 0.97$ mm at 0.301THz, a value of $A_{\text{eff}} = 0.516$ mm² is obtained.

To calculate the power on the effective receiving area of the Nb_5N_6 THz detector, it is necessary to determine the spot size of the THz beam on the focal plane. The spot distribution of the THz beam at the focus is scanned using the Nb_5N_6 THz detector. At a scan range of 20 mm × 20 mm and the step size of 210 μ m, the measured spot distribution of the THz beam at the focus is depicted in **Figure 5.** Figure 5A is the three-dimensional distribution of the THz beam scanned by the Nb_5N_6 THz detector. The ordinate denotes the normalized response voltage value, and the X-axis and Y-axis denote the scanning distance of the Nb_5N_6 THz detector, respectively. A scanned image of the THz beam at the focus is displayed in Figure 5B. It can be seen that the spot diameter of the THz beam is about 10 mm.

The used THz source is composed of AMC-336 frequency multiplier produced by Virginia Diodes, Inc. (VDI) and Agilent E8251D signal generator. A low-frequency signal provided by the signal generator is amplified by AMC-336 frequency multiplier (WR2.8 VDI horn) to obtain the THz frequency in the range of 0.25-0.35 THz. The modulation signal of the signal generator is 1 kHz, which is input to the lock-in amplifier as a reference signal. The maximum output power of THz source is about 0.5 mW measured by a thermal power sensor (OPHIR, 3A-P-THz). The THz radiation is focused to maximize the electrical signal of the detector by two off-axis parabolic mirrors and the threedimensional displacement platform. The responses voltage of the detector is read out by a lock-in amplifier (SR830). The responsivity of the Nb5N6 array detector depends on its voltage value and is calculated by Eqs 1-3. The measured optical voltage responsivity of the Nb₅N₆ THz detector varies with frequency, as shown in Figure 6. The resonant frequency corresponding to the maximum optical responsivity of the Nb5N6 THz detector is 0.308 THz, and the best optical responsivity is 1100 V/W. The



measured responsivity varying with frequency is the same as the variation trend of simulated electric field intensity with the frequency. A small frequency shift is observed at the resonance frequency. This deviation may be due to the approximate simulation model, in which we neither considered the effect of the air-bridge nor the size deviation in the fabrication process.

Besides the RF responsivity, the NEP is crucial parameter for characterizing the sensitivity of THz detector, which can be expressed as

$$NEP = \frac{V_{nosie}}{\Re},$$
 (4)

where V_{noise} is the voltage noise spectral density and \Re is the optical responsivity of the Nb₅N₆ THz detector.



FIGURE 7 | (A) Measured noise voltage under different bias currents and (B) the calculated NEP as a function of the modulation frequency for the Nb₅N₆ THz detector.



The voltage noise spectral density of the Nb₅N₆ THz detector was obtained using a spectrum analyzer (Agilent 35670A) and a low-noise amplifier (LNA-SR560) with a voltage gain of 60 dB, as shown in **Figure 7A**. The noise voltage of the Nb₅N₆ THz detector increases with the bias current. Note that the 1/*f* noise cross point also increases with the increase of the bias current. At lower bias current (0 and 0.2 mA) and modulation frequencies greater than 1 and 4 kHz, the noise spectrum tends to be constant and the device is limited by the thermal noise ($\sqrt{4kTR}$), where *k* is Boltzmann's constant (1.38 × 10⁻²³ J/K), *T* is the Kelvin temperature, and *R* is the DC resistance of the device. However, at 0.6 mA bias current, the 1/*f* noise-decay increases to 12 kHz. This indicates that the low-frequency flicker noise (1/f noise) is related to the bias current, as explained in [17]. The measured noise voltage of the Nb5N6 THz detector with 0.6 mA bias current at 1 kHz is about 17 nV/Hz^{1/2}, and the NEP is 1.5×10^{-11} W/Hz^{1/2}. As shown in Figure 8, the response voltages of the Nb₅N₆ THz detector are a function of modulation frequency of the THz radiation signals. Therefore, the NEP of the Nb₅N₆ THz detector is also related to the modulation frequency of the THz radiation signals. The measured NEP of the Nb₅N₆ THz detector changes with the modulation frequency at the bias current of 0.6 mA, as displayed in Figure 7B. It can be observed that the NEP gradually decreases with the increase of the modulation frequency. The NEP which tends to a constant for the modulation frequency is greater than 10 kHz, and the minimum NEP of the Nb₅N₆ THz detector is 6.4×10^{-12} W/ $Hz^{1/2}$ at the modulation frequency of 12 kHz. The performance comparison between the proposed Nb5N6 THz detector and other published detectors is shown in Table 1. The table contains the relevant performance parameters of THz detector published by our research group and other research groups. Compared with other detectors, the Nb₅N₆ THz detectors also have higher sensitivity and will have greater supplication prospects in the THz imaging field. Compared with our previous work [28], the optical responsivity and NEP of the RF choke-enhanced antenna-coupled Nb5N6 THz detector are increased by 2 times and 3 times, respectively.

The thermal time constant of the THz detector, which determines the speed of image processing, plays an important role in THz imaging technology. The shorter the response time, the faster the image acquisition speed. The measured normalized response voltage of the Nb₅N₆ THz detector as a function of modulation frequency is depicted in **Figure 8**. The normalized amplitude of the response voltage decreases with the increase of the modulation frequency. The measured data can be well fitted by the formula $1/\sqrt{1 + (2\pi f \tau)^2}$, indicating that the detector is well described

TABLE 1	Performance compari	son between the prop	osed Nb∈Ne THz dete	ctors and other detectors.

Detector	Antenna types	Detection frequency (THz)	Responsivity (V/W)	NEP (pW/Hz ^{1/2})	Reference
YBa ₂ Cu ₃ O _{7-x} THz detector	Spiral antenna	0.4	70	50	[17]
Graphene FET THz detector	Bowtie antenna	0.8	535	200	[39]
Graphene FET THz detector	Dipole antenna	0.33	30	51	[40]
GaAs Schottky THz detector	Bowtie antenna	0.22	1,650	3.6	[41]
Nb ₅ N ₆ THz detector	Double-slot antenna	0.65	113	44	[42]
	Dipole antenna	0.28	580	17	[28]
	RF choke-enhanced dipole antenna	0.3	1,100	6.4	This work



by the classical thermal model, where the bolometer is represented by an absorber with heat capacity connected to a thermal bath through a small thermal conductance. By fitting the measured response voltage under different modulation frequencies, the thermal time constant of the Nb_5N_6 THz detector was determined to be 8.46 µs.

APPLICATION OF THE NB₅N₆ THZ DETECTOR

To demonstrate the feasibility of the Nb_5N_6 THz detector in practical application, we applied it in a THz transmission imaging system. **Figure 9** shows the schematic diagram of the transmission imaging system. A pair of off-axis parabolic mirrors is used to collimate and focus the modulated THz radiation, a twodimensional displacement platform is used to move the object to realize the scanning, the THz radiation signals transmitted through the object are detected by the Nb_5N_6 THz detector, the final voltage signals are readout from a lock-in amplifier (SR830) to the computer, and the images of the objects are recovery from these measured raw data without image modification.

Because of the low absorption of THz signals by plastics and the complete reflection of THz signals by metals, THz imaging technology is a very effective way for detecting metal or voids in solid plastic products and other objects containing different THz absorption coefficients. Two objects with different characteristics were used to achieve their images at 0.308 THz with the developed THz imaging system, as shown in Figure 10. The first object is a metal blade, and the THz radiation will fully pass through it or not depending on its shape, resulting in that the transmitted THz radiation at different position has nearly the unity intensity. Figure 10A shows the image of the blade; the red and the black parts represent the areas where the THz radiation can and not transmit, respectively. Details of the blade imaging, such as the edges and irregular holes, are extremely clear. The second object is an access card, whose components inside are more complicated than the blade, resulting in transmitted THz radiation at different part with different intensities. Figure 10B displays the image of the components hidden inside the access card, such as a coil and rectangular chip, which can be observed clearly. These results indicate that the performance of the Nb5N6 microbolometer THz detector



integrated with RF choke-enhanced dipole antenna is good from another perspective and it is sufficient for practical applications.

CONCLUSION

The Nb₅N₆ THz detector integrated with RF choke-enhanced dipole antenna was developed and demonstrated for 0.3 THz detection and imaging at room temperature. The simulated electric field intensity of the proposed RF choke-enhanced dipole antenna showed an improvement of 2.6 times compared to that without RF choke. The measured best optical responsivity of the Nb5N6 microbolometer THz detector integrated with the proposed antenna was 1100 V/W and the corresponding minimum NEP was 6.4×10^{-12} W/Hz^{1/2} at 0.308 THz with 0.6 mA bias current. The response time of the Nb5N6 THz detector was as low as 8.46 µs. The measurement results were in good agreement with the simulation ones. Furthermore, a THz imaging system based on transmission mode was developed using the proposed detector, and the THz imaging of a blade and the access card was investigated. The results show that the obtained THz imaging has high resolution and a clear outline without any image processing. Especially, through the analysis of the transmission image of the access card, its internal structure could be clearly observed. It was demonstrated that the proposed antenna-coupled Nb5N6 THz detector had high sensitivity and fast response speed for THz imaging. The

development of large-scale detector arrays for terahertz imaging using the proposed detector in this work is on the underway.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/Supplementary Material; further inquiries can be directed to the corresponding authors.

AUTHOR CONTRIBUTIONS

CJ designed and simulated the structure. All authors analyzed and discussed the results. CJ and XT wrote the manuscript. The project was conceived and supervised by XT and LK.

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